

Résumé  $\uparrow R_i = R_{i0} + u_i$

$$H = \underbrace{K_e + V_{ee}} + \underbrace{K_i + V_{ii}} + \underbrace{V_{ie}}$$

$$V_{ie} = -\frac{4\pi e^1 z}{|r - R_i|} + \vec{u}_i \cdot \vec{\nabla}_{R_i} V_{ie}$$

$$Z = Z_{0e} Z_{0i} T_e \left[ e^{-\beta(K_e + V_{ee})} e^{-\beta(K_i + V_{ii})} \frac{1}{T_i} e^{-\int V_{ie} d\tau} \right]$$

$$= Z_{0e} Z_{0i} \left\langle T_i e^{\frac{1}{2} \int d\tau dz' \overline{Z_0^e Z_0^i} (V_{ie}(\tau) V_{ie}(z'))} \right\rangle_e$$

$$V_{eff} = V_c(q) \delta(\tau - \tau') + |M_{ij}|^2 \mathcal{D}(q, \tau - \tau')$$



$$\dots = (\dots + \dots) + (\dots + \dots) \circlearrowleft \dots$$

$$\therefore V_c + \frac{1}{1+V_c \chi} \left[ \frac{V_p}{1+V_p \chi} \frac{1}{1+V_c \chi} \right]$$

$$D^e = \frac{D_0(\omega, i\omega_n)}{1 + \frac{|M_g|^2}{1 + V_c X} X D_0(\omega, i\omega_n)}$$

$$D_0 = \frac{-2\omega_{ip}}{\omega_n^2 + \omega_{ip}^2}$$

$$= \frac{-g\omega_{ip}}{\omega_n^2 + \omega_{ip}^2 - \omega_{ip}^2 \frac{V_c X}{1 + V_c X}} = \frac{-g\omega_{ip}}{\omega_n^2 + \frac{\omega_{ip}^2}{1 + V_c X}}$$

$\omega^2(\omega)$

$$+ \dots$$

$$\boxed{|M_g|^2 (2\omega_{ip}) = \omega_{ip}^2 V_c}$$

$$\omega_g^2 = \frac{\omega_{ip}^2}{1 + V_c \chi} = \frac{\omega_{ip}^2}{\epsilon(q)} \propto C_s^2 q^2$$

$$V_{\text{eff}} = V_c^e + \text{---} \downarrow \text{---}$$

$$\frac{M_q}{\epsilon(q, i q_n)} \frac{-2\omega_{ip}}{q_n^2 + \omega_g^2} \frac{M_{-q}}{\epsilon(-q, i q_n)}$$

$$\frac{-\omega_{ip}^2 V_c(q)}{\epsilon(q_n^2 + \omega_g^2) \epsilon}$$

An total.

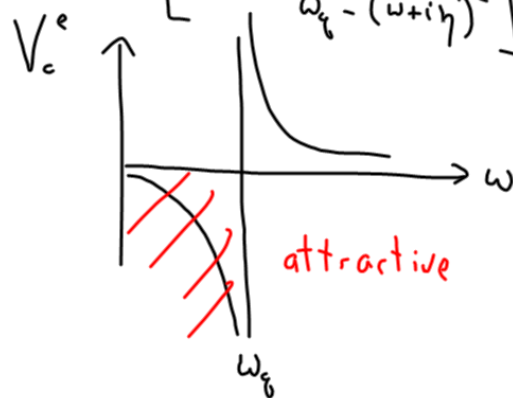
$$\frac{-\omega_g^2 V_c(q)}{(q_n^2 + \omega_g^2) \epsilon(q, i q_n)}$$

$$V_c^e + \frac{-\omega_g^2 V_c^e}{q_n^2 + \omega_g^2}$$

$$= V_c^e \left[ 1 - \frac{\omega_g^2}{\omega_g^2 - (i q_n)^2} \right]$$

Prolongement:

$$= V_c^e \left[ 1 - \frac{\omega_g^2}{\omega_g^2 - (\omega + i\eta)^2} \right]$$



$$\begin{aligned}
 \text{Re } V_{\text{eff}}^{\text{TOT}} &= \frac{V(q)}{\epsilon(q)} \left[ \frac{-\omega^2}{\omega_p^2 - \omega^2} \right] \\
 &= \frac{V(q)}{\epsilon(q) - \frac{\omega_p^2}{\omega^2} \epsilon(q)} \\
 &= \frac{V(q)}{\epsilon(q) - \frac{\omega_{ip}^2}{\omega^2}} \leftarrow \\
 &= \frac{V(q)}{1 + \underbrace{(\epsilon-1)_{cl}} + \underbrace{(\epsilon-1)_{ph}} }
 \end{aligned}$$

## 8) Supraconductivité.

1) Inst. de Cooper

1. Équation de vertex
2. Sol.

2) Phénoménologie.

3) BCS

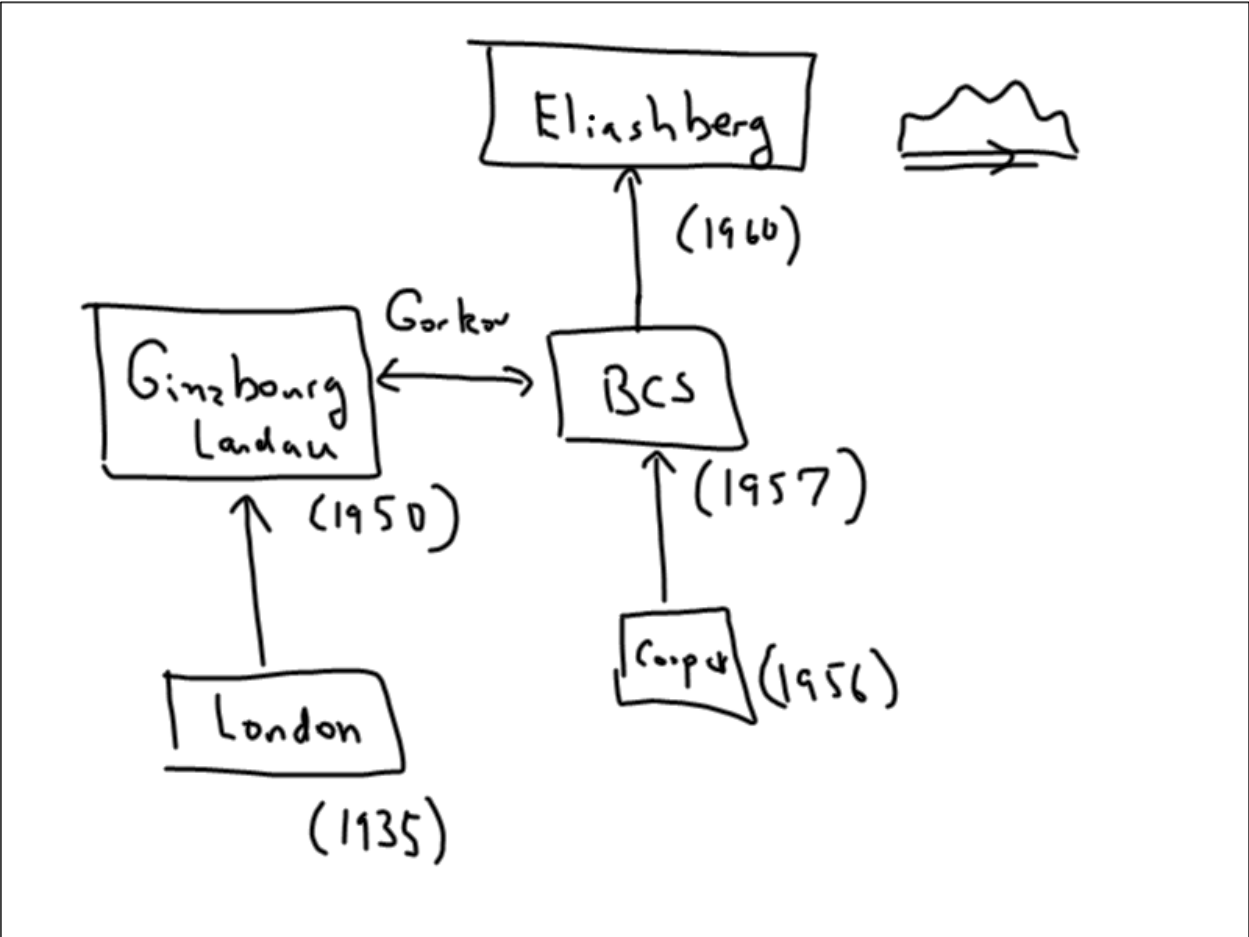
① Cooper



② "Coherence"

$$\prod_k (u_k + v_k c_{k\uparrow}^+ c_{-k\downarrow}^+) |0\rangle$$

$\uparrow$   
 $e^{i\theta_k}$





## Instabilit  de Cooper

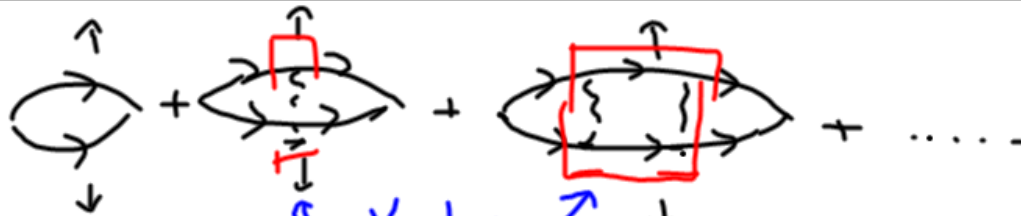
Mahan - 777   787

Richayzen - 203   208

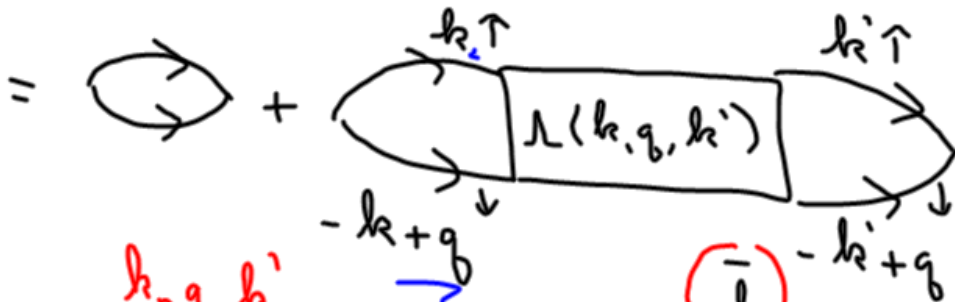
$$\langle S^z(x, \tau) S^z(x', \tau') \rangle = \text{diagram 1} + \text{diagram 2} + \dots$$

$$c_{\uparrow}^{\dagger}(\tau) c_{\uparrow}(\tau) - c_{\downarrow}^{\dagger}(\tau) c_{\downarrow}(\tau)$$

$$\langle T_{\tau} (c_{\uparrow}^{\dagger}(\tau) c_{\downarrow}^{\dagger}(\tau) \quad c_{\downarrow}(0) c_{\uparrow}(0)) \rangle$$



complètement réductible



$$\Lambda(k, q, k') = \underbrace{\dots}_{\dots} + \underbrace{\dots}_{\dots} + \dots$$

$$\boxed{\Lambda} = \underbrace{\dots}_{\dots} + \underbrace{\dots}_{\dots}$$

$$\Lambda(k, q, k') = \left\{ k - k' + \int_{-k+q}^k \Lambda(\bar{k}, q, k') \right\}$$

$$\Lambda_{r_2}(k, q, k') = V(k - k')$$

$$\rightarrow -T \sum_{i\hbar_n} \int \frac{d\bar{k}}{(2\pi)^3} \mathcal{G}_r(\bar{k}) \mathcal{G}_s(-\bar{k}+q) V(k-\bar{k}) \Lambda_{r_2}(\bar{k}, q, k')$$

$$\Lambda_{r_2}(k, q, k') = V(k - k') + \int \frac{d^3\bar{k}}{(2\pi)^3} \frac{[1 - f(S_{\bar{k}}) - f(S_{-\bar{k}+q})]}{i\hbar_n - \int_{-\bar{k}+q}^{\bar{k}} S_{\bar{k}}} V(k - \bar{k}) \Lambda_{r_2}(\bar{k}, q, k')$$

$$1 - f_1 - f_2 = (1 - f_1)(1 - f_2) - \underbrace{f_1 f_2}_{\text{Source}}$$

$$(\omega - S_{\bar{k}} - \int_{-\bar{k}+q}^{\bar{k}} S_{\bar{k}}) \Psi = V + \int V \Psi (1 - f)(1 - f)$$

$$\left[ \omega + \frac{\hbar^2}{2m} \nabla_1^2 + \frac{\hbar^2}{2m} \nabla_2^2 - \omega_m - V(r_1 - r_2) \right] \Psi(r_1, r_2) = V \text{ "Source"}$$

$$T \sum_{i\hbar_n} \frac{1}{i\hbar_n - S_{\bar{k}}} \frac{1}{-i\hbar_n - iq_n - \int_{-\bar{k}+q}^{\bar{k}} S_{\bar{k}}}$$

2. Sol. pour  $\Lambda$ .

$$V(\vec{k}-\vec{k}', i k_n - i k'_n) = \begin{cases} V(\vec{k}-\vec{k}') \text{ si} \\ |k_n| \text{ et } |k'_n| \\ < \omega_D \\ 0 \text{ autrement.} \end{cases}$$

$$V(\vec{k}-\vec{k}') = V(\vec{k} \cdot \vec{k}') \quad |k| = |k'| \sim k_F$$

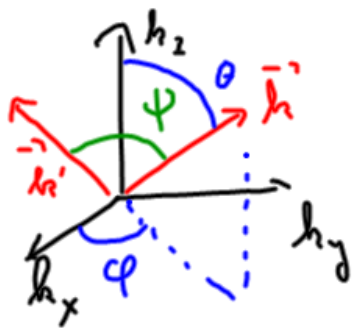
$$\Lambda(k, q, k') \xrightarrow{\gamma_{lm}^*} \Lambda(\hat{k} \cdot \hat{k}', i k_n, i k'_n) \xrightarrow{\gamma_{lm}}$$

$\vec{q} = 0 \quad i q_n = 0$

En général

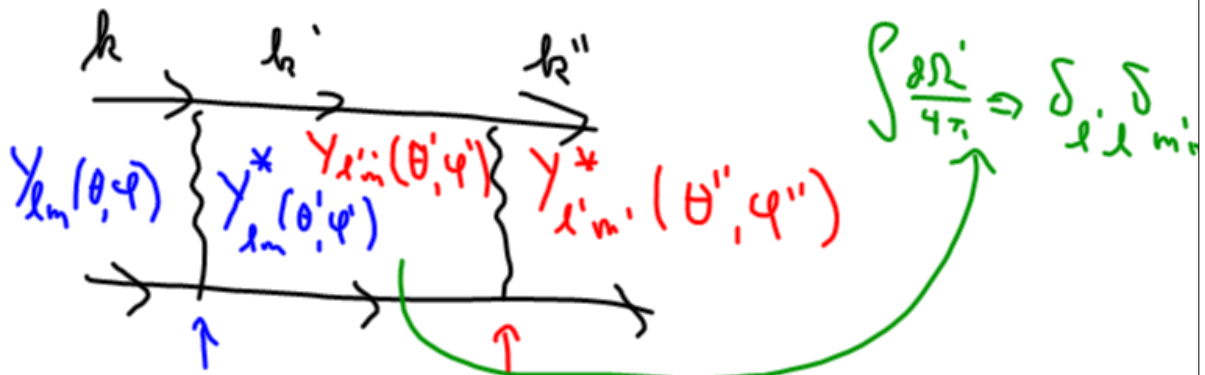
$$V(\hat{h} \cdot \hat{h}') = \sum_l V^l \mathcal{P}_l(\hat{h} \cdot \hat{h}')$$

↓  
 ↑  
 Polynômes de Legendre



$$= \sum_l V^l \mathcal{P}_l(\cos \psi)$$

$$\rightarrow \mathcal{P}_l(\cos \psi) = \frac{4\pi}{2l+1} \sum_{m=-l}^l Y_{lm}(\theta, \varphi) Y_{lm}^*(\theta', \varphi')$$



$$\mathcal{L}_{\uparrow\downarrow}^{\circ\circ}(\underline{ik}_n, \underline{ih}'_n) = V^{\circ} \theta(\omega_0 - |\underline{k}_n|) \theta(\omega_0 - |\underline{k}'_n|)$$

$$-T \sum_{\bar{k}_n} N(E_F) \int_{\tau} d\bar{\xi} \frac{V^{\circ} \theta(\omega_0 - |\underline{k}_n|) \theta(\omega_0 - |\bar{k}_n|)}{(i\bar{k}_n - \bar{\xi})(-i\bar{k}_n - \bar{\xi})}$$

$$\int d\bar{\xi} \frac{1}{\bar{\xi}^2 + \bar{k}_n^2} = \frac{1}{|\bar{k}_n|} \operatorname{arctan} \frac{\bar{\xi}}{|\bar{k}_n|} \Big|_{-\infty}^{\infty} \mathcal{L}_{\uparrow\downarrow}^{\circ\circ}(\underline{ih}_n, \underline{ih}'_n)$$

$$= \frac{\pi}{|\bar{k}_n|}$$

$$\mathcal{L}_{\uparrow\downarrow}^{\circ\circ} = V^{\circ} - \boxed{T N(E_F) \pi \sum_{\bar{k}_n = -\omega_D}^{\omega_D} \frac{1}{|\bar{k}_n|}} V^{\circ} \mathcal{L}_{\uparrow\downarrow}^{\circ\circ}$$

$a$

$$\mathcal{L}_{\uparrow\downarrow}^{\circ\circ} = V^{\circ} - a V^{\circ} \mathcal{L}_{\uparrow\downarrow}^{\circ\circ}$$

$$\mathcal{L}_{\uparrow\downarrow}^{\circ\circ} = \frac{V^{\circ}}{1 + a V^{\circ}} \quad V^{\circ} < 0$$

$$= \infty \quad \text{si } 1 + a V^{\circ} = 0$$

$$a \equiv N(E_F) \frac{2}{\pi} \int_0^{\omega_D} \frac{1}{(2n+1)A_T}$$

$$= N(E_F) \left[ \gamma + \ln \frac{2\omega_D}{\pi T} \right]$$

$$1 + aV^0 = 0 \quad \gamma \equiv 0.5772 \dots \text{Cte d'Euler}$$

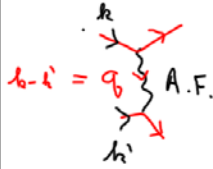
$$1 + N(E_F) \left[ \gamma + \ln \frac{2\omega_D}{\pi T_c} \right] V^0 = 0$$

$$-\frac{1}{N(E_F)V^0} = \gamma + \ln \frac{2\omega_D}{\pi T_c}$$

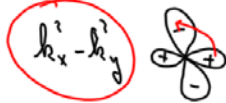
$$e^{-1/N(E_F)V^0} = e^{\gamma} \frac{2\omega_D}{\pi T_c}$$

$$T_c = \left( \frac{2e^{\gamma}}{\pi} \omega_D \right) e^{1/N(E_F)V^0}$$

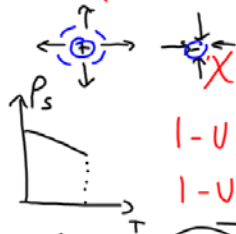
$$T_c = 1.13 \omega_D e^{1/N(E_F)V^0} \quad V^0 < 0$$



$$\vec{h} - \vec{h}' = (\pi, \pi)$$



$$\vec{\Delta}_k = \int \frac{d^2k'}{(2\pi)^2} \frac{\sqrt{|\vec{h}-\vec{h}'|} \vec{\Delta}_{k'}}{\sqrt{\epsilon_{k'}^2 + \Delta_{k'}^2}}$$

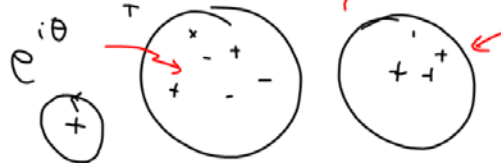


$$X \propto \ln T$$

$$1 - UX = 0$$

$$1 - UN(E_F) \ln \frac{E_F}{T} = 0$$

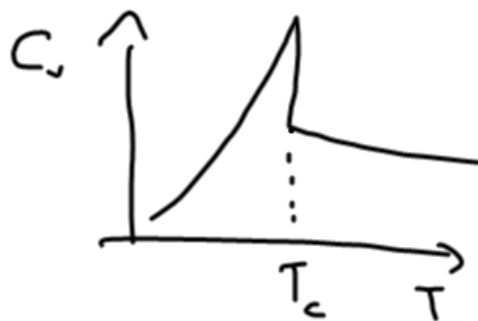
$$T_c \propto e^{-1/UN(E_F)}$$



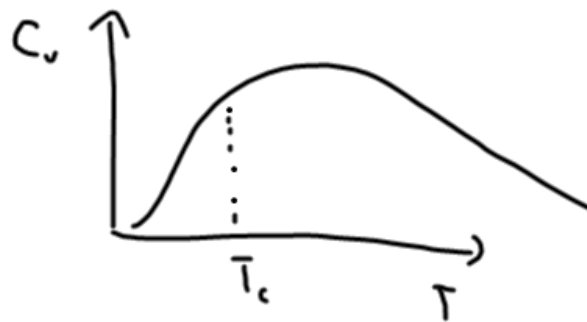
$$\langle j | j \rangle \rightarrow \langle j | \langle j \rangle + e^{-r/3}$$

$$\langle \Delta_{(1)} | \Delta_{(0)} \rangle = 1/r^{2-\eta}$$





$$\xi \propto (T - T_c)^{-1/2}$$



$$\xi \propto (T - T_c)^{1/2}$$