

$$\cdot \chi_{A_j A_i}^{\mathcal{R}}(\omega) = \int \frac{d\omega'}{\pi} \frac{\chi_{A_j A_i}''(\omega')}{\omega' - \omega - i\eta}$$

$z \rightarrow \omega + i\eta$ Retardée

$z \rightarrow \omega - i\eta$ Avancée.

$\omega \chi_{A_j A_i}''(\omega)$ positive def.

$$\cdot S_{A_i A_j}(\omega) = \frac{\hbar}{1 - e^{-\hbar\omega}} \chi_{A_i A_j}''(\omega)$$


$$S_{A_j A_i}(-t) = S_{A_i A_j}(t - i t \hbar \beta)$$
$$S_{A_j A_i}(-\omega) = e^{-\beta \hbar \omega} S_{A_i A_j}(\omega)$$

Menu du jour!

- Règles de somme
- Kubo pour conductivité

Règles de somme.

- Relier exp. les unes aux autres $\text{Re}\sigma$

$$\int \text{Re}\sigma d\omega = \frac{\omega_p^2}{8\pi} = \frac{ne^2}{2m} \omega$$


- Limites $\omega \rightarrow \infty$ fcts de rrp.
- Contraintes sur phénoménologie

Règle de somme T.D.

$$\chi_{A_i A_j} \equiv \frac{\partial \langle A_i \rangle}{\partial a_j} = \lim_{g \rightarrow 0} \chi_{A_i A_j}^{\mathcal{R}}(g, 0)$$

$$= \lim_{g \rightarrow 0} \int_{-\infty}^{\infty} \frac{d\omega}{\pi} \frac{\chi_{A_i A_j}''(\omega)}{\omega}$$

$$\delta \langle A_i \rangle_{\omega} = \chi_{A_i A_j}^{\mathcal{R}}(\omega) \delta a_j(\omega)$$

$$\lim_{g \rightarrow 0} \chi_{nn}^{\mathcal{R}}(g, \omega=0)$$

$$= \lim_{g \rightarrow 0} \int_{-\infty}^{\infty} \frac{d\omega}{\pi} \frac{\chi_{nn}''(g, \omega)}{\omega} = \left(\frac{\partial n}{\partial \mu} \right)_{T,V}$$

Autre preuve

$$\langle NN \rangle - \langle N \rangle^2 = \frac{1}{\beta} \left(\frac{\partial n}{\partial \mu} \right)_{T, V}$$

$$= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} S_{NN}(\omega)$$

F.D.

$$S_{NN}(\omega) = \lim_{\omega \rightarrow 0} \int dt e^{i\omega t} \langle N(t)N \rangle$$

$$\propto \frac{2t}{1 - e^{-\beta t \omega}} \chi''_{NN}(\omega)$$

$$= \frac{2}{\beta \omega} \chi''_{NN}(\omega)$$

Moments, règles de somme
et développement $\omega \rightarrow \infty$

$$\chi_{A_i A_j}^R(q, \omega) = \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} \frac{\chi_{A_i A_j}''(\omega')}{\omega' - \omega - i\eta}$$

$$\stackrel{\omega \rightarrow \infty}{\sim} \sum_{n=1}^{\infty} -\frac{1}{\omega^{2n}} \int_{-\infty}^{\infty} \frac{d\omega'}{\pi} (\omega')^{2n-1} \chi_{A_i A_j}''(q, \omega')$$

Calcul des moments.

$$\begin{aligned} & \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \omega^n \chi_{A_i; A_j}''(\omega) \\ &= \left(i \frac{\partial}{\partial t} \right)^n \left[\int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{-i\omega t} 2\chi_{A_i; A_j}''(\omega) \right] \\ &= \frac{1}{\hbar} \left\langle \left[\left(i \frac{\partial}{\partial t} \right)^n A_i(t), A_j(0) \right] \right\rangle \quad \begin{array}{l} t=0 \\ \swarrow \end{array} \end{aligned}$$

$$= \frac{1}{\hbar} \langle [[A_i(t), \frac{H}{\hbar}], \frac{H}{\hbar}] \dots A_j(t) \rangle_{t=0}$$

$$i \frac{\partial}{\partial t} A_i(t) = [A_i(t), \frac{H}{\hbar}]$$

$$A_i(t) = e^{iHt/\hbar} A_i e^{-iHt/\hbar}$$

Exemple : règle de somme f

$$\int_{-\infty}^{\infty} \frac{d\omega}{\pi} \omega \chi''_{nn}(q, \omega) = \frac{nq^2}{m}$$

à prouver ↗

$$= \frac{i}{\hbar V} \left\langle \left[\frac{\partial}{\partial t} n_q(t), n_{-q}(0) \right] \right\rangle_{t=0}$$

$$= -\frac{i}{\hbar^2 V} \left\langle \left[\underbrace{(H, n_q)}, n_{-q} \right] \right\rangle_{(1)}$$

V ?

$$\int d^3(r-r') \langle n(r)n(r') \rangle e^{-iq \cdot (r-r')}$$
$$= \frac{1}{V} \int d^3 r e^{-iq \cdot r} \int d^3 r' e^{iq \cdot r'} \langle n(r)n(r') \rangle$$

$\frac{r+r'}{2}, r-r' \rightarrow r, r'$

$$\begin{aligned} \rightarrow [p_\beta^x, n_q] &= \frac{\hbar}{i} \left[\frac{\partial}{\partial x_\beta}, \sum_\alpha e^{-iqr_\alpha} \right] \\ n_q &= \int dr e^{-iqr} \sum_\alpha \delta(r - r_\alpha) \\ &= \sum_\alpha e^{-iqr_\alpha} \underbrace{\hspace{10em}}_{n(r)} \\ \rightarrow &= -\hbar q^\beta e^{-iqr_\beta} \end{aligned}$$

$$\left[\sum_{\beta} \frac{p_{\beta}^2}{2m}, n_{\beta} \right] = \frac{2}{2m} \sum_{\beta} \vec{p}_{\beta} \cdot (-\hbar \vec{q} e^{-i\vec{q} \cdot \vec{r}_{\beta}})$$

$$[p \cdot p, n] = p \cdot [p, n] + [p, n] \cdot p$$

$$\begin{aligned} & \frac{1}{3} \sum_{\beta} (-\hbar \vec{q} e^{-i\vec{q} \cdot \vec{r}_{\beta}} \cdot [p_{\beta}, n_{-\beta}]) \\ &= \frac{1}{3} \sum_{\beta} \hbar^2 q^2 e^{-i\vec{q} \cdot \vec{r}_{\beta}} e^{i\vec{q} \cdot \vec{r}_{\beta}} = -\frac{\hbar^2 q^2}{m} N \end{aligned}$$

(2)

Formule de Kubo

$$E_y(q_x, \omega) = i(\omega + i\eta) A_y / c$$

diamagnétique.

À prouver

$$\delta \langle j_m(q, \omega) \rangle = \left[\chi_{j_m j_\omega}^R(q, \omega) - \frac{ne^2}{m\omega} \delta \right] \frac{A_\omega(q, \omega)}{c} - \chi_{j_m \rho}^R(q, \omega) \phi(q, \omega)$$

$$E = -\frac{1}{c} \frac{\partial A}{\partial t} - \nabla \varphi$$

$$B = \nabla \times A$$

$$A(t) = \int \frac{d\omega}{2\pi} e^{-i(\omega + i\eta)t} A(\omega)$$

Inv. de jauge

$$A \rightarrow A + \nabla \Lambda$$

$$\varphi \rightarrow \varphi - \frac{1}{c} \frac{\partial \Lambda}{\partial t}$$

$$\psi \rightarrow e^{ie\Lambda/\hbar c} \psi$$

$$\vec{p}_d = \frac{\hbar}{i} \vec{\nabla}_d \longrightarrow \left(\frac{\hbar}{i} \vec{\nabla}_d - \frac{e}{c} \vec{A}(r_d) \right)$$

Inv. de jauge.

$$\delta H(t) = \int d^3r \varphi(r,t) \rho(r)$$

$$\begin{aligned} -\frac{\hbar^2}{2m} \nabla_{\alpha}^2 \rightarrow & -\frac{\hbar^2}{2m} \nabla_{\alpha}^2 - \frac{e\hbar}{2mc} (A(r_{\alpha}) \cdot \nabla_{\alpha} \\ & + \nabla_{\alpha} \cdot A(r_{\alpha})) + \frac{e^2}{2mc^2} A^2(r_{\alpha}) \end{aligned}$$

$$\delta H(t) = - \sum_{\alpha} \frac{e\hbar}{2mc} (A(r_{\alpha}) \cdot \nabla_{\alpha} + \nabla_{\alpha} \cdot A_{\alpha}(r_{\alpha}))$$

$$= -\frac{1}{c} \int d^3r A(r,t) j(\vec{r}) \quad \leftarrow$$

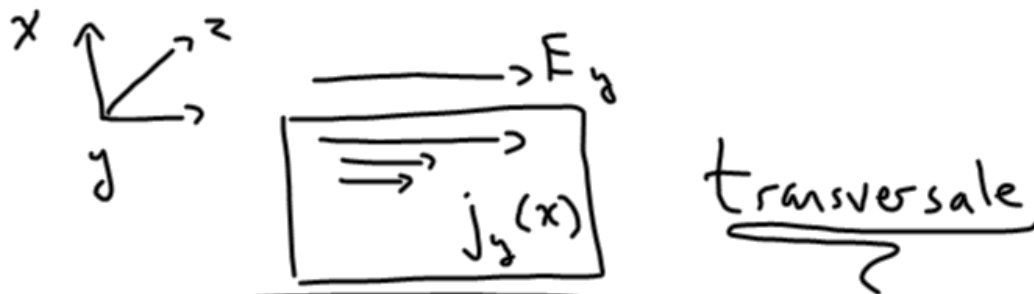
$$j(r) = \frac{e}{2m} \sum_{\alpha} (\delta(r-r_{\alpha}) p_{\alpha} + p_{\alpha} \delta(r-r_{\alpha}))$$

$$\vec{j}^A(\vec{r}) = \vec{j}(\vec{r}) - \frac{e}{mc} A(\vec{r}) \rho(\vec{r})$$

$$\hat{p}_\alpha \rightarrow \frac{\hbar}{i} \nabla_\alpha - \frac{e}{c} A(\vec{r}_\alpha)$$

$$\rho(\vec{r}) = \sum_\alpha \delta(\vec{r} - \vec{r}_\alpha)$$

Conductivité transversale



$$\delta \langle j_y(\rho_x, \omega) \rangle = \sigma_{yy}(\rho_x, \omega) E_y(\rho_x, \omega)$$

def.

$$\begin{array}{ll}
 \text{Transversal} & \vec{\nabla} \cdot \vec{C} = 0 \quad \vec{q} \cdot \vec{C} = 0 \\
 \text{Longitudinal} & \vec{\nabla} \times \vec{C} = 0 \quad \vec{q} \times \vec{C} = 0
 \end{array}$$

$$\vec{A}^L = (\hat{q} \cdot \vec{A}) \hat{q}$$

$$\vec{A}^T = (\overset{\leftrightarrow}{\mathbf{I}} - \hat{q} \hat{q}) \cdot \vec{A}$$

$$\begin{aligned} \overleftrightarrow{\sigma}^{\uparrow}(\mathbf{q}, \omega) &= (\overleftrightarrow{\mathbb{I}} - \hat{\mathbf{q}}\hat{\mathbf{q}}) \cdot \overleftrightarrow{\sigma} (\overleftrightarrow{\mathbb{I}} - \hat{\mathbf{q}}\hat{\mathbf{q}}) \\ \overleftrightarrow{\sigma}^{\downarrow}(\mathbf{q}, \omega) &= \hat{\mathbf{q}}\hat{\mathbf{q}} \cdot \overleftrightarrow{\sigma} \cdot \hat{\mathbf{q}}\hat{\mathbf{q}} \end{aligned}$$

$$\sigma_{yy}(\mathbf{q}_x, \omega) = \frac{1}{i(\omega + i\eta)} \left[\chi_{j_y j_y}^{\mathcal{R}}(\mathbf{q}_x, \omega) - \frac{ne^2}{m} \right]$$

Conductivitate longitudinală dinale

À prouver :

$$\begin{aligned} \delta \langle j_x(q_x, \omega) \rangle &= \frac{1}{i(\omega + i\eta)} \left[\chi_{j_x j_x}^R(q_x, \omega) - \frac{v_1 e^2}{m} \right] \\ &= \frac{1}{i q_x} \chi_{j_x p}^R(q_x, \omega) \left(\frac{i(\omega + i\eta)}{c} A_x(q_x, \omega) - q_x \phi(q_x, \omega) \right) \end{aligned}$$

$$\frac{\partial \rho}{\partial t} = -\nabla \cdot \mathbf{j}$$

$$\boxed{\frac{\partial \rho(\mathbf{q}, t)}{\partial t} = -i\mathbf{q} \cdot \mathbf{j}(\mathbf{q}, t)}$$

$$\frac{\partial}{\partial t} \chi_{\mathbf{j}_x \rho}^{\mathcal{R}}(\mathbf{q}_x, t) = \frac{\partial}{\partial t} \left(\frac{i}{\hbar} \Theta(t) \langle [j_x(\mathbf{q}_x, t), \rho(-\mathbf{q}_x, 0)] \rangle \right)$$

$$\delta(t) \frac{i}{\hbar V} \langle [j_x(\mathbf{q}_x, 0), \rho(-\mathbf{q}_x, 0)] \rangle$$

$$+ \Theta(t) \frac{i}{\hbar V} (-i q_x) \langle [j_x(\mathbf{q}_x, 0), j_x(-\mathbf{q}_x, -t)] \rangle$$

$$\boxed{-i(\omega + i\eta) \chi_{\mathbf{j}_x \rho}^{\mathcal{R}}(\mathbf{q}, \omega) = i q_x \frac{n e^2}{m} - i q_x \chi_{\mathbf{j}_x \mathbf{j}_x}^{\mathcal{R}}(\mathbf{q}, \omega)}$$

$$\frac{c_{ar}}{tV} \frac{i}{\hbar} \langle [j_x(q_x, 0), \rho(-q_x, 0)] \rangle = \frac{i q_x n e^2}{m}$$

$$= i \int \frac{d\omega}{\pi} \chi''_{j_x \rho}(q_x, \omega) = i \int \frac{d\omega}{\pi} \frac{\omega}{q_x} \chi''_{\rho \rho}(q_x, \omega)$$

$$\frac{\partial \rho}{\partial t} = -i q_x \cdot j_x = \frac{i}{q_x} \frac{n e^2}{m} q_x^2$$

$$-i\omega \rho(q, \omega) = -i q_x \cdot j_x(q_x, \omega)$$

$$\sigma_{xx}(q_x, \omega) = \frac{1}{i(\omega + i\gamma)} \left[\sum_{j_x}^R \chi_{j_x} \cdot (q_x, \omega) - \frac{ne^2}{m} \right]$$

$$= \frac{1}{iq_x} \chi_{j_x} \rho(q_x, \omega)$$

$\varphi=0$ \wedge indep. det \Rightarrow

$$\rightarrow \left[\sum_{j_x} \chi_{j_x} (q_x, 0) - \frac{ne^2}{m} \right] = 0$$

Règle de somme f

$$A = 0 \quad \delta \langle j_x(q, \omega) \rangle = -\chi_{j \times p}^R(q, \omega) \phi(q, \omega)$$

Λ indép. de x , dép. de t .

$$A \rightarrow A + \nabla \Lambda = 0$$

$$\varphi \rightarrow \varphi - \frac{1}{c} \frac{\partial \Lambda}{\partial t} \neq 0$$

$$\Rightarrow \boxed{\chi_{j \times p}^R(0, \omega) = 0}$$

	D	D _s
Métal	D	0 [↙]
Isolant	0	0
S.c.	D	D _s [↙]

$$D_s = \pi \lim_{q_x \rightarrow 0} \left[\frac{ne^2}{m} - \chi^R_{ij}(q_x, 0) \right]$$

$$D = \pi \lim_{\omega \rightarrow 0} \left[\frac{ne^2}{m} - \text{Re} \chi_{jxjx}^R(0, \omega) \right]$$



si pas
d'interaction

