

4.6 Signif. physique de $A(k, \omega)$

1. Sans int.

2. Repr. de Lehman

3. Interp. probabiliste

4. ARPES

5. Quasiparticules.

6. Liquide de Fermi et ARPES.

7. $n_k = \langle C_k^\dagger C_k \rangle$ dans liquide de Fermi.

Comportement asymptotique.

4.7 3 théorèmes.

1. Wick 2. Graphs connexes 3. Principe var.

$$A(k, \omega) = \delta(\omega - \epsilon_k) 2\pi$$

$$A(k, \omega) = e^{\beta \Omega} \sum_{m, n} (e^{-\beta K_m} + e^{-\beta K_n})$$



$$|\langle n | c_k | m \rangle|^2 2\pi \delta(\omega - (\epsilon_m - \epsilon_n))$$

$$|m\rangle = c_{k_1}^+ c_{k_2}^+ \dots c_{k_n}^+ |0\rangle$$

$$\langle n| = \langle 0| c_{k_1} \dots c_{k_n}$$

$$|m\rangle = a_1 c_{k_1}^+ \dots c_{k_n}^+ |0\rangle + a_2 c_{k_1}^+ \dots c_{k_n}^+ |1\rangle + \dots$$

$$\frac{d^2 \sigma}{d\Omega d\omega} \propto f(\omega) A(k, \omega)$$

4.6.5 Quasiparticules.

$$G^R(k, \omega) = \frac{1}{\omega + i\eta - \epsilon_k - \Sigma^R(k, \omega)}$$

$$\frac{a + b' - ib''}{a + b - ib''} \frac{1}{a + b' + ib''}$$

$$\begin{aligned} \text{Im} &= -b'' \\ &= \frac{-b''}{(a+b')^2 + (b'')^2} \end{aligned}$$

$$A(k, \omega) = -2 \text{Im} G^R(k, \omega)$$

$$G^R(k, \omega) = \int \frac{d\omega'}{2\pi} \frac{A(k, \omega')}{\omega + i\eta - \omega'}$$

$$A(k, \omega) = \frac{-2 \text{Im} \Sigma^R(k, \omega)}{(\omega - \epsilon_k - \text{Re} \Sigma^R(k, \omega))^2 + (\text{Im} \Sigma^R(k, \omega))^2}$$

$$(\epsilon_k - \mu) - \epsilon_k - \text{Re} \Sigma^R(k, \epsilon_k - \mu) = 0$$

$$\omega = \epsilon_k - \mu$$

Quasiparticule

$$\epsilon_k = \epsilon_k + \text{Re} \Sigma^R(k)$$

Hartree-

Développe autour de

$$\omega = \epsilon_k - \mu$$

Fock
 Σ^R indep. de ω

$$\omega - \xi_h - \text{Re} \Sigma^R(h, \omega) = 0 + \frac{\partial}{\partial \omega} [\omega - \xi_h - \text{Re} \Sigma^R] (\omega - (E_h - \mu))$$

$$= \left(1 - \frac{\partial \text{Re} \Sigma^R(h, \omega)}{\partial \omega} \right) (\omega - (E_h - \mu))$$

$$\equiv Z_h^{-1} (\omega - (E_h - \mu)) \quad \omega = E_h - \mu$$

$$A(h, \omega) = \frac{-2 \text{Im} \Sigma^R}{(Z_h^{-1})^2 (\omega - (E_h - \mu))^2 + (\text{Im} \Sigma^R)^2}$$

$$= 2\pi Z_h \frac{1}{\pi} \frac{-Z_h \text{Im} \Sigma^R}{(\omega - (E_h - \mu))^2 + (Z_h \text{Im} \Sigma^R)^2}$$

$$\Gamma_h \equiv -Z_h \text{Im} \Sigma^R \quad \Gamma_h > 0$$

$$A(h, \omega) = Z_h 2\pi \left[\frac{1}{\pi} \frac{\Gamma_h}{(\omega - (E_h - \mu))^2 + \Gamma_h^2} \right]$$

Renormalisation de Ψ + incohérent

Quasiparticule

$$\nabla_h [(E_h - \mu) - \xi_h - \text{Re} \Sigma^R(h, E_h - \mu)] = 0$$

$$N_h = \nabla_h (E_h - \mu) \quad \text{Bandes}$$

$$N_h^* = \nabla_h (E_h - \mu) = \text{vitesse renormalisée par interaction.}$$

$\xi_h \leftrightarrow h$

$$N_h^* - N_h - \nabla_h \text{Re} \Sigma^R(h, (E_h - \mu)) = 0$$

$$- N_h^* \frac{\partial \text{Re} \Sigma^R(h, \omega)}{\partial \omega} \Big|_{\omega = E_h - \mu} = 0$$

$$N_h^* \left(1 - \frac{\partial \text{Re} \Sigma^R(h, \omega)}{\partial \omega} \right) \Big|_{\omega = E_h - \mu} - N_h \left(1 + \frac{\partial \text{Re} \Sigma^R}{\partial \xi_h} \right) = 0$$

Z_h^{-1}

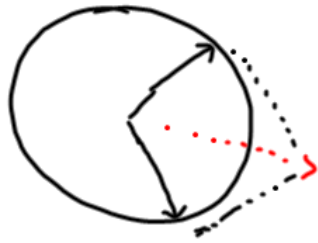
$$\nabla_h \text{Re} \Sigma^R(h, \omega)$$

$$= \nabla_h \xi_h \frac{\partial \text{Re} \Sigma^R(h, \omega)}{\partial \xi_h}$$

$$\frac{m}{m^*} = \frac{N_h^*}{N_h} = \left(1 + \frac{\partial \text{Re} \Sigma^R(h, E_h - \mu)}{\partial \xi_h} \right) Z_h$$

$$m N_h^* = \hbar k_F \quad k_F \text{ ne dépend pas des interactions}$$

4.6.6. Liquides de Fermi



$$\text{Im} \Sigma^R(\hbar, \omega) = 0 \text{ à } \omega = 0$$

Dépendance en ω

Analytique:

$$\text{Im} \Sigma^R(\hbar, \omega) = a\omega - \gamma\omega^2 + \dots$$

$a=0$ par causalité.

$$\rightarrow \boxed{\text{Im} \Sigma^R(\hbar, \omega) = -\gamma\omega^2}$$

$$G^R = G_0^R + G_0^R \Sigma^R G^R$$

$$\Sigma^R(\hbar, \omega) - \Sigma^R(\hbar, \infty) = \int \frac{d\omega'}{\pi} \frac{\text{Im} \Sigma^R(\hbar, \omega')}{\omega' - \omega - i\eta}$$

$$\text{Re} \Sigma^R(\hbar, \omega) - \Sigma^R(\hbar, \infty) = \int \frac{d\omega'}{\pi} \frac{\text{Im} \Sigma^R}{\omega'}$$

$$\rightarrow Z_h^{-1} = \left(1 - \frac{\partial}{\partial \omega} \text{Re} \Sigma^R\right) + \omega \int \frac{d\omega'}{\pi} \frac{\text{Im} \Sigma^R}{\omega'^2} + \dots$$

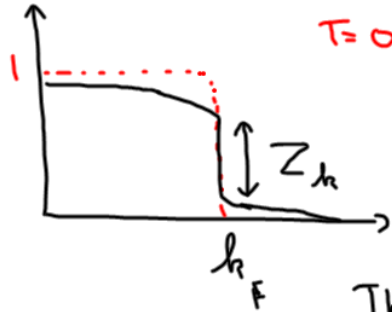
$$\text{Re} \Sigma^R = \text{cte} + \omega(-\gamma)$$

$$Z_h^{-1} = (1 + \gamma) > 1$$

$$\boxed{Z_h < 1}$$

4.6.7 Distribution de h

$$n_h = \langle c_h^+ c_h \rangle$$



Théorème de
Luttinger

$$\langle c_h^+ c_h \rangle = \lim_{\tau \rightarrow 0^-} G(h, \tau)$$

$$G(h, \tau) = - \langle T_\tau c_h(\tau) c_h^+ \rangle$$

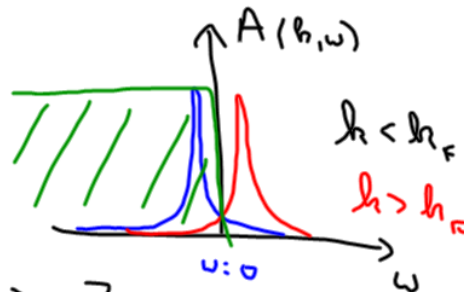
$$G(h, \tau) = T \sum_n e^{i h_n \tau} \int \frac{d\omega}{2\pi} \frac{A(h, \omega)}{i h_n - \omega}$$

$$= \int \frac{d\omega}{2\pi} \frac{A(h, \omega) e^{-\omega \tau}}{e^{\beta \omega} + 1}$$

$$\lim_{\tau \rightarrow 0^-} G(h, \tau) = \int \frac{d\omega}{2\pi} f(\omega) A(h, \omega) = n_h$$

$T=0$

$$n_h = \int_{-\infty}^0 \frac{d\omega}{2\pi} A(h, \omega)$$



$$\langle c_{h-}^+ c_{h-} \rangle - \langle c_{h+}^+ c_{h+} \rangle = Z_h$$

Comportement asymptotique.

$$G(h, ih_n) = \int \frac{d\omega}{2\pi} \frac{A(h, \omega)}{ih_n - \omega}$$

$$\lim_{ih_n \rightarrow \infty} G(h, ih_n) = \frac{\int \frac{d\omega}{2\pi} A(h, \omega)}{ih_n} = \frac{1}{ih_n}$$

$$G(h, ih_n) = \frac{1}{ih_n - \mathcal{P}_h - \Sigma(h, ih_n)}$$

$$\lim_{ih_n \rightarrow \infty} \Sigma(h, ih_n) = \text{cte.} \quad \text{Hartree-Fock}$$

4.7 3 théorèmes.

4.7.1 Théorème de Wick

$$\langle T_{\tau} \psi(\tau) \psi^{\dagger}(0) \rangle = T_{\tau} \left[e^{-\beta H_0} \hat{U}(\beta, 0) \hat{U}(0, \tau) \hat{\psi} \hat{U}(\tau, 0) \right]$$

$$\hat{U}(\tau, 0) = T_{\tau} e^{-\int_0^{\beta} \hat{V}(\tau) d\tau} T_{\tau} \left[e^{-\beta H_0} \hat{U}(\beta, 0) \right]$$

$$\langle T_{\tau} \psi(\tau) \psi^{\dagger} \int_0^{\beta} d\tau' \psi^{\dagger}(\tau') \psi(\tau') \frac{1}{|\tau - \tau'|} \rangle$$

Simple: 1 electron:

$$\langle a, a^\dagger \rangle = \frac{\text{Tr} e^{-\beta K_0} a, a^\dagger}{\text{Tr} e^{-\beta K_0}}$$

$$= \frac{\langle 0 | a a^\dagger | 0 \rangle + \langle 0 | a \rangle e^{-\beta \epsilon_1} a, a^\dagger \langle a^\dagger | 0 \rangle}{\langle 0 | 0 \rangle + \langle 0 | a \rangle e^{-\beta \epsilon_1} \langle a^\dagger | 0 \rangle}$$

$$= \frac{1}{1 + e^{-\beta \epsilon_1}}$$

2 electrons:

$$|0\rangle \quad a_1^\dagger |0\rangle \quad a_2^\dagger |0\rangle \quad a_1^\dagger a_2^\dagger |0\rangle$$

$$\rightarrow \langle 0 | (1+a_1)(1+a_2) \Theta (1+a_1^\dagger)(1+a_2^\dagger) | 0 \rangle$$

$$\langle 0 | (1+a_1) \Theta (1+a_1^\dagger) | 0 \rangle$$

$$\langle 0 | \Theta | 0 \rangle$$

$$\langle 0 | a_1 \Theta a_1^\dagger | 0 \rangle$$

$$\begin{aligned} \langle a_1 a_1^\dagger a_2 a_2^\dagger \rangle &= \frac{1}{1+e^{-\beta \epsilon_1}} \frac{1}{1+e^{-\beta \epsilon_2}} \frac{1+e^{-\beta \epsilon_3}}{1+e^{-\beta \epsilon_3}} \\ &= \langle a_1 a_1^\dagger \rangle \langle a_2 a_2^\dagger \rangle \end{aligned}$$

$$\langle a_i(\tau_i) a_j(\tau_j) a_h^+(\tau_h) a_l^+(\tau_l) \rangle$$

$$a_\alpha(\tau) = e^{-\tau \mathcal{L}_\alpha} a_\alpha \quad \text{si } H = \sum_\alpha \mathcal{L}_\alpha a_\alpha^+ a_\alpha$$

$$a_i(\tau) = \langle i | \alpha \rangle a_\alpha(\tau) \quad \text{avec } \sum_\alpha \text{ implicite.}$$

$$= e^{-\tau \mathcal{L}_\alpha} \langle i | \alpha \rangle a_\alpha$$

$$\langle i | \alpha \rangle \langle j | \beta \rangle \langle a_\alpha a_\beta a_h^+ a_l^+ \rangle \langle \beta | h \rangle \langle \alpha | l \rangle$$

$$e^{-\mathcal{L}_\alpha \tau_i - \mathcal{L}_\beta \tau_j + \mathcal{L}_h \tau_h + \mathcal{L}_l \tau_l}$$

$$\langle a_\alpha a_\beta a_h^+ a_l^+ \rangle = \langle a_\alpha a_\alpha^+ \rangle \langle a_\beta a_\beta^+ \rangle \delta_{\alpha\beta} \delta_{hl}$$

$$\begin{matrix} \beta = \gamma \\ \alpha = \delta \\ \alpha \neq \beta \end{matrix} \quad - \langle a_\alpha a_\alpha^+ \rangle \langle a_\beta a_\beta^+ \rangle \delta_{\alpha\beta} \delta_{hl}$$