

Figure 41-1. Real and imaginary part of the RPA self-energy for three wave vectors, in units of the plasma frequency. The chemical potential is included in $\text{Re}\Sigma$. The straight line that appears on the plots is $\omega - \epsilon_k$. Taken from B. I. Lundqvist, *Phys. Kondens. Mater.* **7**, 117 (1968). $r_s = 57$

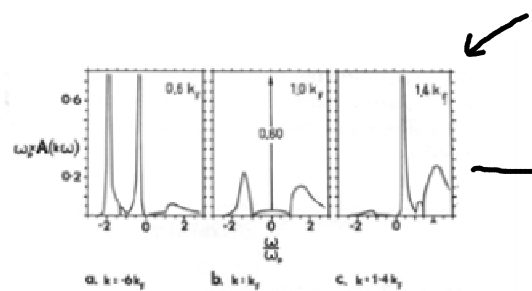


Figure 41-2. RPA spectral weight, in units of the inverse plasma frequency. Taken from B. I. Lundqvist, *Phys. Kondens. Mater.* **7**, 117 (1968).

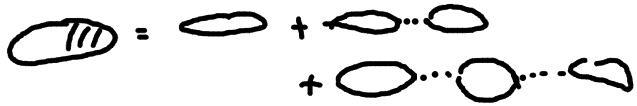
41. Physique: propriétés à 1 particule

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- 2. Interp. physique
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5. Comparaison avec expérience

VI. Fermions sur réseau: Hubbard et Mott

- 44. Densité fonctionnelle
 - 1. Fondamental est une fonctionnelle de $n(\vec{r})$
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 - 1. $U=0$ ←
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- 46. TPSC - ACDP
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41.2 Interp. physique



$$\Sigma(k, i\hbar_n) = \Sigma_{HF} + \int \frac{d^3q}{(2\pi)^3} T \sum_{i q_n} V_q \chi_{nn}(q, i q_n) V_q \frac{1}{i\hbar_n + i q_n - \epsilon_{k+q}}$$

$$= \Sigma_{HF} + \int \frac{d^3q}{(2\pi)^3} T \sum_{i q_n} \int \frac{d\omega'}{\pi} \frac{V_q \chi''_{nn}(q, \omega') V_q}{\omega' - i q_n} \frac{1}{i\hbar_n + i q_n - \epsilon_{k+q}}$$

$$T \sum_{i q_n} \frac{1}{\omega' - i q_n} \frac{1}{i\hbar_n + i q_n - \epsilon_{k+q}} = T \sum_{i q_n} \left[\frac{1}{\omega' - i q_n} + \frac{1}{i\hbar_n + i q_n - \epsilon_{k+q}} \right] \frac{1}{\omega' + i\hbar_n - \epsilon_{k+q}}$$

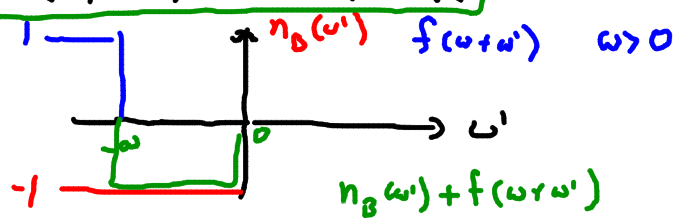
$$= \frac{n_B(\omega') + f(\epsilon_{k+q})}{i\hbar_n + \omega' - \epsilon_{k+q}} \quad \Sigma_{HF} + \int \frac{d^3q}{(2\pi)^3} \int \frac{d\omega'}{\pi} \frac{\chi''_{nn}(q, \omega') V_q^2}{i\hbar_n + \omega' - \epsilon_{k+q}} [n_B(\omega') + f(\epsilon_{k+q})]$$

$i\hbar_n \rightarrow \frac{1}{2} + i\eta$

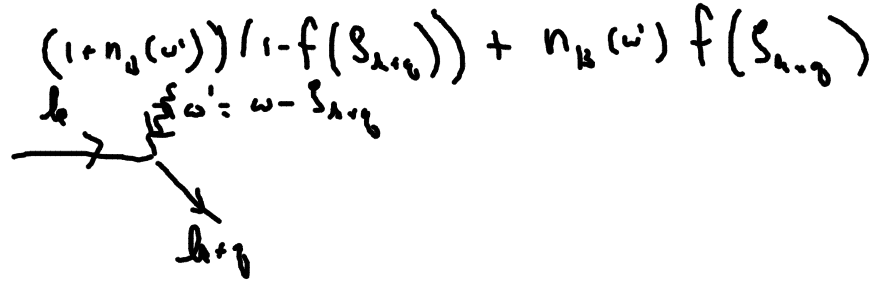
$$\Sigma''(k, \omega) = -\pi \int \frac{d^3q}{(2\pi)^3} \int \frac{d\omega'}{\pi} V_q^2 \chi''_{nn}(q, \omega') [n_B(\omega') + f(\omega + \omega')] \delta(\omega + \omega' - \epsilon_{k+q})$$

à $T=0$ et $\omega=0$ $\Sigma''(k, \omega=0) = 0 \quad \forall k$

$$n_B(\omega) = \frac{1}{e^{\beta\omega} - 1}$$



$\omega > 0 \quad \omega' \rightarrow -\omega'$



$$\delta(\omega + \omega' - \omega_{n,2}) = \delta\left(\omega + \omega' - \frac{\hbar^2 k^2}{2m} - \frac{\hbar^2 q^2}{2m} - \frac{q^2 \omega'}{2m} + \dots\right) \quad \boxed{k = k_F}$$

$$\delta(ax) = \frac{1}{|a|} \delta(x)$$

$$\Sigma^n(k_F, \omega) = -\frac{\pi}{2\pi} \frac{m}{\hbar_F} \int \frac{d^2 q_\perp}{(2\pi)^2} \int \frac{d\omega'}{\omega'} v_F^2 \left[\frac{\chi'_{nn}(q_\perp, \omega')}{\omega'} [n_B(\omega') + f(\omega + \omega')] \right] \omega'$$

$$\int \frac{d^2 q_\perp}{(2\pi)^2} \left[\frac{\chi''_{nn}(q_\perp, \omega')}{\omega'} \right] = A(k_F)$$

$$\int \frac{d\omega'}{\pi} \omega' [n_B(\omega') + f(\omega + \omega')] \xrightarrow{T=0} \int_{-\omega}^0 \frac{d\omega'}{\pi} \omega' = -\frac{\omega^2}{2\pi} \quad \frac{1}{2k_F}$$

$$= -\frac{1}{2\pi} [\omega^2 + (nT)^2]$$



41.4 Energie libre

$$Z = \text{Tr} \left[e^{-\beta(\hat{T} + \lambda \hat{V} - \mu N)} \right]$$

$$\frac{\partial F}{\partial \lambda} = \langle V \rangle_\lambda \Rightarrow F = F_0 + \int_0^\lambda d\lambda' \frac{1}{\lambda'} \langle \lambda' V \rangle_{\lambda'}$$

$$\langle \lambda V \rangle = \int d1 d2 \langle T_r \psi^\dagger(1) \psi^\dagger(2) \lambda V(1-2) \psi(2) \psi(1) \rangle_\lambda$$

$$\text{eye} \rightarrow \sum_q V_q X_{nn}^\lambda(q)$$

$$X_{nn}^\lambda(q) = \frac{X_{nn}^0}{1 + \lambda V_q X_{nn}^0}$$

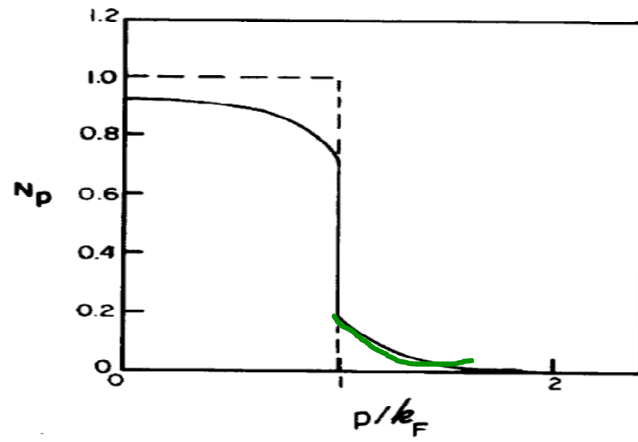
$$E_{\mathbf{k}} = \varepsilon_{\mathbf{k}} - 0.17 r_s (\ln r_s + 0.2) \frac{k_F k}{2m} + \text{cst}$$

appearing in this expression is now obvious

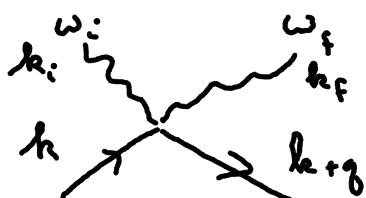
$$m^* = \frac{m}{1 - 0.08 r_s (\ln r_s + 0.2)}$$

cattering rate for $\omega = E_{\mathbf{k}} - \mu$ we find

$$\Gamma_{\mathbf{k}}(E_{\mathbf{k}} - \mu) = 0.25 r_s^{1/2} \frac{(k - k_F)^2}{2m}$$



r_s	Z_{RPA}
0	1
1	0.859
2	0.768
3	0.700
4	0.646
5	0.602
6	0.568



$$\omega_f - \omega_i = \omega$$

$$k_f - k_i = q$$

Effet Compton

$$\frac{d^2\sigma}{d\omega d\Omega} \propto \int \frac{d^3k}{(2\pi)^3} n_k \delta(\omega + \epsilon_k - \epsilon_{k+q})$$

$$\propto \int k^2 dk \int d(\cos\theta) n_k \delta\left(\omega - \frac{q^2}{2m} - \frac{kq}{m} \cos\theta\right)$$

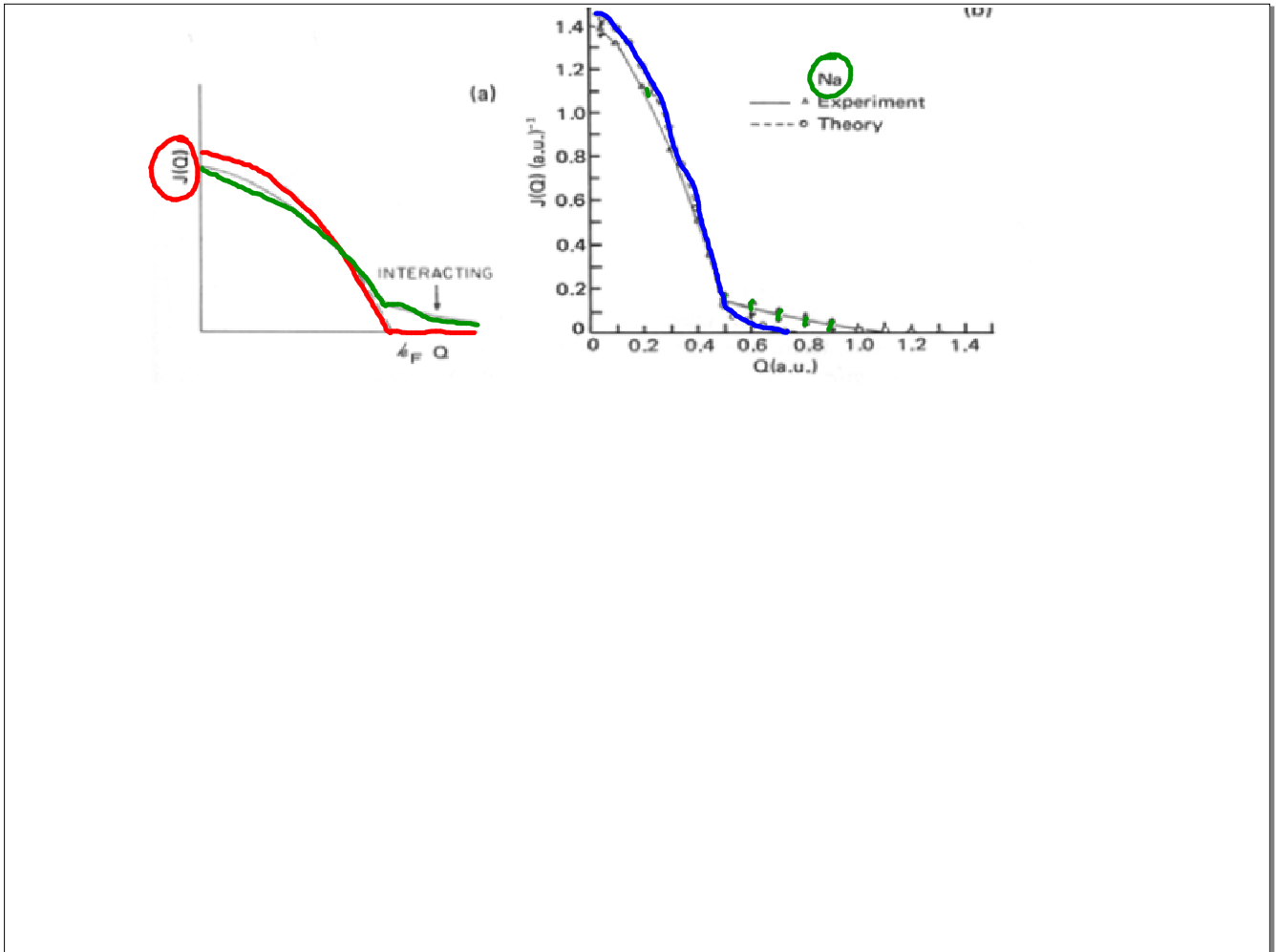
$$\int k^2 dk \int d(\cos\theta) n_k \frac{m}{kq} \delta\left(\frac{\omega - \frac{q^2}{2m}}{\frac{kq}{m}} - \cos\theta\right)$$

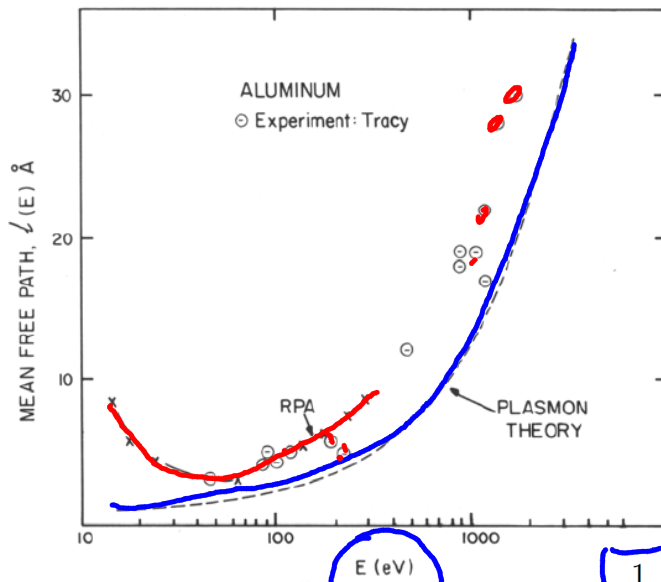
$$\frac{|Q|}{k} < 1 \quad |Q| = \left| \left(\omega^2 - \frac{q^2}{2m} \right) \frac{m}{q} \right|$$

$$\int k dk n_k \left(\frac{m}{q} \right) \Theta(k - |Q|)$$

Cas sans interaction:

$$\int_{|Q|}^{k_f} k dk n_k \frac{m}{q} = \frac{k_f^2 - |Q|^2}{2q} m \Theta(k_f - |Q|)$$





$$\frac{1}{l_{\mathbf{k}}} = \frac{\Gamma_{\mathbf{k}}}{v_{\mathbf{k}}} = \frac{1}{\tau_{\mathbf{k}} v_{\mathbf{k}}} = -\frac{2}{v_{\mathbf{k}}} \text{Im} \Sigma(\mathbf{k}, \zeta_{\mathbf{k}})$$

44. Densité fonctionnelle

1. DFT

$$E = \min_{\Psi} \left[\frac{\langle \Psi | H | \Psi \rangle}{\langle \Psi | \Psi \rangle} \right]$$

$$H = \hat{T}_e + \hat{V} + \hat{V}_e$$

$$\langle \Psi | \hat{V}_e | \Psi \rangle = \langle \Psi | \int d^3r \int d^3r' \sum_{\sigma} \psi_{\sigma}^{\dagger}(\mathbf{r}) \psi_{\sigma}(\mathbf{r}) \frac{e^2}{4\pi\epsilon_0 |\mathbf{r} - \mathbf{r}'|} n_{\lambda}(\mathbf{r}') | \Psi \rangle$$

$$= \int d^3r \underbrace{n(\mathbf{r})}_{\text{density}} \underbrace{V_{\lambda}(\mathbf{r})}_{\text{potential}}$$

$$|\Psi\rangle = \sum_i a_i |i\rangle$$

$$\begin{aligned} \langle \Psi | H | \Psi \rangle &= \sum_{ij} a_j^{\dagger} a_i E_i \langle j | i \rangle \\ &= \sum_i |a_i|^2 E_i \\ &\geq \left(\sum_i |a_i|^2 \right) E_0 \end{aligned}$$

$$\frac{e^2}{4\pi\epsilon_0 |\mathbf{r} - \mathbf{r}'|} n_{\lambda}(\mathbf{r}') | \Psi \rangle$$

$$n(\mathbf{r}) = \langle \Psi | \sum_{\sigma} \psi_{\sigma}^{\dagger}(\mathbf{r}) \psi_{\sigma}(\mathbf{r}) | \Psi \rangle$$

$$E = \min_{n(\vec{r})} \left[\min_{\Psi \rightarrow n(\vec{r})} \left[\langle \Psi | \hat{T}_e + \hat{V} | \Psi \rangle + \int d^3r n(\vec{r}) V_e(\vec{r}) \right] \right]$$

$$E = \min_{n(\vec{r})} \left[F[n] + \int d^3r n(\vec{r}) V_e(\vec{r}) \right]$$

$$F[n] = \min_{\Psi \rightarrow n(\vec{r})} \langle \Psi | \hat{T}_e + \hat{V} | \Psi \rangle$$

$$\langle \Psi | \hat{T} | \Psi \rangle$$

$$\frac{\hbar^2}{2m}$$

$$\int \frac{d^3k}{(2\pi)^3} \frac{\hbar^2 k^2}{2m}$$

example

$$F[n]$$

$$= \int d^3r A(\vec{r}) n(\vec{r})$$

$$+ \int d^3r d^3r' B(\vec{r}-\vec{r}') n(\vec{r}) n(\vec{r}') + \int d^3r C(\vec{r}) (\nabla n(\vec{r}))^2$$

44.2 Kohn-Sham

$$n(r) = \sum_{i=1}^{N/2} \sum_{\sigma} |\phi_{i,\sigma}^{KS}(r)|^2$$

Approximation pour $\langle \Psi | \hat{T}_e | \Psi \rangle$

$$= \sum_{i=1}^{N/2} \sum_{\sigma} \int d^3r \phi_{i,\sigma}^{KS} \left(\frac{-\nabla^2}{2m} \right) \phi_{i,\sigma}^{KS}$$

$$\langle \Psi | \sum_{\sigma} \psi_{\sigma}^{\dagger}(r) \psi_{\sigma}(r) | \Psi \rangle$$

$$|\Psi\rangle = a_1^{\dagger} a_2^{\dagger} \dots |0\rangle$$

↑ ↑

$$\vec{B} = 0$$

ϕ réels

$$F[n] = \min_{KS} \left[\langle \Psi | \hat{T}_e | \Psi \rangle \right]_{KS} + V_{Hartree}[n] + E_{xc}[n]$$

$$E_{xc} = \min_{KS} \left(-\frac{1}{2} \int d^3r d^3r' \sum_{\sigma, \sigma'} \phi_{i,\sigma}^{KS}(r) \phi_{j,\sigma'}^{KS}(r) \frac{e^2}{4\pi\epsilon_0 |r-r'|} \phi_{j,\sigma'}^{KS}(r') \phi_{i,\sigma}^{KS}(r') \right)$$

$$+ \int d^3r C^x n^{4/3}(r)$$