

### III - Intro. propagateur à une particule $G$

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✓ 16. Info. dans  $G$

✓ 1. Représentation en opérateur

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18. Théorie des perturbations pour  $G$ 
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## 15. Définition

$$\Psi(\vec{r}, t) \theta(t-t') = i \int d^3r' G^R(\vec{r}, t; \vec{r}', t') \Psi_0(\vec{r}', t')$$

$$G^R(\vec{r}, t; \vec{r}', t') = -i \langle r | e^{-iH(t-t')} | r' \rangle \theta(t-t')$$

16.1 Info dans  $G$ :

$$G^R(\vec{r}, \vec{r}'; \omega) = \sum_n \frac{\Psi_n(\vec{r}) \Psi_n^*(\vec{r}')}{\omega + i\eta - E_n} = \langle r | \hat{G}^R(\omega) | r' \rangle$$

$$\hat{G}^R(\omega) = \frac{1}{\omega + i\eta - H}$$

$$\hat{G}^R(t) = -i e^{-iHt} \theta(t)$$

$$\hat{G}^A(\omega) = \frac{1}{\omega - i\eta - H}$$

$$\hat{G}^A(t) = i e^{-iHt} \theta(-t)$$

Pas d'interaction, invariant sous transl.

$$\langle k | \hat{G}^R(\omega) | k \rangle = \frac{\langle k | k \rangle}{\omega + i\eta - E_k}$$

$$H | k \rangle = E_k | k \rangle$$

16.2. Relation à DOS

$$\rho(E) = -\frac{1}{\pi} \text{Tr} \left[ \text{Im} \hat{G}^R(\omega) \right]$$

### 16.3 / Repr. spectrale et K.K.

$$G^R(r, r'; \omega) = \sum_n \frac{\psi_n(r) \psi_n^*(r')}{\omega + i\eta - E_n} = \int \frac{d\omega'}{2\pi} \frac{\sum_n \psi_n(r) \psi_n^*(r') 2\pi \delta(\omega' - E_n)}{\omega + i\eta - \omega'}$$

poiss spectrale

$$\sigma = \frac{1}{i(\omega + i\eta)} \left[ \chi_{jj}^R(\omega) - \frac{n \epsilon^2}{m} \right]$$

$$\chi_{jj}^R(\omega) = \int \frac{d\omega'}{\pi} \frac{\chi_{jj}''(\omega')}{\omega' - \omega - i\eta}$$

$$G^R(r, r'; \omega) = \int \frac{d\omega'}{2\pi} \frac{A(r, r'; \omega')}{\omega + i\eta - \omega'} \Rightarrow A(r, r'; \omega) = -2 \text{Im} G^R(r, r'; \omega)$$

Repr Spectrale  $\frac{1}{\omega + i\eta - \omega'} = \mathcal{P} \left( \frac{1}{\omega - \omega'} \right) - i\pi \delta(\omega - \omega')$  Poiss spectral

Relations K.K. pour G

$$\text{Re} [G^R(r, r'; \omega)] = \mathcal{P} \int \frac{d\omega'}{\pi} \frac{\text{Im} G^R(r, r'; \omega')}{\omega' - \omega}$$

$$\text{Im} [G^R(r, r'; \omega)] = - \mathcal{P} \int \frac{d\omega'}{\pi} \frac{\text{Re} G^R(r, r'; \omega')}{\omega' - \omega}$$

### 16.3.2 Règles de somme

$$A(r, r'; \omega) = \sum_n \Psi_n(r) \Psi_n^*(r') \delta(\omega - E_n) 2\pi$$

$E_n = \langle n | H | n \rangle$

$$\int \frac{d\omega}{2\pi} A(r, r'; \omega) = \sum_n \Psi_n(r) \Psi_n^*(r') = \sum_n \langle r | n \rangle \langle n | r' \rangle = \langle r | r' \rangle = \delta^3(r - r')$$

$$\int d^3r \int \frac{d\omega}{2\pi} \omega^n A(r, r; \omega) = \int d^3r \sum_m E_m^n \Psi_m(r) \Psi_m^*(r) = \int d^3r \langle r | H^n | r \rangle = \text{Tr}[H^n]$$

### 16.3.3 Dév. haute fréquence

$\langle \omega^n \rangle$

$$G^R(r, r'; \omega) = \int \frac{d\omega'}{2\pi} \frac{A(r, r'; \omega')}{\omega + i\eta - \omega'} = \sum_{n=0}^{\infty} \int \frac{d\omega'}{2\pi} \frac{\omega'^n A(r, r'; \omega')}{\omega^{n+1}}$$

$$\text{Tr}[\hat{G}^R(\omega)] = \int \frac{d\omega'}{2\pi} \frac{\text{Tr}[-2\text{Im}\hat{G}^R]}{\omega + i\eta - \omega'} = \sum_{n=0}^{\infty} \int \frac{d\omega'}{2\pi} \frac{\omega'^n \text{Tr}[2\pi\delta(\omega' - H)]}{\omega^{n+1}}$$

$$\hat{G}^R(\omega) = \frac{1}{\omega + i\eta - H} \quad \text{Im}\hat{G}^R = -\pi\delta(\omega - H) \quad \left. \vphantom{\hat{G}^R(\omega)} \right\} \sum_{n=0}^{\infty} \frac{\text{Tr}[H^n]}{\omega^{n+1}}$$

$$\frac{1}{\omega - \omega'} = \frac{1}{\omega(1 - \frac{\omega'}{\omega})} = \frac{1}{\omega} \left( 1 + \frac{\omega'}{\omega} + \left(\frac{\omega'}{\omega}\right)^2 + \dots \right)$$

## 16.9 Relation aux fluctuations

$$\begin{aligned} G^R &= -i e^{-iHt} \theta(t) \\ G^A &= i e^{-iHt} \theta(-t) \end{aligned}$$

$$\begin{aligned} S_{P_h P_{-h}}(\epsilon) &= \frac{1}{V} \langle P_h(t) P_{-h}(0) \rangle \\ &= \frac{1}{V} \langle e^{+iHt} P_h e^{-iHt} P_{-h} \rangle \\ &= \frac{1}{V} \langle i (G^R(-t) - G^A(-t)) P_h (G^R(t) - G^A(t)) i \rangle \\ & \quad \underbrace{\hspace{10em}}_{P_{-h}} \end{aligned}$$

$$G^R(-t) G^R(t) = 0$$

$$\begin{aligned} &= \frac{1}{V} \langle G^R(-t) P_h G^A(t) P_{-h} \rangle \\ & \quad + \frac{1}{V} \langle G^A(-t) P_h G^R(t) P_{-h} \rangle \end{aligned}$$

## 16.5 Relation aux E.D.

$$\Psi(r,t)\theta(t-t') = i \int d^3r' G^R(r,t; r',t') \Psi_0(r',t')$$

$$i \frac{\partial}{\partial t} [\Psi(r,t)\theta(t-t')] = H(r) \Psi(r,t)\theta(t-t') + i \Psi(r,t) \delta(t-t')$$

$$(i \frac{\partial}{\partial t} - H) [\Psi(r,t)\theta(t-t')] = i \Psi(r,t) \delta(t-t')$$

$$(i \frac{\partial}{\partial t} - H) i \int d^3r' G^R(r,t; r',t') \Psi_0(r',t') = \cancel{i} \Psi_0(r,t') \delta(t-t')$$

$$\Psi_0(r',t') = \delta^3(r'-r'')$$

$$(i \frac{\partial}{\partial t} - H) G^R(r,t; r'',t') = \delta^3(r-r'') \delta(t-t')$$



Opérateurs.

$$\hat{G}^R(t) = -i e^{iHt} \theta(t)$$

$$i \frac{\partial}{\partial t} \hat{G}^R(t) = H \hat{G}^R(t) + \delta(t)$$

$$\left( i \frac{\partial}{\partial t} - H \right) \hat{G}^R(t) = \delta(t) \quad \leftarrow \langle r | \quad | r' \rangle$$

$$\hat{G}^R(t) = \frac{1}{i \frac{\partial}{\partial t} - H} \delta(t)$$

$$i \frac{\partial}{\partial t} \psi = H \psi \\ = \left( -\frac{\hbar^2}{2m} \nabla^2 + V \right) \psi$$

$$i \frac{\partial}{\partial t} \langle r | \psi \rangle \\ = \int d^3r' \langle r | H | r' \rangle \langle r' | \psi \rangle$$

$$\int \frac{d\omega'}{\pi} \frac{1}{\omega - i\eta - \omega'} G^R(r, r'; \omega) = G^R(r, r'; \omega)$$

$$\frac{1}{\omega - i\eta - \omega'} = \mathcal{P} \frac{1}{\omega - \omega'} + i\pi \delta(\omega - \omega')$$

# 17. Phénoménologie pour $\Sigma^R$ self-énergie

Cas sans int.

$$\langle h | \hat{G}_0^D | h' \rangle = \frac{\langle h | h' \rangle}{\omega + i\eta - \epsilon_h} = G_0^R(h, \omega) \langle h | h' \rangle$$

$$G_0^R(h, \omega) = \frac{1}{\omega + i\eta - \epsilon_h}$$

$$A_0(h, \omega) = -2 \text{Im} G^R(h, \omega) = 2\pi \delta(\omega - \epsilon_h) \checkmark$$

$$G_0^R(h, \omega) = \int \frac{d\omega'}{2\pi} \frac{A_0(h, \omega')}{\omega + i\eta - \omega'}$$

Cas non inv. sous translation

$$G^R(h, \omega) = \int \frac{d\omega'}{2\pi} \frac{A(h, \omega')}{\omega + i\eta - \omega'}$$

Soit

$$A(h, \omega') = \frac{2\Gamma}{(\omega' - \tilde{\epsilon}_h)^2 + \Gamma^2}$$

$$A(h, \omega) = -2 \text{Im} G^R(h, \omega)$$

$$G^R(h, \omega) = \frac{1}{\omega + i\eta - \tilde{\epsilon}_h + i\Gamma}$$

$$\text{Im} G^R(h, \omega) = \text{Im} \left[ \frac{\omega - \tilde{\epsilon}_h - i\Gamma}{(\omega - \tilde{\epsilon}_h + i\Gamma)(\omega - \tilde{\epsilon}_h - i\Gamma)} \right]$$

$$G^R(h, \omega) = \frac{1}{\omega - \epsilon_h - \Sigma^R(h, \omega)}$$

Définition de  $\Sigma^R$

$$\Sigma^R(h, \omega) = \tilde{\epsilon}_h - \epsilon_h - i\Gamma$$

Cas particulier

$$G^R(h, \omega) = \frac{1}{G_0^{R-1}(h, \omega) - \Sigma^R(h, \omega)}$$

Général

$$G^{R-1}(h, \omega) = G_0^{R-1}(h, \omega) - \Sigma^R(h, \omega)$$

# 18. Théorie des perturbations

$$H = H_0 + V(\cdot) \quad H_0 |h\rangle = \frac{\hbar^2}{2m} |h\rangle$$

$$\hat{G}^R = \frac{1}{\omega + i\eta - H} = \frac{1}{\omega + i\eta - H_0 - V} \rightarrow (\omega + i\eta - H_0 - V) \hat{G}^R = 1$$

$$(\omega + i\eta - H_0) \hat{G}^R = 1 + V \hat{G}^R$$

$$\hat{G}_0^R \hat{G}^R = 1 + V \hat{G}^R$$

$$\hat{G}^R = \hat{G}_0^R + \hat{G}_0^R V \hat{G}^R = \hat{G}_0^R + \hat{G}_0^R V \hat{G}_0^R + \dots$$

$$= \hat{G}_0^R + \hat{G}_0^R V \hat{G}_0^R + \hat{G}_0^R V \hat{G}_0^R V \hat{G}_0^R$$

$$+ \hat{G}_0^R V \hat{G}_0^R V \hat{G}_0^R V \hat{G}_0^R + \dots$$

$$\langle h | \hat{G}^R | h' \rangle = \langle h | \hat{G}_0^R | h' \rangle + \int \frac{d^3 h_1}{(2\pi)^3} \int \frac{d^3 h_2}{(2\pi)^3} \langle h | \hat{G}_0^R | h_1 \rangle$$

$$\langle h_1 | V | h_2 \rangle \langle h_2 | \hat{G}_0^R | h' \rangle + \dots$$

$$= G_0^R(h, \omega) \langle h | h' \rangle + G_0^R(h, \omega) \langle h | V | h' \rangle G_0^R(h', \omega)$$

$$\int \frac{d^3 h_1}{(2\pi)^3} G_0^R(h_2, \omega) \langle h | V | h_1 \rangle G_0^R(h_1, \omega) \langle h_1 | V | h' \rangle G_0^R(h', \omega)$$

$$\Rightarrow = \rightarrow + \rightarrow \overset{x}{\vdots} \Rightarrow = \rightarrow + \rightarrow \overset{x}{\vdots} \rightarrow + \dots$$

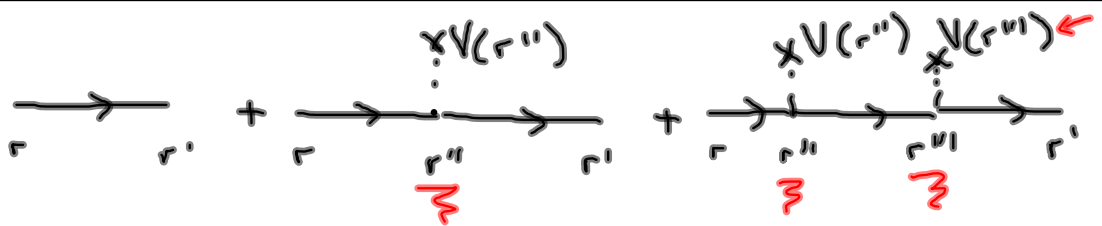
$$= \rightarrow_k + \rightarrow_k \overset{x}{\uparrow} \overset{h-h'}{h'} + \rightarrow_k \overset{x}{\uparrow} \overset{h-h_1}{h_1} + \dots$$

①  $\rightarrow$  devient  $G_0^R(h, \omega)$

②  $\overset{x}{\uparrow} h-h'$  devient  $\langle h | V | h' \rangle$

③  $\int \frac{d^3 h_i}{(2\pi)^3}$  de tous les  $h_i$  nous déterminés par conservation de  $h$

④  $h$  est conservé à chaque vertex



### Série de Born

Si on se tance vite.

$$\langle h | G^R | h' \rangle = \langle h | G_0^R | h, \omega \rangle + G_0^R | h, \omega \rangle V | h - h' \rangle G_0^R | h', \omega \rangle$$

$$-2 \text{Im} G = \sum_n \Psi_n(k) \Psi_n^*(k') 2\pi \delta(\omega - E_n)$$

$$\text{Im} \left[ \left( \frac{1}{\omega + i\eta - E_h} \right)^2 \right] = \frac{\partial}{\partial \omega} \text{Im} \left[ \frac{1}{\omega + i\eta - E_h} \right]$$

dérivée