

$$H_{BCS} = \sum_{p, \sigma} \sum_p c_{p\sigma}^\dagger c_{p\sigma}$$

$$+ \sum_p \left( \Delta_p^\dagger c_{p\uparrow}^\dagger c_{-p\downarrow}^\dagger + \Delta_p c_{-p\downarrow} c_{p\uparrow} \right)$$

$$= \sum_p \left( c_{p\uparrow}^\dagger \quad c_{-p\downarrow} \right) \begin{pmatrix} \epsilon_p & \Delta_p \\ \Delta_p^\dagger & -\epsilon_{-p} \end{pmatrix} \begin{pmatrix} c_{p\uparrow} \\ c_{-p\downarrow}^\dagger \end{pmatrix}$$

$$\begin{pmatrix} \alpha_{p\uparrow} \\ \alpha_{-p\downarrow}^\dagger \end{pmatrix} = \begin{pmatrix} u_p & v_p \\ -v_p^\dagger & u_p \end{pmatrix} \begin{pmatrix} c_{p\uparrow} \\ c_{-p\downarrow}^\dagger \end{pmatrix} \quad \leftarrow$$

$$\begin{cases} u_p = \frac{1}{\sqrt{2}} \left( 1 + \frac{\epsilon_p}{E_p} \right)^{1/2} e^{-i\phi_{1p}} \\ v_p = \frac{1}{\sqrt{2}} \left( 1 - \frac{\epsilon_p}{E_p} \right)^{1/2} e^{-i\phi_{2p}} \end{cases}$$

$$\Delta_p = |\Delta_p| e^{-i(\phi_{1p} + \phi_{2p})}$$

$$E_p = \sqrt{\epsilon_p^2 + |\Delta_p|^2}$$

$$\Delta_p = \frac{1}{V} \sum_{p'} U(\vec{p} - \vec{p}') \langle c_{-p'\downarrow} c_{p'\uparrow} \rangle$$

$$\alpha_{p\sigma} |BCS\rangle = 0$$

- Solutions de BCS
    - Cohérence
    - Singulet, Triplet,  $s, p, d, f$ .
  - $T_c, \Delta_0$  equ. G.L.
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## Effet Josephson.

- Ham. de Josephson
- Base  $\theta, N$
- Courant, effet Josephson

$$\langle \alpha_{p\sigma}^\dagger \alpha_{p\sigma} \rangle = \frac{1}{e^{\beta E_p} + 1} \quad ?$$

$$Z = \text{Tr} e^{-\beta \sum_{p\sigma} E_p \alpha_{p\sigma}^\dagger \alpha_{p\sigma}}$$

$$\langle \sigma \rangle = \frac{\text{Tr} e^{-\beta H_{BCS}} \sigma}{\text{Tr} e^{-\beta H_{BCS}}}$$

$$\langle \alpha_{p\sigma}^\dagger \alpha_{-p\sigma}^\dagger \rangle = 0$$

$$\alpha_{p\sigma}^\dagger \alpha_{p\sigma}^\dagger |BCS\rangle$$

$$\alpha_{p\sigma} |BCS\rangle = 0$$

$$\langle c_{-p'\downarrow} c_{p'\uparrow} \rangle =$$

$$\langle (v_{p'} \alpha_{p'\uparrow}^+ + u_{p'} \alpha_{-p'\downarrow}^+) (u_{p'} \alpha_{p'\uparrow} - v_{p'} \alpha_{-p'\downarrow}^+) \rangle$$

$$= v_{p'} u_{p'} \langle (\alpha_{p'\uparrow}^+ \alpha_{p'\uparrow} - \alpha_{-p'\downarrow}^+ \alpha_{-p'\downarrow}^+) \rangle$$

$$= \frac{1}{2} \left(1 - \frac{v_{p'}^2}{E_{p'}^2}\right)^{1/2} e^{-i\phi_{1p'} - i\phi_{2p'}} (2n(E_{p'}) - 1)$$

$$= \frac{1}{2} \frac{\Delta_{p'}}{E_{p'}} (2n(E_{p'}) - 1)$$

$$\Delta_p = -\frac{1}{V} \sum_{p'} U(p-p') \frac{|\Delta_{p'}| e^{-i\phi_{1p'} - i\phi_{2p'}}}{E_{p'}} (1 - 2n(E_{p'}))$$

$$\Delta_p = -\frac{1}{V} \sum_{p'} U(p-p') \frac{\Delta_{p'}}{E_{p'}} (1 - 2n(E_{p'}))$$

Supra. de type s, d, p... cohérence.

$$\text{BCS } U(p-p') = -|U_0|$$

$$\Delta_p = |U_0| \int \frac{d^3 p'}{(2\pi)^3} \frac{\Delta_{p'}}{E_{p'}} (1 - 2n(E_{p'}))$$

$$\Delta = |U_0| \int \frac{d^3 p'}{(2\pi)^3} \frac{\Delta}{E_{p'}} (1 - 2n(E_{p'}))$$

$$\Delta = |\Delta| e^{i\phi} \text{ où } \phi \text{ est indep. de } \vec{p}$$

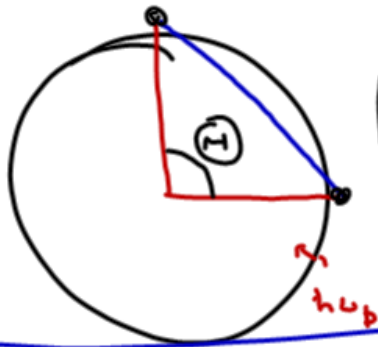
$$\prod_h \left( 1 + \frac{v_h}{u_h} c_{-h\downarrow}^+ c_{h\uparrow}^+ \right) |0\rangle$$

$$\uparrow e^{-i\phi_{2p} - i\phi_{1p}} = e^{i\phi}$$

$$\begin{aligned}
 \langle c_{-p'\downarrow} c_{p'\uparrow} \rangle &= \langle c_{p'\downarrow} c_{-p'\uparrow} \rangle \\
 &= - \langle c_{-p'\uparrow} c_{p'\downarrow} \rangle
 \end{aligned}$$

$\Delta$  spatial  
 Singulet de spin

$$U(P-P') = U(\cos \Theta)$$

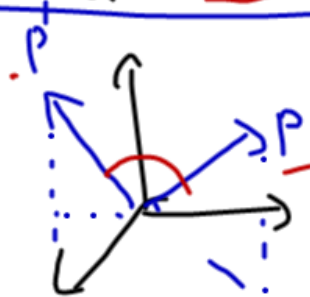


$$|P-P'| = \sqrt{P_F'^2 + P_F^2 - 2P_F P_F' \cos \Theta}$$

$$U(\cos \Theta) = \sum_{l=0}^{\infty} U_l P_l(\cos \Theta)$$

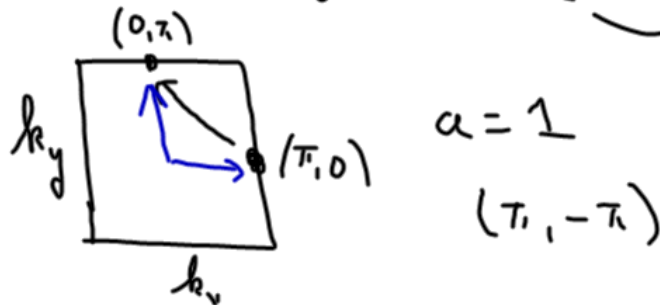
$$P_l(\cos \Theta) = \frac{4\pi}{2l+1} \sum_{m=-l}^l Y_{lm}(\theta, \varphi)$$

$Y_{lm}(\theta', \varphi')$



$$\Delta_{l,m}^{\checkmark} = - \int \frac{d^3 p}{(2\pi)^3} U_l \frac{4\pi}{2l+1} \frac{\Delta_{l,m}^{\checkmark}}{2E_{p'}} \quad (1 - 2n(E_{p'}))$$

$$U(\cos(\phi - \phi')) = U_0 - V \cos^2(\phi - \phi')$$



$$\cos^2 \theta = \frac{1 + \cos 2\theta}{2}$$

$$= U_0 - V \left( \frac{1 + \cos 2\phi \cos 2\phi' + \sin 2\phi \sin 2\phi'}{2} \right)$$

$$\Delta_P \rightarrow - \frac{1}{2} \int \frac{d^3 p'}{(2\pi)^3} U(r-r') \frac{\Delta_P'}{E_{p'}} \quad (1 - 2n(\bar{E}_{p'}))$$

Solution BCS,  $T_c$ ,  $\Delta_0$ , L.G.

$$\int \frac{d^3 p'}{(2\pi)^3} \Rightarrow \int d\Omega \int_0^\infty \frac{p'^2 dp'}{(2\pi)^3}$$

$$= 4\pi \int_0^\infty \frac{p'^2 dp'}{(2\pi)^3} \rightarrow \int_{-\hbar\omega_D}^{\hbar\omega_D} N(\xi') d\xi'$$

$$\Delta = |U_0| \int_{-\hbar\omega_D}^{\hbar\omega_D} N(\xi') \frac{\Delta}{E_{p'}} (1 - 2n(E_{p'})) d\xi'$$

$$= |U_0| D(E_F) \int_0^{\hbar\omega_D} \frac{\Delta}{E_{p'}} (1 - 2n(E_{p'})) d\xi'$$

$$D(E_F) = 2N(E_F)$$

$$\Delta = \frac{\Delta |U_0| D(E_F)}{2} \int_0^{\hbar\omega_D} d\xi' \left\{ \frac{\tanh\left(\frac{1}{2}\beta\sqrt{\xi'^2 + \Delta^2}\right)}{\sqrt{\xi'^2 + \Delta^2}} \right\}$$

Pics de  $T_c$

$$\Delta = a'' \Delta + b'' \Delta |\Delta|^2$$

$$(1 - a'') = b'' |\Delta|^2$$



$$1 = \frac{|U_0| D(E_F)}{2} \int_0^{\hbar\omega_D/\beta/2} \frac{\tanh(\beta_c S/2)}{\beta S/2} d\frac{\beta S}{2}$$

$$= \frac{|U_0| D(E_F)}{2} \left( \ln\left(\frac{\beta_c \hbar\omega_D}{2}\right) - \int_0^\infty \frac{\ln x}{\cosh^2 x} dx \right)$$

$\beta_c \hbar\omega_D \gg 1$

$$\boxed{\hbar\omega_D \gg T_c}$$

$$1 = \frac{|U_0| D(E_F)}{2} \left( \ln\left(\frac{\beta_c \hbar\omega_D}{2}\right) - \ln\left(\frac{4\gamma^4}{\pi}\right) \right)$$

$$\boxed{k_B T_c = \frac{2\gamma^4}{\pi} \hbar\omega_D e^{-2/|U_0| D(E_F)}}$$

à  $T=0$

$$D_0 = \Delta_0 \frac{|U_0| D(E_F)}{2} \int_0^{\hbar\omega_D} \frac{dS}{\sqrt{S^2 + \Delta_0^2}}$$

$$S = |\Delta_0| \sinh \theta$$

$$\frac{2}{|U_0| D(E_F)} = \sinh^{-1} \left( \frac{\hbar\omega_D}{|\Delta_0|} \right)$$

$$|\Delta_0| = \frac{\hbar\omega_D}{\sinh \left( \frac{2}{|U_0| D(E_F)} \right)}$$

$$\sinh x = \frac{e^x + e^{-x}}{2}$$

$$\boxed{|\Delta_0| = 2\hbar\omega_D e^{-2/|U_0| D(E_F)}}$$

$$\frac{2 \Delta_0}{k_B T_c} = 3.53$$

Al

$$3.37 \pm 0.1$$

Cd

$$3.2 \pm 0.1$$

Sn

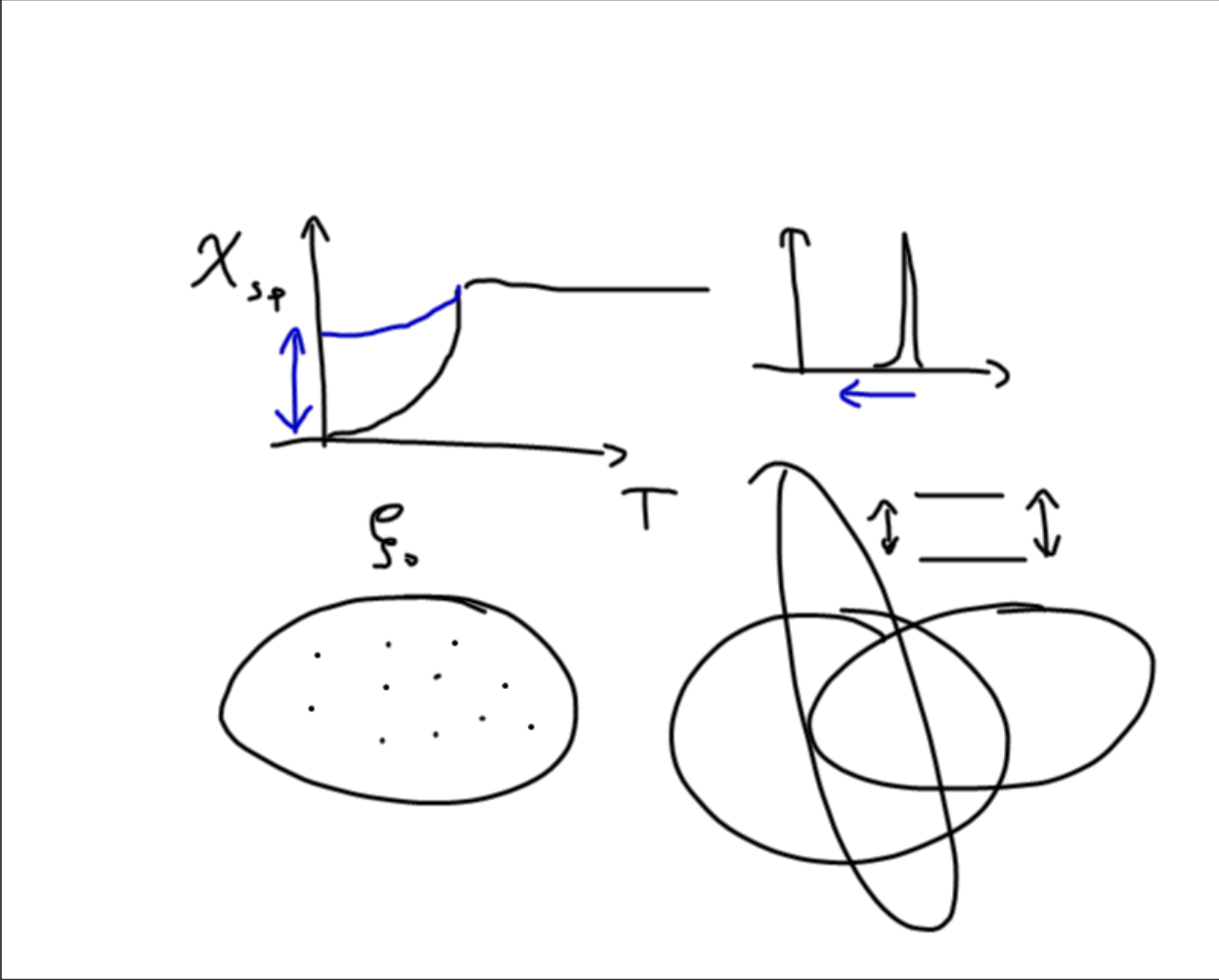
$$3.46 \pm 0.1$$

Pb

$$4.29 \pm .4$$

Hg

$$4.6$$



Josephson

$\phi_G$   $\phi_D$   $e^{-\frac{2eVt}{\hbar}}$

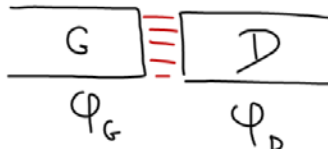
$$|BCS\rangle = \prod_k \left( 1 + \frac{v_k}{u_k} e^{-i(\phi_{2k} + \phi_k)} c_{-k\downarrow}^+ c_{k\uparrow}^+ \right) |0\rangle$$

$$\left( 1 + g(k) e^{-i\phi} c_{-k\downarrow}^+ c_{k\uparrow}^+ \right) |0\rangle$$

$$\Psi_\phi = \sum_{n=0}^N e^{-in\phi} \Psi_n$$

$$\Psi = \int_0^{2\pi} \frac{d\phi}{2\pi} e^{im\phi} \Psi_\phi$$

$$\int_0^{2\pi} \frac{d\phi}{2\pi} e^{i(p-m)\phi} = \delta_{p,m}$$



Base in tunnel = 0

$$\Psi_{\phi_G} \Psi_{\phi_D} = \sum_{n_G=0}^{N_G} e^{-in_G \phi_G} \Psi_{n_G} \times \sum_{n_D=0}^{N_D} e^{-in_D \phi_D} \Psi'_{n_D}$$

$$n_G = \frac{N_T}{2} + n \quad n_D = \frac{N_T}{2} - n$$

$$= \sum_{N_T=0}^{N_G+N_D} \sum_{n=-\frac{N_T}{2}}^{\frac{N_T}{2}} e^{-i(\frac{N_T}{2}+n)\phi_G} e^{-i(\frac{N_T}{2}-n)\phi_D}$$

$$= \sum_{N_T=0}^{N_G+N_D} e^{-i\frac{N_T}{2}(\phi_G + \phi_D)} \sum_{n=-\frac{N_T}{2}}^{\frac{N_T}{2}} e^{-in(\phi_G - \phi_D)} \Psi_{\frac{N_T}{2}+n} \Psi'_{\frac{N_T}{2}-n}$$

$$|\theta\rangle = \sum_{n=-\infty}^{\infty} e^{-in\theta} |n\rangle$$

Base pour calcul

$$\hat{H}_J = -\frac{E_J}{2} \sum_{n=-\infty}^{\infty} (|n\rangle\langle n+1| + |n+1\rangle\langle n|)$$

$$|n\rangle = \int_0^{2\pi} e^{in\theta} |\theta\rangle \frac{d\theta}{2\pi}$$

$$\sum_{n=-\infty}^{\infty} e^{-in\theta} = 2\pi \delta(\theta)$$

$$\hat{H}_J = -\frac{E_J}{2} \sum_{n=-\infty}^{\infty} \int_0^{2\pi} \int_0^{2\pi} e^{in\theta} e^{-i(n+1)\theta'} |\theta\rangle \frac{d\theta}{2\pi} \langle\theta'| e^{-i(n+1)\theta'} \frac{d\theta'}{2\pi} + \text{c.h.}$$

$$= -\frac{E_J}{2} \int_0^{2\pi} \frac{d\theta}{2\pi} \left( e^{-i\theta} |\theta\rangle \langle\theta| + \text{c.h.} \right)$$

$$\hat{H}_J = -E_J \int_0^{2\pi} \frac{d\theta}{2\pi} \cos\theta |\theta\rangle \langle\theta|$$

$$\hat{H}_J |\theta'\rangle = -E_J \cos\theta' |\theta'\rangle$$

Base  $\theta, N$  :



$$e^{-i\hat{\theta}} = \int_0^{2\pi} e^{-i\theta} |\theta\rangle \langle \theta| \frac{d\theta}{2\pi}$$

$$\hat{H}_J = -E_J \cos \hat{\theta}$$

$$\hat{N} = \sum_{n=-\infty}^{\infty} n |n\rangle \langle n|$$

$$\hat{N} |n'\rangle = n' |n'\rangle$$

$$e^{-i\hat{\theta}} |\theta'\rangle = e^{-i\theta'} |\theta'\rangle$$

$$\begin{aligned} e^{\pm i\hat{\theta}} |n\rangle &= e^{\pm i\hat{\theta}} \int_0^{2\pi} \frac{d\theta}{2\pi} e^{in\theta} |\theta\rangle \\ &= \int_0^{2\pi} \frac{d\theta}{2\pi} e^{i(n\pm 1)\theta} |\theta\rangle \end{aligned}$$

$$\hat{N} e^{\pm i\hat{\theta}} |n\rangle = (n\pm 1) |n\pm 1\rangle$$

$$= (n\pm 1) e^{\pm i\hat{\theta}} |n\rangle$$

$$e^{\pm i\hat{\theta}} \hat{N} |n\rangle = n |n\pm 1\rangle$$

$$= n e^{\pm i\hat{\theta}} |n\rangle$$

$$\boxed{[\hat{N}, e^{\pm i\hat{\theta}}] = \pm e^{\pm i\hat{\theta}}}$$

$$[\hat{N}, \hat{\theta}] = -i$$

## Courant et effet Josephson

$$\frac{d\hat{N}}{dt} = \frac{i}{\hbar} [\hat{H}, \hat{N}]$$

$$= \frac{i}{\hbar} \left( -\frac{E_J}{2i} \right) \left[ e^{i\hat{\theta}} + e^{-i\hat{\theta}}, \hat{N} \right]$$

$$= \frac{1}{\hbar} \left( +\frac{E_J}{2i} \right) \left[ -e^{i\hat{\theta}} + e^{-i\hat{\theta}} \right]$$

$$= -\frac{E_J}{\hbar} \sin \hat{\theta}$$

$$\hat{I} = -2e \frac{d\hat{N}}{dt} = \frac{2e}{\hbar} E_J \sin \hat{\theta}$$

Si on applique une diff. de potentiel  $V = V_D - V_G$

$$\hat{H} = (-2eV)(-\hat{N}) + \hat{H}_J$$

$$\frac{d\hat{\theta}}{dt} = \frac{i}{\hbar} [\hat{H}, \hat{\theta}]$$

$$= \frac{i}{\hbar} 2eV [\hat{N}, \hat{\theta}] = \frac{2e}{\hbar} V$$

$$\frac{d\hat{\theta}}{dt} = \frac{2eV}{\hbar} = \frac{2e}{\hbar} V \approx$$



$$\frac{d c_{p,\sigma}^+}{dt} = \frac{i}{\hbar} \left[ H_1, c_{p,\sigma}^+ \right]$$

$$-eV N_{ptcl} = -eV \sum_{p,\sigma} c_{p,\sigma}^+ c_{p,\sigma}$$

$$= -eV \frac{i}{\hbar} c_{p,\sigma}^+$$

$$c_{p,\sigma}^+(t) = c_{p,\sigma}^+(0) e^{-\frac{i e V}{\hbar} t}$$

$$\frac{\hbar}{i} \nabla \rightarrow \frac{\hbar}{i} \nabla - eA$$

$$i\hbar \frac{\partial}{\partial t} \rightarrow i\hbar \frac{\partial}{\partial t} + eV$$

$$|\psi| e^{i\phi}$$

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