

③ Graphene

$$h_{\text{eff}}(\vec{\tau}) = v(\tau^x \sigma^x q_x + \sigma^y q_y)$$

3.1 TRS

$$\mathcal{H} = \sum_{\vec{k}} c_{\vec{k}}^\dagger h(\vec{k}) c_{\vec{k}}$$

$$[\mathcal{H}, \Theta] = 0 \Leftrightarrow h(\vec{k}) = \Theta^{-1} h(-\vec{k}) \Theta$$

$$\Theta = U \mathbb{K}$$

$\uparrow$
complex conjugation.

* TRS in graphene?

First guess: $\Theta = \mathbb{K}$

(spinless graphene)

$$TRS \Leftrightarrow h_{\text{eff}}(\vec{q}) = h_{\text{eff}}(-\vec{q})^*$$

$$\Leftrightarrow v(\tau^z \sigma^x q_x - \sigma^y q_y)$$

$$= v(-\tau^z \sigma^x q_x - \sigma^y q_y)$$

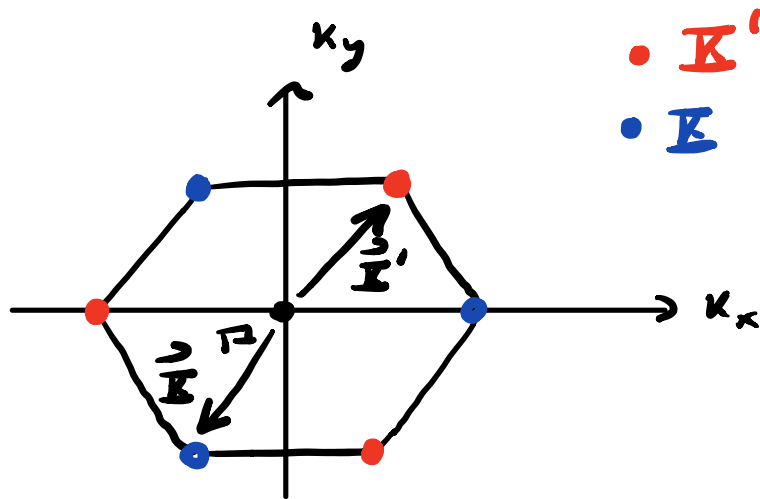
Doesn't work. $\Rightarrow$ graphene
breaks TRS ! ?

Mistake: $\Theta = \mathbb{K}$ was too naïve.
We forgot to act w/ Θ on the
valley d.o.f.

Indeed, valley $\sim$ momentum.

TR inverts momentum.

How does TR act on valley?



Momentum near K valley:

$$\vec{k} = \vec{K} + \vec{q} \quad (q \text{ small})$$

Momentum near K' valley:

$$\vec{k} = \vec{K}' + \vec{q} \quad (q \text{ small})$$

$\vec{k}$ = momentum measured from Γ .

Under TR , $\vec{k} \rightarrow -\vec{k}$.

Then,

$$(i) \quad \vec{k} = \vec{k} + \vec{q} \quad (\text{near } \mathbb{K})$$

$$\downarrow \text{TR}$$

$$-\vec{k} = -\vec{k} - \vec{q} = \vec{k}' - \vec{q} \quad (\text{near } \mathbb{K}')$$

$$(ii) \quad \vec{k} = \vec{k}' + \vec{q} \quad (\text{near } \mathbb{K}')$$

$$\downarrow \text{TR}$$

$$-\vec{k} = -\vec{k}' - \vec{q} = \vec{k} - \vec{q} \quad (\text{near } \mathbb{K})$$

So, TR flips valley and

takes $\vec{q} \rightarrow -\vec{q}$.

To take this into account,

$$\Theta = \tau^* \overline{\mathbb{K}} \quad \leftarrow \begin{array}{l} \text{complex} \\ \text{conjugation} \end{array}$$

Graphene has TRS

$$\Leftrightarrow h_{\text{eff}}(\vec{q}) = \tau^x h_{\text{eff}}(-\vec{q})^* \tau^x$$

$$\Leftrightarrow v(\tau^z \sigma^x q_x + \sigma^y q_y)$$

$$= v \tau^x (-\tau^z \sigma^x q_x + \sigma^y q_y) \tau^x$$

$$= v(\tau^z \sigma^x q_x + \sigma^y q_y) \quad \checkmark$$

$\uparrow$

$$\tau^x \tau^z \tau^x = -\tau^z$$

3.3 symmetries (II): space inversion

$$\vec{r} \rightarrow -\vec{r}$$

* Inversion operator: Π

$$\pi \vec{r} \pi^{-1} = -\vec{r} \quad (1)$$

$\uparrow$
 position
 operator

$$\pi \vec{p} \pi^{-1} = -\vec{p} \quad (2)$$

$\uparrow$
 momentum
 operator

(1) and (2) :

$$\pi [x, p_x] \pi^{-1} = [x, p_x]$$

Since $[x, p_x] = i\hbar$,

$$\pi i \pi^{-1} = i$$

$\Rightarrow \pi$ is unitary $(\pi^{-1} = \pi^\dagger)$

* $\boxed{\pi^2 = 1}$ irrespective of
spin of particles.

* A crystal has space inversion
symmetry (SIS)

$$\Leftrightarrow [\mathcal{H}, \pi] = 0.$$

$$* \quad \mathcal{H} = \sum_{\vec{k}} \sum_{\sigma\sigma'} \sum_{\alpha\alpha'}$$

$$+ c_{\vec{k}\alpha\sigma}^\dagger h_{\sigma\sigma'}^{\alpha\alpha'}(\vec{k}) c_{\vec{k}\alpha'\sigma'}$$

σ, σ' : spin

α, α' : everything else

$$SIS \Leftrightarrow \Pi \mathcal{H} \Pi^{-1} = \mathcal{H}$$

$$* \quad \Pi \mathcal{H} \Pi^{-1} = \sum_{\vec{k}} \sum_{\alpha \alpha'} \sum_{\sigma \sigma'}$$

$$\underbrace{\Pi C_{\vec{k} \alpha \sigma}^{\dagger} \Pi^{-1}}_{?} \underbrace{\Pi h_{\sigma \sigma'}^{\alpha \alpha'}(\vec{k}) \Pi^{-1}}_{?} \underbrace{\Pi}_{C_{\vec{k} \alpha' \sigma'} \Pi^{-1}}_{?}$$

$$* \quad \Pi C_{j \sigma \alpha} \Pi^{-1}$$

$$= \sum_{\alpha'} \Pi_{\alpha \alpha'} C_{-j \sigma \alpha'}$$

$$(\vec{n}_j = -\vec{n}_{-j})$$

$$\Rightarrow \Pi C_{\vec{k} \sigma \alpha} \Pi^{-1} \stackrel{\text{same as}}{=} \text{lecture 16}$$

$$= \sum_{\alpha'} \Pi_{\alpha \alpha'} C_{-\vec{k}, \sigma \alpha'}$$

$$* \quad \pi \underbrace{h_{\sigma\sigma'}^{\alpha\alpha'}(\vec{k})}_{\substack{\uparrow \\ \text{a number}}} \pi^{-1} = h_{\sigma\sigma'}^{\alpha\alpha'}(\vec{k})$$

$$* \quad \pi \mathcal{H} \pi^{-1}$$

$$= \sum_{\vec{k}} c_{-\vec{k}}^+ \pi^+ \underbrace{h(\vec{k})}_{\substack{\uparrow \\ \text{matrix}}} \underbrace{\pi}_{\substack{\uparrow \\ \text{matrix}}} c_{-\vec{k}}$$

$\uparrow$
 same as
 Lecture 16

$$= \sum_{\vec{k}} c_{+\vec{k}}^+ \pi^+ h(-\vec{k}) \pi c_{\vec{k}}$$

$$* \quad \pi \mathcal{H} \pi^{-1} = \mathcal{H} \Leftrightarrow$$

$$\boxed{h(\vec{k}) = \pi^+ h(-\vec{k}) \pi}$$

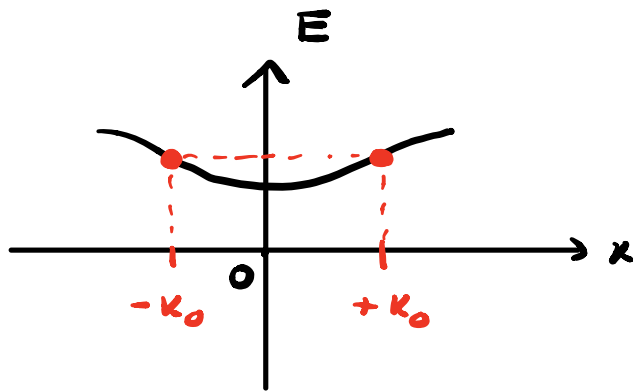
condition for SIS.

* Implication for electronic structure:

In a crystal w/ SIS, if

$|\psi_{\vec{k}}\rangle$ is an eigenstate of $h(\vec{k})$,

then $\pi |\psi_{\vec{k}}\rangle$ is an eigenstate of $h(-\vec{k})$, with same energy.



* In general,

$$h(\vec{k}) = \pi^{-1} h(-\vec{k}) \pi$$

does not imply $[h(\vec{k}), \pi] = 0$.

But, $[\ln(\vec{k}_{tr}), \pi] = 0.$

Are $|\mu_{\vec{k}_{tr}}\rangle$ and

$$\pi |\mu_{\vec{k}_{tr}}\rangle$$

the same state? yes,

because $\pi^2 = \mathbb{1}.$

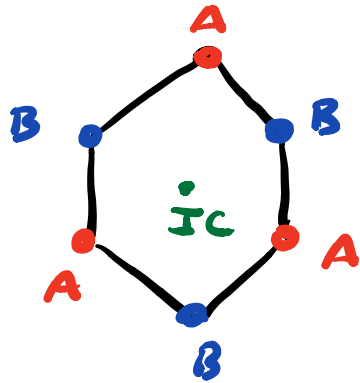
$\Rightarrow$ no Kramers degeneracy
associated to SIS.

* Does graphene have SIS?

what's π in graphene?

$$\pi \propto \tau^x \quad (t/c \mathbb{R} \leftrightarrow \mathbb{R}')$$

under space
inversion)



IC = inversion
center.

(about which
the graphene
lattice is
invariant under

$$\vec{r} \rightarrow -\vec{r})$$

Space inversion exchanges sublattices:

$$A \leftrightarrow B.$$

$$\Pi \propto \sigma^x$$

$$\text{Then, } \Pi = \sigma^x \tau^x$$

Graphene has SIS $\Leftrightarrow$

$$h_{\text{eff}}(\vec{q}) = \sigma^x \tau^x h_{\text{eff}}(-\vec{q}) \sigma^x \tau^x$$

check as an exercise that this
is indeed satisfied for

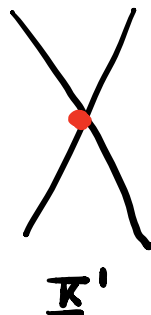
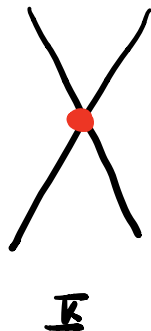
$$h_{\text{eff}}(\vec{q}) = v(\tau^z \sigma^x q_x + \sigma^y q_y)$$

$$\left(\begin{aligned} \text{use } \sigma^x \sigma^y \sigma^x &= -\sigma^y \\ \sigma^x \sigma^x \sigma^x &= \sigma^x \\ \tau^x \tau^z \tau^x &= -\tau^z \end{aligned} \right)$$

3.4 How to open a gap in graphene

Unperturbed graphene:

$$h_0(\vec{q}) = v(\tau^z \sigma^x q_x + \sigma^y q_y)$$



• Dirac
points.

Consider perturbations that preserve translational symmetry of graphene (diagonal in $\vec{q}$ and diagonal in valley index). Then, the most general perturbation:

$$\delta h(\vec{q}) = \mu_0 + \vec{u} \cdot \vec{\sigma} + \left(w_0 + \vec{w} \cdot \vec{\sigma} \right) \tau^z$$

where

$$\mu_i, w_i \in \mathbb{R}$$

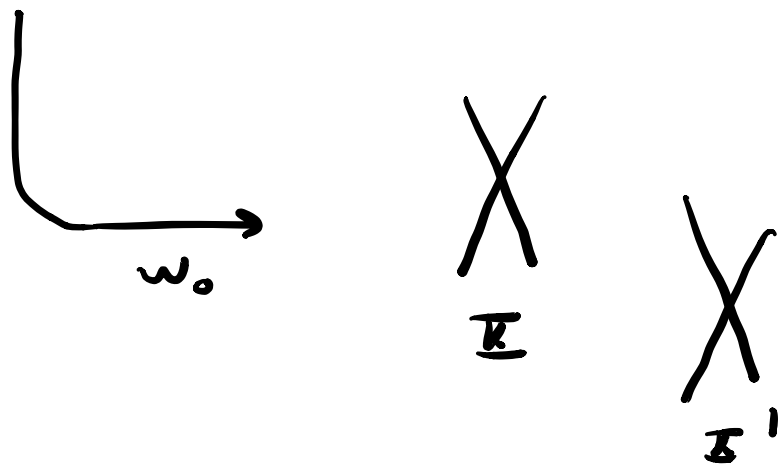
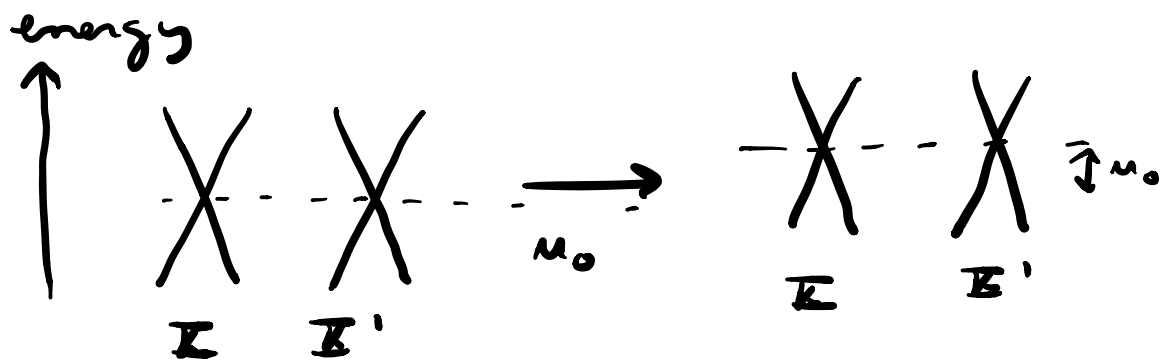
$$i = 0, x, y, z$$

$$\mu_i = \mu_i(\vec{q})$$

$$w_i = w_i(\vec{q})$$

* Which terms open a gap at Dirac points?

(i) μ_0, ω_0 : shift of Dirac points w/o opening gap.



(ii) $m_x(\vec{q})$:

shifts position of Dirac point. No gap.

$$h_0 + \delta h =$$

$$= \begin{cases} v \left[\sigma^x \left(q_x + \frac{m_x}{v} \right) + \sigma^y q_y \right] & (\mathbb{K}) \\ v \left[\sigma^x \left(-q_x + \frac{m_x}{v} \right) + \sigma^y q_y \right] & (\mathbb{K}') \end{cases}$$

Unperturbed Dirac points

are at $q_x = 0$, $q_y = 0$.

Perturbed Dirac points are at

$$q_x + \frac{m_x}{v} = 0, \quad q_y = 0 \quad (\mathbb{K})$$

$$q_x - \frac{\mu_x}{v} = 0, \quad q_y = 0 \quad (\mathbb{K}')$$

(iii) similarly, μ_y, w_x and w_y shift Dirac points but do not open gaps.

(iv) μ_z, w_z : these open gaps.

$$h = \begin{cases} v(\sigma^x q_x + \sigma^y q_y) + m_{\mathbb{K}} \sigma^z & (\mathbb{K}) \\ v(-\sigma^x q_x + \sigma^y q_y) + m_{\mathbb{K}'} \sigma^z & (\mathbb{K}') \end{cases}$$

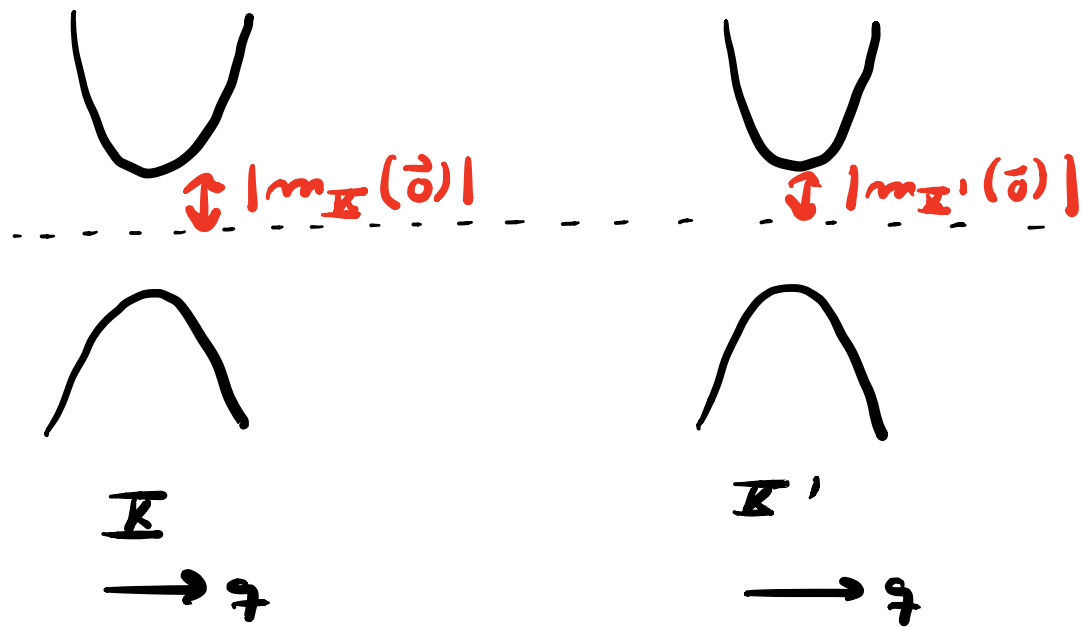
$m_{\mathbb{K}}$ and $m_{\mathbb{K}'}$: mass terms

$$m_{\mathbb{K}} = \mu_z + w_z$$

$$m_{\mathbb{K}'} = \mu_z - w_z$$

$$E_{q, \pm}^{(\mathbf{K})} = \pm \sqrt{v^2(q_x^2 + q_y^2) + m_{\mathbf{K}}^2}$$

$$E_{q, \pm}^{(\mathbf{K}')} = \pm \sqrt{v^2(q_x^2 + q_y^2) + m_{\mathbf{K}'}^2}$$



From here on, let's concentrate only on u_z and w_z .

$$\delta h(\vec{q}) = u_z(\vec{q}) \sigma^z + w_z(\vec{q}) \sigma^z \tau^z$$

$u_z(\vec{0})$: Semenoff mass

G. Semenoff, PRL 53, 2449 (1984)

$w_z(\vec{0})$: Haldane mass.

D. Haldane, PRL 61, 2015 (1988)

$$* \quad \text{TRS} \Leftrightarrow h(\vec{q}) = \tau^x h(-\vec{q})^* \tau^x$$

$$\Leftrightarrow \delta h(\vec{q}) = \tau^x \delta h(-\vec{q})^* \tau^x$$

$$\Leftrightarrow \begin{cases} u_z(\vec{q}) = u_z(-\vec{q}) \\ w_z(\vec{q}) = -w_z(-\vec{q}) \end{cases}$$

$$\Rightarrow w_z(\vec{0}) = 0 \Rightarrow \text{no Haldane mass.}$$

Can still have a Semenov mass.

$$* \text{ SIS} \Leftrightarrow h(\vec{q}) = \tau^x \sigma^x h(-\vec{q}) \tau^x \sigma^x$$

$$\Leftrightarrow \delta h(\vec{q}) = \tau^x \sigma^x \delta h(-\vec{q}) \tau^x \sigma^x$$

$$\Leftrightarrow \begin{cases} u_z(\vec{q}) = -u_z(-\vec{q}) \\ w_z(\vec{q}) = w_z(-\vec{q}) \end{cases}$$

$$\Rightarrow u_z(\vec{0}) = 0 \Rightarrow \text{no Semenov mass.}$$

Can still have a Haldane mass.

* If both SIS and TRS are present:

$$u_z = w_z = 0 \text{ all } \vec{q}.$$

$$\Rightarrow \text{no gap.}$$

Need to break either SIS
or TRS to get graphene.

Break only $SIS \Rightarrow$ Semeroff mass

" only $TRS \Rightarrow$ Haldane "

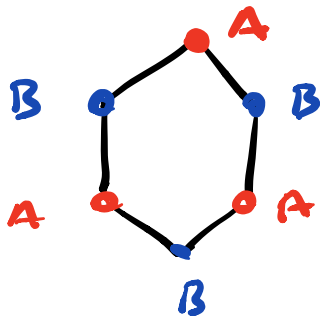
* Physical origin of Semeroff mass?

$$\delta h = u_z(\vec{q}) \sigma^z$$

Staggered potential

(opposite sign for A and B
sublattices). Realised if

A and B
are different
atoms.



example: BN.

* Physical origin of Haldane mass?

Imaginary second neighbor
hopping.

$$\mathcal{H} = i \sum_{\langle\langle i,j \rangle\rangle} \underbrace{t_2}_{\substack{\uparrow \\ \text{real}}} c_i^\dagger c_j + \text{h.c.}$$