

Résumé

$$\sum_{i=1}^n O_i \underline{|\alpha_1\rangle |\alpha_2\rangle |\alpha_3\rangle \dots |\alpha_n\rangle}$$

$$\langle \alpha' | O_i | \alpha \rangle = \langle \alpha' | \alpha \rangle \langle \alpha | O_i | \alpha \rangle$$

$$\sum_{i=1}^n \langle \alpha_i | O_i | \alpha_i \rangle \left(\langle \alpha'_1 | \langle \alpha'_2 | \dots \langle \alpha'_n | \right) \underline{|\alpha_1\rangle |\alpha_2\rangle \dots |\alpha_n\rangle}$$

$C_\alpha^+ C_\alpha$

$$\sum_\alpha \langle \alpha | O | \alpha \rangle C_\alpha^+ C_\alpha$$
$$|\alpha\rangle = \int dx |x\rangle \langle x | \alpha \rangle$$
$$C_\alpha^+ = \int dx \Psi^+(x) \langle x | \alpha \rangle$$

$\int dx dy \Psi^+(y) \langle y | O | x \rangle \Psi(x)$

$$\langle y | O | x \rangle = \delta(x-y)$$
$$\int dx \Psi^+(x) \Psi(x)$$

Exercices:

$$\int dx dy \psi^+(y) \langle y | 0 \rangle \psi(x)$$

i) 0 est un potentiel à 1 corps.

$$\langle y | V | x \rangle = V(x) \delta(x-y)$$

$$\int dx \psi^+(x) V(x) \psi(x)$$

$$ii) \langle y | \frac{p^2}{2m} | x \rangle = \int \frac{dp}{2\pi} \langle y | \frac{p^2}{2m} | p \rangle \langle p | x \rangle$$

$$= \int \frac{dp}{2\pi} \frac{\hbar^2 p^2}{2m} \langle y | p \rangle \langle p | x \rangle$$

$$= \int \frac{dp}{2\pi} \frac{\hbar^2 p^2}{2m} e^{i p y - i p x}$$

$$= \int \frac{dp}{2\pi} \frac{\hbar^2}{2m} \left(-\frac{\partial^2}{\partial x^2} \right) e^{i p (y-x)}$$

$$= -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \delta(x-y)$$

$$\int dx dy \psi^+(y) \left(-\frac{\hbar^2}{2m} \right) \frac{d^2}{dx^2} \delta(x-y) \psi(x)$$

$$-\frac{\hbar^2}{2m} \int dx \psi^+(x) \frac{d^2}{dx^2} \psi(x)$$

$$H = \int d^3r \psi^\dagger(\vec{r}) \mathcal{D} \psi(\vec{r})$$

$$\mathcal{D} = -\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r})$$

$$\mathcal{D} \varphi_n(\vec{r}) = E_n \varphi_n(\vec{r})$$

$$\psi^\dagger(\vec{r}) = \sum_n c_n^\dagger \varphi_n^*(\vec{r}) \leftarrow$$

$$|\vec{r}\rangle = \sum_n |n\rangle \langle n|\vec{r}\rangle \leftarrow$$

$$\langle \vec{r}|n\rangle = \varphi_n(\vec{r})$$

$$\begin{aligned} \mathcal{D}\psi &= \mathcal{D} \sum_n c_n \varphi_n(\vec{r}) \\ &= \sum_n c_n E_n \varphi_n(\vec{r}) \end{aligned}$$

$$\psi^\dagger \mathcal{D}\psi = \sum_{n_1 n_2} c_{n_1}^\dagger \varphi_{n_1}^*(\vec{r}) E_{n_2} c_{n_2} \varphi_{n_2}(\vec{r})$$

$$\sum_{n_1 n_2} E_{n_2} c_{n_1}^\dagger c_{n_2} \int d^3r \varphi_{n_1}^*(\vec{r}) \varphi_{n_2}(\vec{r})$$

$$= \sum_{n_1 n_2} E_{n_2} c_{n_1}^\dagger c_{n_2} \delta_{n_1 n_2}$$

$$= \sum_n E_n c_n^\dagger c_n$$

$$|\Psi\rangle = c_{n_1}^\dagger c_{n_2}^\dagger |0\rangle$$

$$|n\rangle = \int d^3r |\vec{r}\rangle \langle \vec{r} | n \rangle$$

$$c_m^\dagger = \int d^3r \psi^\dagger(\vec{r}) \varphi_m(\vec{r})$$

$$c_n = \int d^3r \psi(\vec{r}) \varphi_n^*(\vec{r})$$

$$[c_n, c_m^\dagger] =$$

$$\int d^3r d^3r' \varphi_n^*(\vec{r}) \varphi_m(\vec{r})$$

$$\delta(\vec{r}-\vec{r}') \left[\psi(\vec{r}') \psi^\dagger(\vec{r}) - \psi^\dagger(\vec{r}) \psi(\vec{r}') \right]$$

$$= \int d^3r d^3r' \varphi_n^*(\vec{r}) \varphi_m(\vec{r}) \delta(\vec{r}-\vec{r}')$$

$$= \delta_{m,n} = \int d^3r \varphi_n^*(\vec{r}) \varphi_m(\vec{r})$$

$$U_{r_n} = \langle \vec{r} | n \rangle$$

$$U_{m_r}^\dagger = \langle n | \vec{r} \rangle$$

$$\int d^3r U_{m_r}^\dagger U_{r_n} = \delta_{mn}$$

5.? Opérateurs à 2 corps.

exemple 1^{ère} quant.

$$\frac{1}{2} \sum_{\substack{i=1 \\ j=1 \\ i \neq j}}^{n,n} V(x_i - x_j) \quad \text{énergie potentielle}$$

Sans antisymétrie

$$\frac{1}{2} \sum_{\substack{i,j \\ i \neq j}} \int d\vec{x}_i d\vec{x}_j \frac{|\Psi(\vec{x}_i)|^2 e e |\Psi(\vec{x}_j)|^2}{|\vec{x}_i - \vec{x}_j|}$$

Base diagonale:

$$\langle \alpha' | \langle \beta' | V | \alpha \rangle | \beta \rangle$$

$$= V_{\alpha\beta} \langle \alpha' | \alpha \rangle \langle \beta' | \beta \rangle$$

ex.

$$\langle x' | \langle y' | V | x \rangle | y \rangle$$

$$= V(x-y) \langle x' | x \rangle \langle y' | y \rangle$$

$$\langle \alpha'_1 | \langle \alpha'_2 | \dots \langle \alpha'_n |$$

$$\frac{1}{2} \sum_{\substack{i,j \\ i \neq j}} V_{ij} | \alpha_i \rangle | \alpha_j \rangle \dots | \alpha_n \rangle$$

$$= \frac{1}{2} \sum_{\substack{\alpha_i, \alpha_j \\ i \neq j}} V_{\alpha_i, \alpha_j} \left(\langle \alpha'_1 | \langle \alpha'_2 | \dots \langle \alpha'_n | | \alpha_i \rangle \dots | \alpha_j \rangle \right)$$

$$= \frac{1}{2} \sum_{\substack{\alpha_i, \alpha_j \\ i \neq j}} V_{\alpha_i, \alpha_j} \left(\begin{array}{ccc|ccc} 1 & 1 & 1 & & & \\ & & & 1 & 1 & 1 \\ \hline & & & 2 & 2 & 2 \\ & & & 2 & 2 & 2 \end{array} \right)$$

$$= \langle \alpha'_1 \dots \alpha'_n | \frac{1}{2} \sum_{\alpha\beta} V_{\alpha\beta} \left(\begin{array}{cccc} c_\alpha^+ & c_\alpha^+ & c_\beta^+ & c_\beta^+ \\ \hline & & -\delta_{\alpha\beta} & c_\alpha^+ c_\alpha^+ \end{array} \right) | \alpha_1 \dots \alpha_n \rangle$$

$$V_{ij} | \alpha_i \rangle | \alpha_j \rangle \dots | \alpha_n \rangle$$

$$V_{\alpha_i, \alpha_j} | \alpha_i \rangle | \alpha_j \rangle \dots | \alpha_n \rangle$$

$$c_{\alpha}^{+} c_{\alpha} c_{\beta}^{+} c_{\beta} - \delta_{\alpha\beta} c_{\alpha}^{+} c_{\alpha}$$

$$c_{\alpha}^{+} (\delta_{\alpha\beta} + c_{\beta}^{+} c_{\alpha}) c_{\beta} - \delta_{\alpha\beta} c_{\alpha}^{+} c_{\alpha}$$

$$= \pm c_{\alpha}^{+} c_{\beta}^{+} c_{\alpha} c_{\beta}$$

$$= c_{\alpha}^{+} c_{\beta}^{+} c_{\beta} c_{\alpha}$$

Dans la base diagonale:

Opérateur à 2 corps

$$\frac{1}{2} \sum_{\alpha\beta} V_{\alpha\beta} c_{\alpha}^{\dagger} c_{\beta}^{\dagger} c_{\beta} c_{\alpha}$$

$$c_{\alpha}^{\dagger} = \int dx' \Psi^{\dagger}(x') \langle x' | \alpha \rangle$$

$$= \frac{1}{2} \sum_{\alpha\beta} V_{\alpha\beta} \int dx' dy' dx dy \Psi^{\dagger}(x') \langle x' | \alpha \rangle$$

$$\Psi^{\dagger}(y') \langle y' | \beta \rangle \langle \beta | x \rangle \Psi(x) \langle x | y \rangle \Psi(y)$$

Note:

$$\frac{1}{2} \sum_{\alpha\beta} \langle x' | \alpha \rangle \langle y' | \beta \rangle V_{\alpha\beta} \langle \beta | x \rangle \langle x | y \rangle$$

$$= \frac{1}{2} \langle x' | \langle y' | V | y \rangle | x \rangle$$

$$\frac{1}{2} \int dx' dy' dx dy \langle x' | \langle y' | V | y \rangle | x \rangle$$

$$\Psi^{\dagger}(x') \Psi^{\dagger}(y') \Psi(x) \Psi(y)$$

$$\langle x' | \langle y' | V | y \rangle | x \rangle = V(x-y) \langle x' | y \rangle \langle y' | x \rangle$$

$$= \frac{1}{2} \int dx dy V(x-y) \Psi^{\dagger}(y) \Psi^{\dagger}(x) \Psi(x) \Psi(y)$$

$$\langle \alpha' | \langle \beta' | V | \alpha \rangle | \beta \rangle = V_{\alpha\beta} \frac{\langle \alpha' | \alpha \rangle}{\langle \beta' | \beta \rangle}$$

$$= V_{\alpha' \beta', \alpha \beta}$$

Formel:

$$H = \sum_a \int d^3r \underbrace{\Psi_a^\dagger(\vec{r})}_{\text{red}} \left(-\frac{\hbar^2}{2m} \nabla^2 + V(r) \right) \underbrace{\Psi_a(r)}_{\text{red}}$$

Méc. class:

$$\dot{p} = -\frac{\partial H}{\partial q} \quad \dot{q} = \frac{\partial H}{\partial p}$$

$$[Q, P] = i\hbar$$

Essayons

$$q \rightarrow \Psi$$

$$\dot{p} = -\frac{\partial H}{\partial q} = - \left[-\frac{\hbar^2}{2m} \nabla^2 + V(r) \right] \Psi^\dagger$$

$$= i\hbar \frac{\partial \Psi^\dagger}{\partial t}$$

$$\dot{p} \rightarrow i\hbar \frac{\partial \Psi^\dagger}{\partial t}$$

$$p \rightarrow i\hbar \Psi^\dagger$$

$$[\Psi_r, i\hbar \Psi_{r'}^\dagger] = \delta_{r,r'} i\hbar$$

$$[\Psi(r), \Psi^\dagger(r')] = \delta(\vec{r} - \vec{r}')$$

Approximation de Hartree-Fock

"Champ moyen" pour les systèmes à N particules

Intuitif:

Une particule Hartree:

$$-\frac{\hbar^2}{2m} \nabla^2 \varphi_n(\vec{r}) + \left[V(r) + U(r) \right] \varphi_n(r) = \epsilon_n \varphi_n(r)$$

Particules en interaction:

Interaction de Coulomb

$$\begin{aligned} \rightarrow U(\vec{r}) &= \int d^3 r' \frac{e^2}{|\vec{r} - \vec{r}'|} \rho(\vec{r}') \\ &= \int d^3 r' \frac{e^2}{|\vec{r} - \vec{r}'|} \sum_n^{\text{occupés}} |\varphi_n(\vec{r}')|^2 \end{aligned}$$

Principe variationnel (Ritz)

$$\langle \Psi | H | \Psi \rangle \geq \langle \Psi_0 | H | \Psi_0 \rangle$$

$$s: \langle \Psi | \Psi \rangle = 1 \quad \leftarrow$$

$$H | \Psi \rangle = \sum_{i=0}^{\infty} a_i | \Psi_i \rangle E_i$$

$$\langle \Psi | H | \Psi \rangle = \sum_{i=0}^{\infty} E_i |a_i|^2 \geq E_0$$

$$\geq \sum_{i=0}^{\infty} E_0 |a_i|^2$$

$$= E_0 \sum_{i=0}^{\infty} |a_i|^2$$

$$= E_0$$

$$H = \int d^3r \underbrace{\Psi^\dagger(\vec{r})}_{\downarrow} \left[-\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r}) \right] \underbrace{\Psi(\vec{r})}_{\downarrow}$$

$$+ \frac{1}{2} \int d^3r d^3r' \underbrace{\Psi^\dagger(\vec{r})}_{\uparrow} \underbrace{\Psi^\dagger(\vec{r}')}_{\uparrow} \underbrace{U(|\vec{r}-\vec{r}'|)}_{\uparrow} \underbrace{\Psi(\vec{r})}_{\uparrow} \underbrace{\Psi(\vec{r}')}_{\uparrow}$$

$$|\Phi_0\rangle = \overbrace{a_1^\dagger a_2^\dagger \dots a_n^\dagger} |0\rangle$$

$$\Psi^\dagger(\vec{r}) = \sum_n a_n^\dagger \varphi_n^*(\vec{r})$$

$$\Psi(\vec{r}) = \sum_n \varphi_n(\vec{r}) a_n$$

$$|\phi\rangle = \sum_i C_i |\phi_i\rangle$$

A faire

$$\frac{\partial}{\partial \varphi_n(\vec{r}_i)} \langle \phi_0 | H | \phi_0 \rangle = 0$$

Contrainte

$$\int d^3r \varphi_n^*(\vec{r}) \varphi_m(\vec{r}) = \delta_{nm}$$