

Introduction to Ω *MaxEnt*, a tool for analytic continuation of Matsubara data

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Analytic Continuation

- $G(\tau)$ or $G(i\omega_n) \Rightarrow A(\omega)$?

\Rightarrow invert

$$G(\tau) = - \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \frac{e^{-\omega\tau} A(\omega)}{e^{-\beta\omega} + 1} .$$

or

$$G(i\omega_n) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \frac{A(\omega)}{i\omega_n - \omega} .$$

Conditioning Problem:

- If we discretize ω : $G = KA \Rightarrow A = K^{-1}G$
- error on A :

$$\frac{1}{\|K\| \|K^{-1}\|} \frac{\|\delta G\|}{\|G\|} \leq \frac{\|\delta A\|}{\|A\|} \leq \|K\| \|K^{-1}\| \frac{\|\delta G\|}{\|G\|}$$

- $\|K\| \|K^{-1}\|$ is large $\Rightarrow \frac{\|\delta A\|}{\|A\|}$ not bounded
- Analytic continuation unique in principle (Baym and Mermin, J.Math.Phys.1961), but unstable numerically

\Rightarrow need constraints on $A(\omega)$

Maximum entropy

- Different strategy: minimize

$$Q = \chi^2 - \alpha S$$

$$\chi^2 = \sum_{mn} (G_m - K_m A)^T C_{mn}^{-1} (G_n - K_n A)$$

$$S = - \int d\omega A(\omega) \ln \frac{A(\omega)}{D(\omega)}$$

$$C_{mn} = \langle (G_m - \langle G_m \rangle)(G_n - \langle G_n \rangle) \rangle$$

Ω MaxEnt

- Use
$$G(i\omega_n) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \frac{A(\omega)}{i\omega_n - \omega} .$$

if $A(\omega)$ is a piecewise polynomial \Rightarrow analytical integration in intervals

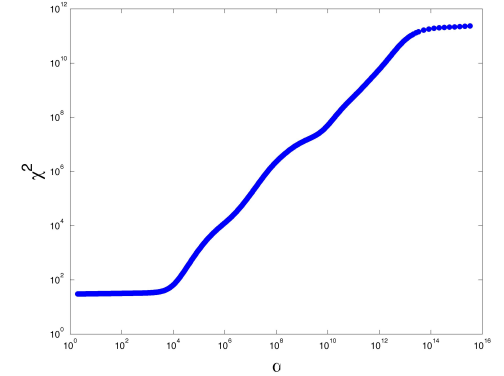
- Replace high Matsubara frequencies with constraints on moments

$$M_j = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \omega^j A(\omega)$$

- Use adapted ω grid
- Input can also be $G(\tau)$
- Treats fermionic ($A(\omega) > 0$) and bosonic ($A(\omega)/\omega > 0$) data
- General covariance matrix

Ω MaxEnt: how to choose α ?

- Compute $A(\omega)$ for large range of α :
 - Three regimes in χ^2 vs α
 \Rightarrow Optimal α located on χ^2 vs α in *log-log*
- Additional diagnostic tools to assess quality of the result:
 - $A(\omega_{sample})$ vs α
 - $\Delta G = (G_{in} - G_{out}) / \sigma$ vs ω_n
 - $\langle \Delta G_m \Delta G_{m+n} \rangle$



Why three regimes in χ^2 vs α ?

$$Q = \chi^2 - \alpha S \quad \nabla_A Q = 0$$

- Large α : $A(\omega) \approx D(\omega) \Rightarrow \chi^2 \approx \text{const}$
- Intermediate: $\alpha \searrow \Rightarrow \chi^2 \searrow$
- Small α : $\chi^2 \searrow$ very slowly (why?)

$$\alpha_j = \alpha_{j-1} - \Delta\alpha \quad \Rightarrow \quad A_j = A_{j-1} + \delta A_j$$

$$\delta A_j = \frac{\Delta\alpha}{2} \left[\tilde{\mathbf{K}}^T \tilde{\mathbf{K}} + \frac{\alpha_j}{2} \Delta\omega \mathbf{A}_{j-1}^{-1} \right]^{-1} (\Delta\omega \ln(\mathbf{D}^{-1} A_{j-1}) + d\omega)$$

All quantities are smooth in RHS at optimal α ,
but smooth δA_j cannot fit noise

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