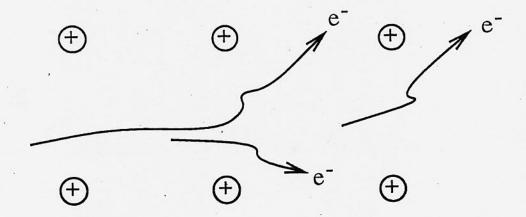
The Search for Unconventional Superconductors: A Long Story with a Happy End

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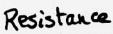
Basis of Condensed Matter Physics

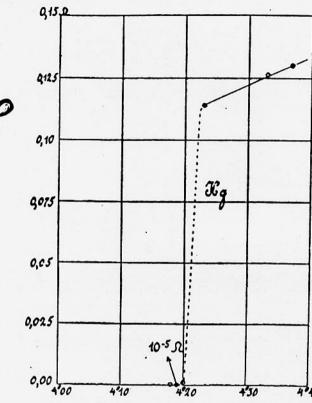
- Quantum Mechanics
- Electrons + Ions + Coulomb Interaction



• 1 cm 3 contains $\sim 10^{22}$ electrons + ions!

. Discovery of Superconductivity





· Hg.

Figure 1 Resistance in ohms of a specimen of mercury versus absolute temperature. This plot by Kamerlingh Onnes marked the discovery of superconductivity.

Temperature



Wolfgang Pauli



Felix Bloch

"Pauli thought that superconductivity was the only remaining matter of some interest in the theory of metals and that I should get on with it so as to be finally done with all these 'dirt effects'."

"Once in a while I thought that I had indeed found such states but it never took Pauli long to point to some error in the calculations. While he did not object to my approach he became rather annoyed at my continued failure to come out with the desired answer to such a simple question."

"After the fog, which so long enveloped the phenomenon, had begun to lift many years later, I could not resist reminding Pauli that the problem was not quite as easy to solve as he thought when he gave it to me. Since that time he had become more mellow – so much more, in fact, that he agreed."

F. Bloch, Proc. Roy. Soc. A371 (1980)

- Starting point: Normal Metal with repulsive (Coulomb) and attractive (electron-phonon-electron) forces
- Low Energy Scale: repulsive forces screened attractive forces enhanced
- leads to formation of bound electron pairs in a spin singlet
- below a transition temperature T_c there is a macroscopic coherent occupation of pairs to form $\Psi(r)$.
- Explains all features of classical superconductor

Can <u>Superconductivity</u> be reached thru' electron-electron interactions?

Yes!: Kohn-Luttinger '65

Landau Fermi Liquid

Sharp Fermi Surface causes attractive interactions in higher (ℓ >0) angular momentum channels, λ_{ℓ} <0 for pairing

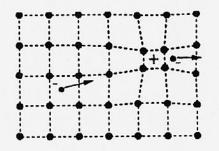
$$T_c \approx E_F e^{-1/|\lambda_\ell|}$$
 but $|\lambda_\ell| << 1$ and T_c is very small

T

Unconventional superconductivity

Cooper Pairing, if the electrons try to avoid each other

Conventional pairing attractive interaction by electron-phonon interaction





angular momentum I =0 spin singlet

Bardeen Cooper Schrieffer '57

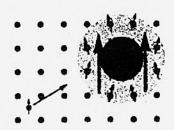
Cooper pairing from purely repulsive interactions: Kohn & Luttinger '65

I > O higher angular momentum; extremely small T_c

"unconventional pairing"

Pairing interaction mediated by spin fluctuations: Berk & Schrieffer '66, ---

most polarizable near magnetic instability



i > 0: avoiding Coulomb repulsion

Symmetry of Cooper Pairs

Pair wavefunction:

$$F_{ss'}(\vec{k}) = \langle \hat{c}_{\vec{k}s} \hat{c}_{-\vec{k}s'} \rangle = \Phi(\vec{k}) \chi (s, s')$$
orbital spin

totally antisymmetric under electron exchange

$$\vec{k} \rightarrow -\vec{k}$$
 $s \leftrightarrow s'$

even parity
$$\Phi(-\vec{k}) = \Phi(\vec{k})$$
 \longrightarrow S=0 singlet odd parity $\Phi(-\vec{k}) = -\Phi(\vec{k})$ \longrightarrow S=1 triplet L=1,3,5,...

Superconductivity of strongly correlated electrons

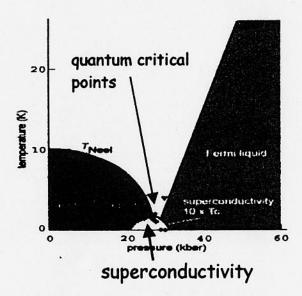
Heavy Fermion systems:

CeCu₂Si₂ Steglich '79

UBe₁₃ OH '83

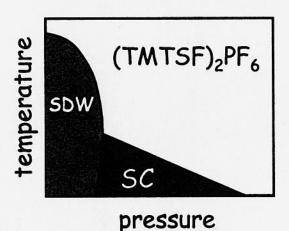
UPt₃ Stewart '84

CeIn₃ Mathur '98

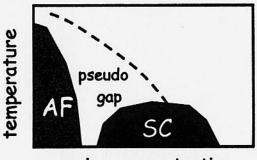


Relation to magnetism

Organic conductors: Bechgard, Jerome '80



High-Tc superconductors: Müller & Bednorz '86



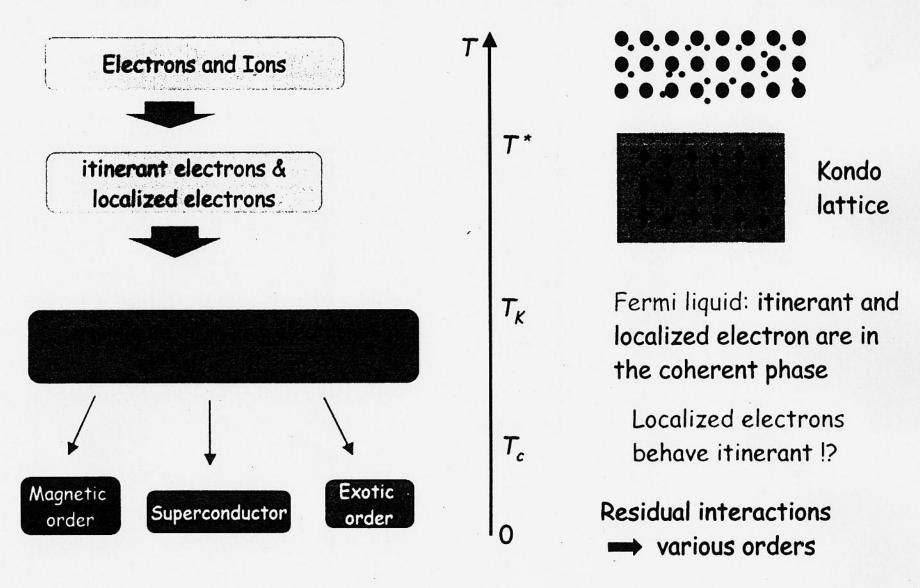
carrier concentration

cuprates YBa₂Cu₃O₇ La_{2-x}Sr_xCuO₄

 $T_{cmax} = 133K$

Heavy Fermion materials

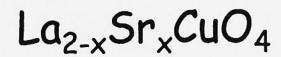
Intermetallic compounds with rare earth elements (partially filled f-shells)

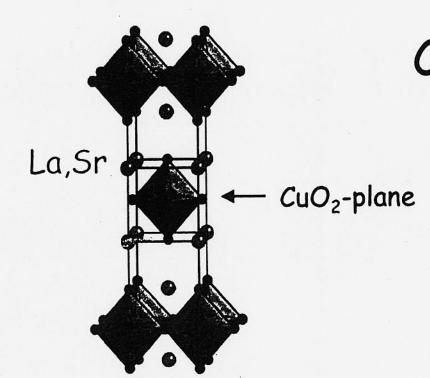


High-temperature superconductors

K.A. Müller & G. Bednorz 1

1986





Layered perovskite structure

Electronic configuration

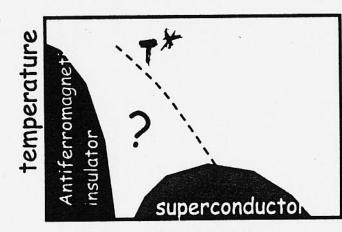


one 3d-eq hole

t₂₉

spin S = 1/2

magnetic Mott insulator -> superconductor

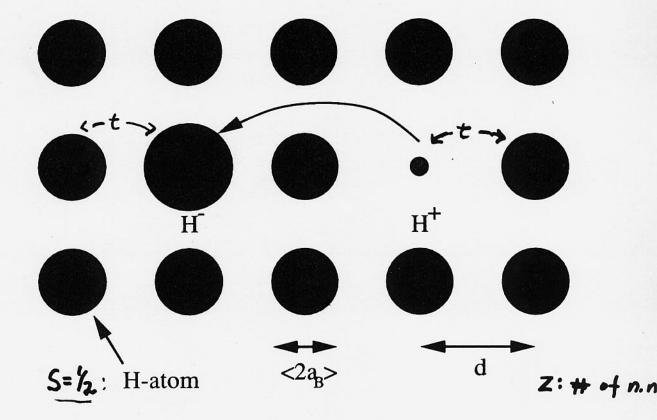


 T_c up to 133K

Carrier concentration x

Dilute Limit $(d >> a_B)$

⇒ Localized H-atoms.

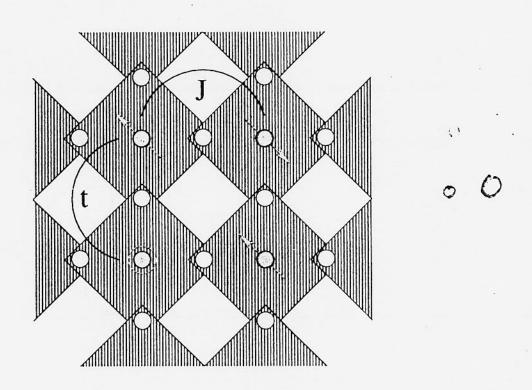


- \Rightarrow Charge Excit. Energy: U 2zt > 0. $[U \rightarrow E(H^- H^+) \sim 0.95 \text{Ry}, d \rightarrow \infty]$.
- → Mott Insulator with gap for charge excitations.
- \Rightarrow Low Energy Sector purely spin, $H_{Heis} = J \sum_{\langle ij \rangle} \mathbf{S}_i \cdot \mathbf{S}_j \Rightarrow \mathsf{AF}$ order.
- Local Gauge Invariance in Real Space (Kohn '64): $\Phi = \prod_i c_{i,\sigma_i}^{\dagger} | \text{vac} \rangle, \qquad c_{i,\sigma_i}^{\dagger} \to \mathrm{e}^{i\phi_i} c_{i,\sigma_i}^{\dagger}.$

Hole Doping introduces Cust

t-J Model

• restrict Hilbert space to Cu^{2+} and Cu^{3+} oxidation states.



ullet overlapping CuO_4 -squares lead to processes:

$$Cu^{2+} \leftrightarrow Cu^{2+}$$
, AF Heisenberg coupling, $J\vec{S}_{i}.\vec{S}_{j}$
 $Cu^{3+} \leftrightarrow Cu^{2+}$, hopping of electrons, t. $t = \frac{J}{3}$
 $Cu^{2+}: S = \frac{J}{2}$ $u^{3+}: S = 0$ [3d8: Zhang - Rice Singlet]

Numerical Simulation of the t-J Model

severely limits Quantum Monte Carlo method

Results:

- Superconducting state must have d_{x²-y²} pairing
- Competing ground states upon doping
 - \rightarrow d_{x2-y2} superconductivity
 - → stripe-phase with charge and spin order
 - → orbital antiferromagnetism or flux phase

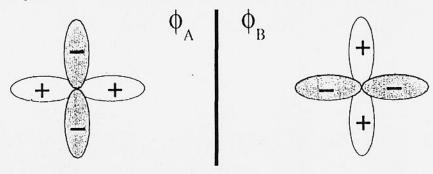
Questions:

- which is the ground state of the t-J model?
- which additional terms will favor which ground state?

Internal Symmetry of Cooper Pairs

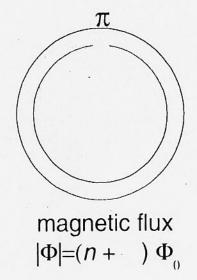
Possible sign changes in pairing amplitude around the Fermi surface tested

by π -junctions



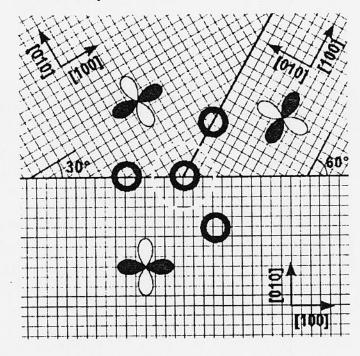
Energy in barrier: $\Delta F \propto -\cos(\phi_A - \phi_B + \pi)$

 ΔF causes a phase slip of π across the barrier



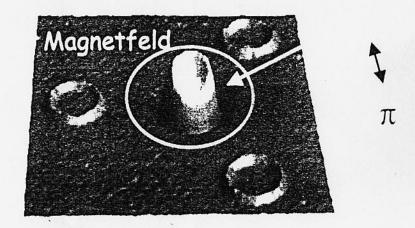
- \rightarrow Frustration in a loop with an odd number of $\pi-$ junctions (Geshkenbein, Larkin, Barone '87)
- → First signs in random ceramic samples, paramagnetic Meissner effect (Braunisch et al '92), d_{x²-y²} symmetry (Sigrist, Rice '92).
- → Controlled geometries (Wollmann et al. '93, Brawner, Ott '94, Mathai et al. '95, Tsuei, Kirtley '95)

Tri-crystal-configuration



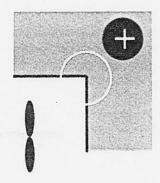
Tsuei, Kirtley et al. (IBM)

Superconducting loops $Ø = 60 \mu m$ $YBa_2Cu_3O_7$ $T_c = 92 K$ spontaneous currents and magnetic field

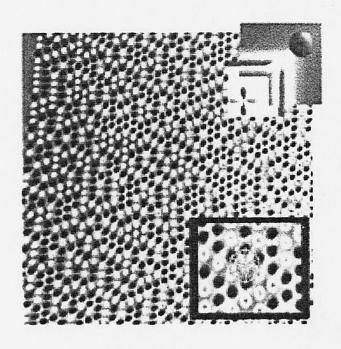


SQUID scanning microscope

Basic unit

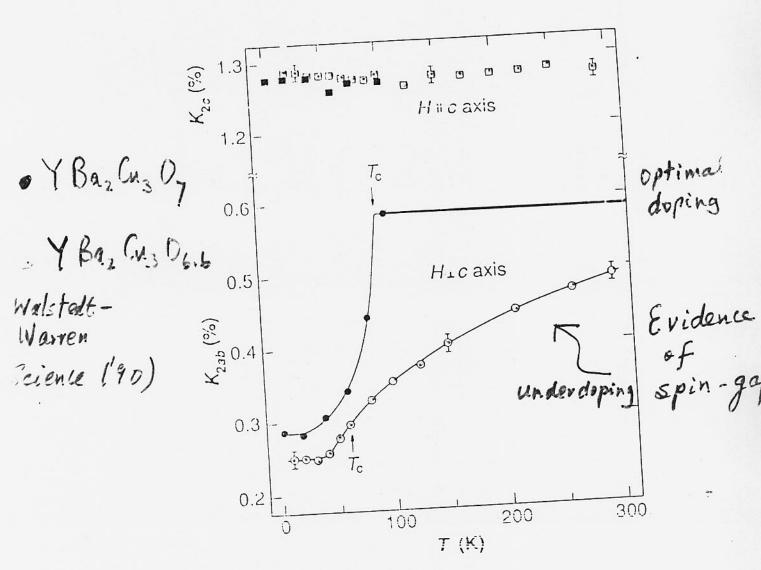


flux at the corner junction



H. Hilgenkamp et al., Nature 422, 50 (2003)

Cu-Knight Shifts K x X(T): Spin susceptibil



continuous onset of pairing.

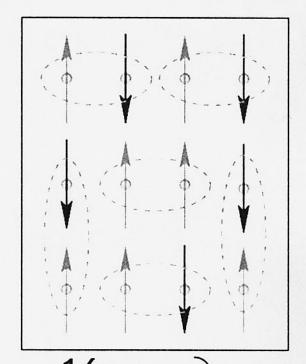
of Cu2+ spins at T>>> Tc

in the underdoped regime

Quantum Fluctuations can lead to a Spin Liquid of Singlet pairs

• Energy of a Singlet: $-\frac{3}{4}J$

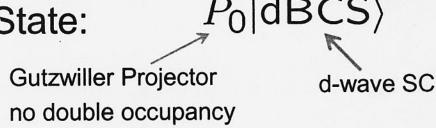
very favorable



But AF order has a lower energy at ½-filling in 2-Dimensions

Will hole doping stabilize RVB and make singlet pairs mobile leading to superconductivity?

dRVB Variational State:



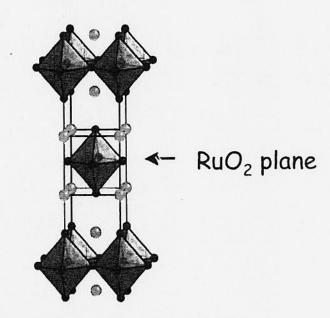
- Low Energy States of t-J model
 - $d_{x^2-y^2}$ pairing symmetry
- Density of superconducting carrier $\sim x$: hole density
 - * $T_c \sim x$ for x small due to nodal quasiparticles
- Spin pairing at higher temperatures $T^* > T$
- Explains main features of ARPES experiments

See: P.W. Anderson, P.A. Lee, M. Randeria, T.M. Rice, N. Trivedi and F.C. Zhang, to be published.

Sr2RuO, - a spin triplet superconductor

There is no better characterized unconventional superconductor - except may be high-T_c superconductors.

Transition metal oxide

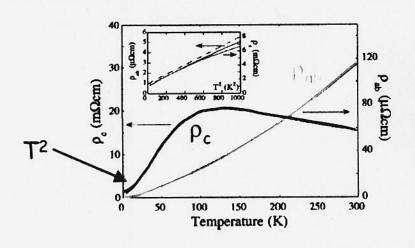


Maeno, Bednorz et al. ('94)

Superconductor with $T=1.5~\mathrm{K}$

Quasi-two-dimensional Fermi liquid

Strong correlation effects



Analogy to Fermi liquid ³He

Spin triplet superconductivity

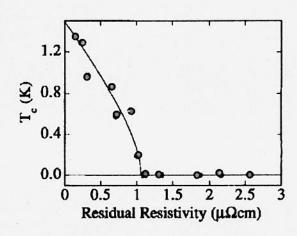
Superconductivity in Sr2RuO4

- strongly correlated 2D Fermi liquid
- superconducting with $T_c = 1.5 \text{ K}$

Maeno et al.,

Nature 372, 532 (1994)

T_c highly sensitive to non-magnetic impurities



Mackenzie et al., Phys. Rev. Lett. 80, 161 (1998)



unconventional superconductivity

pairing in higher-angular momentum channel

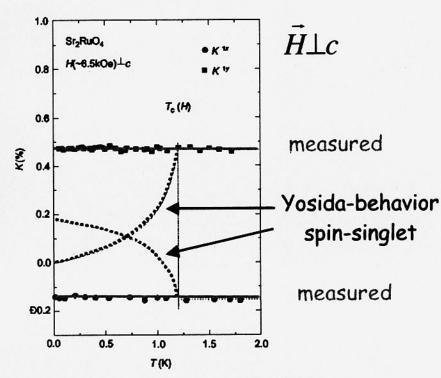


electronic analog to ³He? pairing in the p-wave spin-triplet channel (A- or B-phase?)

Two crucial experiments for Sr₂RuO₄

Spin susceptibility

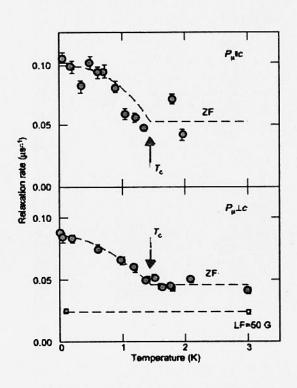
¹⁷O-NMR Knight shift



Ishida et al., Nature 396, 242 (1998)

inplane equal-spin pairing

Muon-spin relaxation



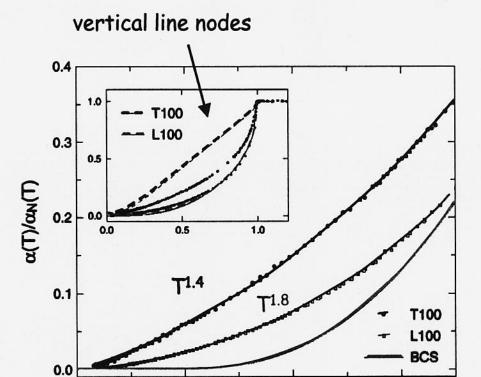
Luke et al., Nature 394, 558 (1998)

intrinsic magnetism T-violation

triplet A-phase: $\Delta(\vec{k}) = (i\sigma'\sigma^2) \cdot d_2(\vec{k})$ with $\vec{d} = \hat{z}(k_x \pm ik_y)$

Ultrasound absorption

0.6



Lupien, Taillefer et al.

0.2

0.0

Real powerlaw or multi-band effect?

0.4

T/T_c

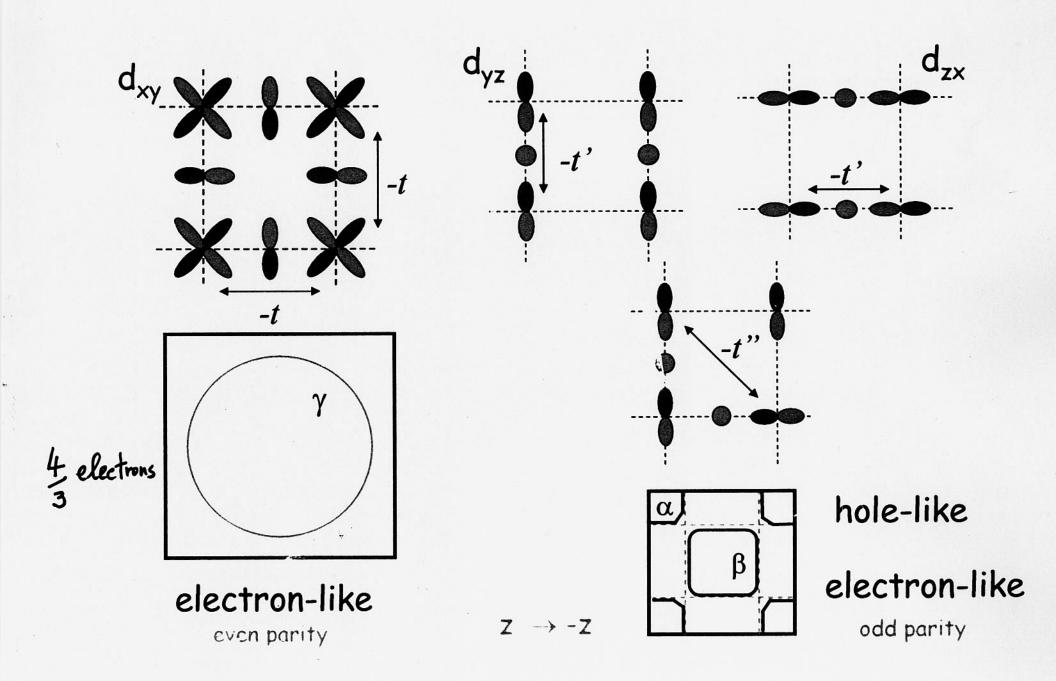
Propagation direction and polarization of ultrasound is important!

Line nodes:

$$\alpha(T) \longrightarrow \begin{cases} T^{1.5} & \text{active nodes} \\ T^{3.5} & \text{inactive nodes} \end{cases}$$

Moreno & Coleman

Electronic structure



Microscopic Model for St. Ru Dy

· Y-band active band

* Agterberg, Rice Signist
for superconductivity

* Zhitomirsky-Rice

. d.B bands passive bands

8-band: - 2-Dim. (xy-orbital) and 4/3 el. / Ru

- Fermi Surface very close to saddle-pits.

(van Hove singalanities

The Superconductor that stood on the shelf.

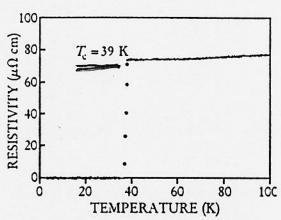


FIGURE 2. SUPERCONDUCTIVITY in magnesium diboride appears at a temperature of 39 K. (Adapted from ref. 1).

Nagamatsu
Akimitsu
Nature Mar. 101

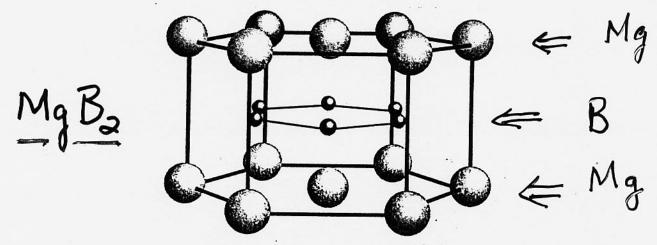


FIGURE 1. MAGNESIUM DIBORIDE belongs to the AlB, family of structures. The magnesium atoms, shown here in gray, form a hexagonal layer, while the boron atoms, shown in brown, from a graphite-like honeycomb layer.

Fermi Surface in Tr- and T- Bands

· A BCS electron-phonon Superconductor

Endless search for ever more exotic species

