

1- Condensation de Bose-Einstein
 T finie.

2- Effet des interactions ?

3- OHDLP - ODLRO
état cohérents

4- Superfluidité

5- Pièges magnétiques + optiques

6- Equ. Gross - Pitaevskii

1. B. F.

$$|\Omega\rangle = \frac{(a_{k=0}^+)^N}{\sqrt{N!}} |0\rangle \leftarrow$$

$$n_k = \frac{1}{e^{\beta(\epsilon_k - \mu)} - 1} = \langle a_k^+ a_k \rangle$$

$$\epsilon_k = \frac{\hbar^2 k^2}{2m} \leftarrow$$

$$N = \sum_k n_k = \sum_k \frac{1}{e^{\beta(\epsilon_k - \mu)} - 1}$$

$N_0 \equiv$ # dans fond.

$$N_0 = \frac{1}{e^{-\beta\mu} - 1} = \frac{1}{-\beta\mu}$$

$$\mu = - \frac{k_B T}{N_0} \rightarrow 0$$

$$\epsilon_{k_1} = \frac{\hbar^2}{2m} \left(\frac{2\pi}{L} \right)^2$$



$$n_{k_1} = \frac{1}{e^{-\beta \left(\frac{\hbar^2 (2\pi)^2}{2m L^2} - \frac{k_B T}{L^3 n_0} \right)} - 1}$$

$$N = N_0 + \sum_{\substack{k \\ \neq 0}} \frac{1}{e^{\beta(\epsilon_k - \mu)} - 1}$$

$T > T_c$ $N_0 = 0$

$$N = N_0 + V \int \frac{d^3k}{(2\pi)^3} \frac{1}{e^{\beta\epsilon_k} - 1}$$

$$= N_0 + V \frac{4\pi}{(2\pi)^3} \int_0^\infty k^2 dk \frac{1}{e^{\beta\epsilon_k} - 1}$$

$$\epsilon = \frac{\hbar^2 k^2}{2m} \quad d\epsilon = \frac{\hbar^2 k}{m} dk \leftarrow$$

$$k = \sqrt{\frac{2m\epsilon}{\hbar^2}}$$

$$N = N_0 + V \frac{1}{2\pi^2} \int_0^\infty \frac{m}{\hbar^2} d\epsilon \sqrt{\frac{2m}{\hbar^2} \epsilon}$$

$$n = n_0 + \left(\frac{m}{\hbar^2}\right)^{3/2} \frac{1}{\sqrt{2}\pi^2} \int_0^\infty \frac{e^{\beta\epsilon} - 1}{\beta^{3/2} e^{\beta\epsilon} - 1} d\epsilon$$

$$y = \beta\epsilon$$

$$= n_0 + \left(\frac{m\hbar_0 T}{\hbar^2 2\pi}\right)^{3/2} \zeta\left(\frac{3}{2}\right)$$

$$\int_0^\infty dy y^{1/2} \frac{e^{-y}}{1 - e^{-y}}$$

$$= \int_0^\infty dy y^{1/2} e^{-y} (1 + e^{-y} + e^{-2y} + e^{-3y} + \dots)$$

$$= \int_0^\infty dy y^{1/2} e^{-y} \sum_{p=0}^\infty e^{-py} \quad \frac{p^{-3/2}}{p^{3/2}}$$

$$= \sum_{p=1}^\infty \frac{1}{p^{3/2}} \int_0^\infty dx x^{1/2} e^{-x}$$

$$= \zeta\left(\frac{3}{2}\right) \Gamma\left(\frac{3}{2}\right)$$

$$= 2.612 \frac{\sqrt{\pi}}{2}$$

$$n = n_0 + \left(\frac{k_B T m}{2\pi \hbar^2} \right)^{3/2} \zeta(3/2)$$

fixe n_0 \downarrow *change avec T*

$$\frac{n_0}{n} = 1 - \frac{1}{n} \left(\frac{k_B T m}{2\pi \hbar^2} \right)^{3/2} \zeta(3/2)$$

$= 1$

$$\frac{k_B T_c m}{2\pi \hbar^2} = \left(\frac{n}{\zeta(3/2)} \right)^{2/3}$$

$$k_B T_c = \frac{2\pi \hbar^2}{m} \left(\frac{n}{\zeta(3/2)} \right)^{2/3}$$

$$\frac{n_0}{n} = 1 - \left(\frac{T}{T_c} \right)^{3/2}$$

\leftarrow

2. Stabilité sous l'effet des interactions

$$|\Omega_N\rangle = \frac{(a_{k=0}^+)^N}{\sqrt{N!}} |0\rangle \quad \text{CBE}$$

$$|\Omega_{N/2, N/2}\rangle = \left(\frac{a_{k=0}^+}{\sqrt{N/2}} \right)^{N/2} \left(\frac{a_{k=k_1}^+}{\sqrt{N/2}} \right)^{N/2} |0\rangle$$

$$H_{\text{int}} = \frac{1}{2} \int d^3r d^3r' U(\vec{r}-\vec{r}') \psi^\dagger(\vec{r}) \psi^\dagger(\vec{r}') \psi(\vec{r}) \psi(\vec{r}')$$

$$\langle \Omega_N | H_{\text{int}} | \Omega_N \rangle$$

$$\langle \Omega_{N/2, N/2} | H_{\text{int}} | \Omega_{N/2, N/2} \rangle$$

$$\psi(\vec{r}) = \frac{1}{\sqrt{V}} \sum_{\vec{k}} e^{i\vec{k}\cdot\vec{r}} a_{\vec{k}}$$

$$\begin{aligned} a|n\rangle &= \sqrt{n} |n-1\rangle \\ a^\dagger |n\rangle &= \sqrt{n+1} |n+1\rangle \end{aligned}$$

$$\psi(\vec{r}) |\Omega_N\rangle = \sqrt{\frac{N}{V}} |\Omega_{N-1}\rangle$$

$$\psi(\vec{r}') \psi(\vec{r}) |\Omega_N\rangle = \frac{N}{V} |\Omega_{N-2}\rangle$$

$$\langle \Omega_N | \psi^\dagger \psi^\dagger \psi \psi | \Omega_N \rangle = \frac{N^2}{V^2}$$

$$\left(\frac{1}{2} \int d^3r d^3r' U(\vec{r}-\vec{r}') \right) n^2 = V g n^2$$

$$\psi(\vec{r}) |\Omega_{N/2, N/2}\rangle = \sqrt{\frac{N}{2}} |\Omega_{N/2-1, N/2}\rangle$$

$$+ \sqrt{\frac{N}{2}} e^{i\vec{k}_1 \cdot \vec{r}} |\Omega_{N/2, N/2-1}\rangle$$

$$\psi^\dagger(\vec{r}) \psi(\vec{r}') |\Omega_{N/2, N/2}\rangle$$

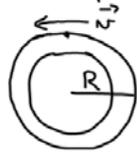
$$= \frac{N}{2} |\Omega_{N/2, N/2}\rangle + \frac{N}{2} e^{-i\vec{k}_1 \cdot \vec{r}} |\Omega_{N/2-1, N/2+1}\rangle$$

$$+ \frac{N}{2} e^{i\vec{k}_1 \cdot \vec{r}} |\Omega_{N/2+1, N/2-1}\rangle + \frac{N}{2} e^{i\vec{k}_1 \cdot \vec{r} - i\vec{k}_1 \cdot \vec{r}'} |\Omega_{N/2, N/2}\rangle$$

$$\left(\langle \Omega_{N/2, N/2} | \psi^\dagger(\vec{r}') \psi(\vec{r}') \right) \left(\psi(\vec{r}) \psi(\vec{r}') | \Omega_{N/2, N/2} \rangle \right)$$

$$= n^2 + \left(\frac{n}{2} \right)^2 e^{i\vec{k}_1 \cdot (\vec{r}' - \vec{r})} + \left(\frac{n}{2} \right)^2 e^{-i\vec{k}_1 \cdot (\vec{r}' - \vec{r})}$$

4. Superfluidité, vitesse critique



$$L = \frac{1}{2} I \omega^2$$

$$I = m R^2$$

Du pt de vue des parois \vec{v}

$$(\vec{p} - m\vec{v})^2 = (\vec{p}' - m\vec{v}')^2$$

↑ particule dans le Lab.

$$p^2 - 2\vec{p} \cdot m\vec{v} + (m\vec{v})^2 = p'^2 - 2\vec{p}' \cdot m\vec{v}' + (m\vec{v}')^2$$

$p = 0$ pour condensat.

$$p'^2 = 2\vec{p}' \cdot m\vec{v}'$$

Pour obtenir un superfluide rajouter les interactions

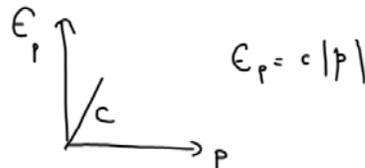
$$\Psi(\vec{r}) = \Psi_0 + \delta\Psi(\vec{r})$$

↑ Condensat ↑ Partie add.

Chs Amett à l'ordre $(\delta\Psi)^2$ (e.g. $\Psi_0^2 (\delta\Psi)^2$)
 $a^\dagger a$, $a^\dagger a^\dagger$, $a a$

Transformation de Bogolubov

→ on trouve

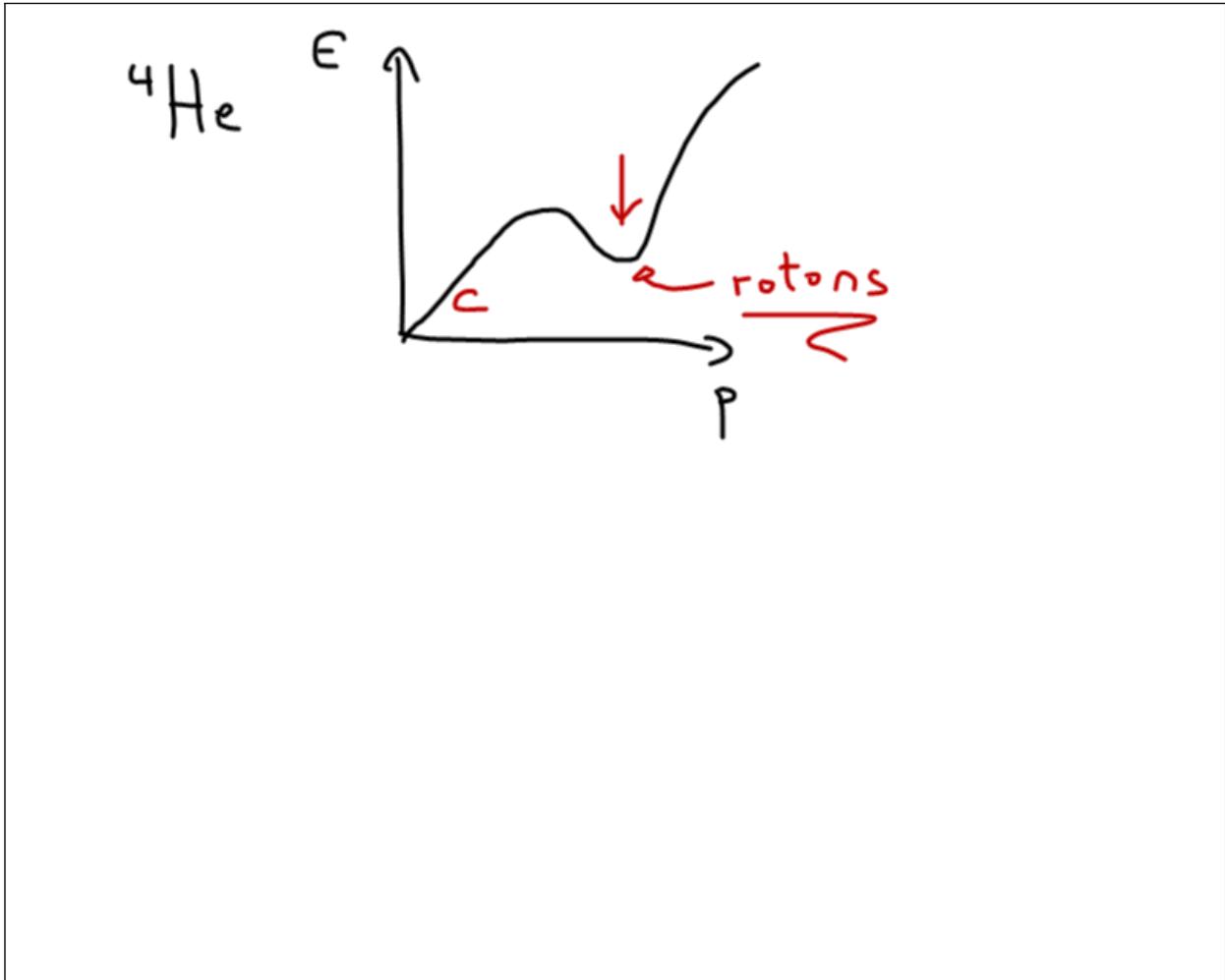


$$\text{où } c = \left(\frac{\hbar^2 n_0 g}{m} \right)^{1/2}$$

$$U(\vec{r}-\vec{r}') = g \delta(\vec{r}-\vec{r}')$$

Vitesse critique $v_c \equiv c$

Pour $v < v_c$



3. Ordre hors diagonal à longue portée

Condensat de B. E.

$$\lim_{r, r' \rightarrow \infty} \langle \Omega_N | \psi^\dagger(\vec{r}) \psi(\vec{r}') | \Omega_N \rangle$$

$$\langle \Omega_N | \frac{1}{V} \sum_{\mathbf{k}} \sum_{\mathbf{k}'} e^{-i\mathbf{k} \cdot \mathbf{r} + i\mathbf{k}' \cdot \mathbf{r}'} a_{\mathbf{k}}^\dagger a_{\mathbf{k}'} | \Omega_N \rangle = n_0$$

Fluide normal :

$$\langle \Omega | \frac{1}{V} \sum_{\mathbf{k}} e^{-i(\mathbf{r}-\mathbf{r}') \cdot \mathbf{k}} a_{\mathbf{k}}^\dagger a_{\mathbf{k}} | \Omega \rangle$$

$$= \frac{1}{V} \sum_{\mathbf{k}} e^{-i\mathbf{k} \cdot (\mathbf{r}-\mathbf{r}')} n_{\mathbf{k}} = n(\vec{r}-\vec{r}') \underset{r, r' \rightarrow \infty}{=} 0$$

Dans l'état superfluide, $\psi(\vec{r})$ a une valeur moyenne non nulle

$$\langle \Omega_s | \psi(\mathbf{r}) | \Omega_s \rangle = \sqrt{n_0} e^{i\varphi}$$

Avec ce paramètre d'ordre on brise la conservation du # total de particules

5. Pièges magnétiques et optiques
pour les atomes froids
et observation de la cond.
de B.E.

^4He $T_c \approx 2.2\text{K}$

^3He $T_c \sim 3\text{mK}$

Pièges magnétique.

