

Mott Physics in Strongly Correlated Superconductivity

A.-M. Tremblay

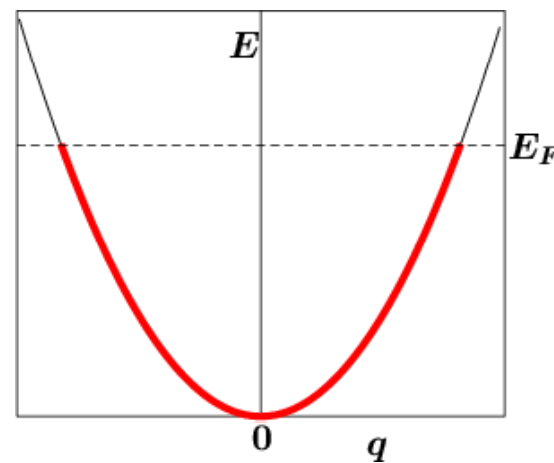
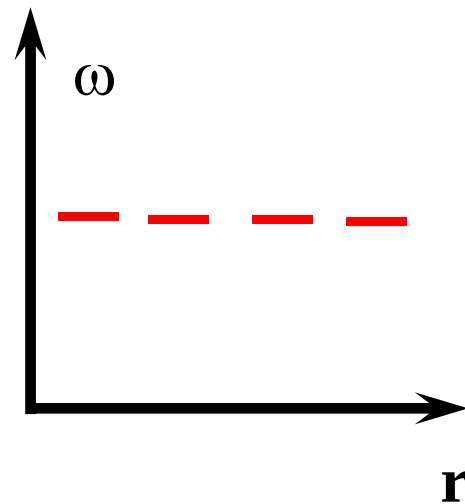
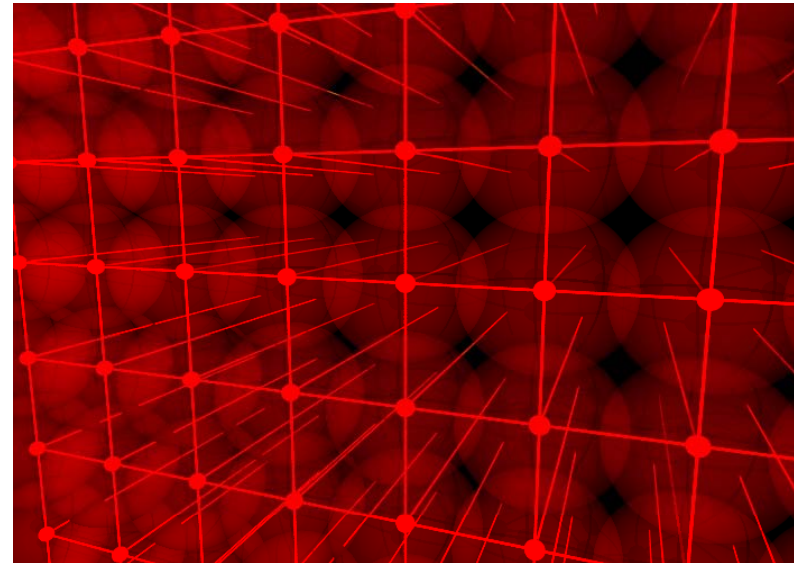
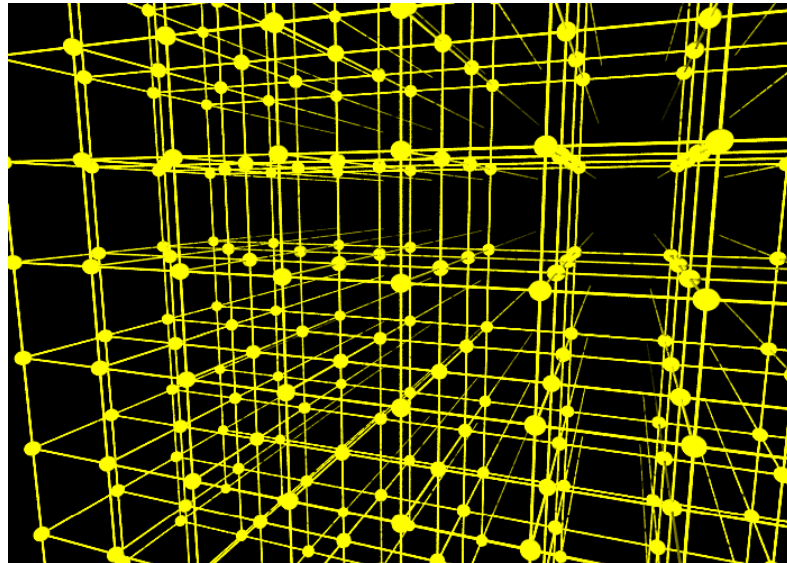
G. Sordi, D. Sénéchal, K. Haule,
S. Okamoto, B. Kyung, M. Civelli



Sousse, 6 November 2010

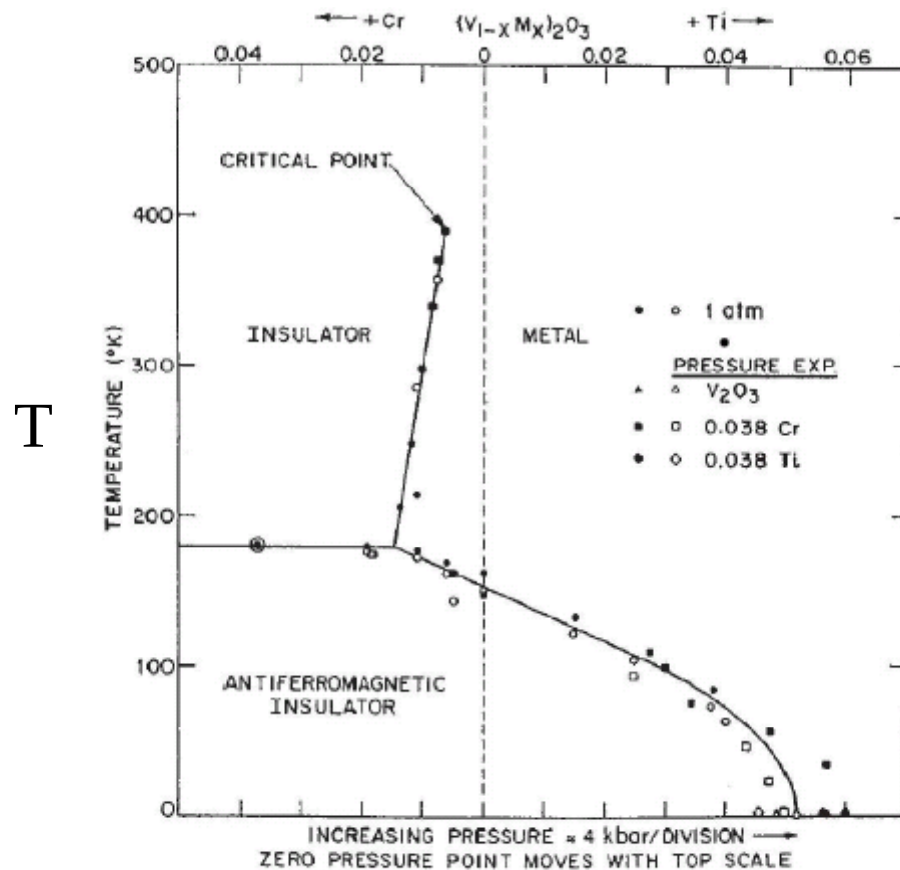


How to make a metal



Courtesy, S. Julian

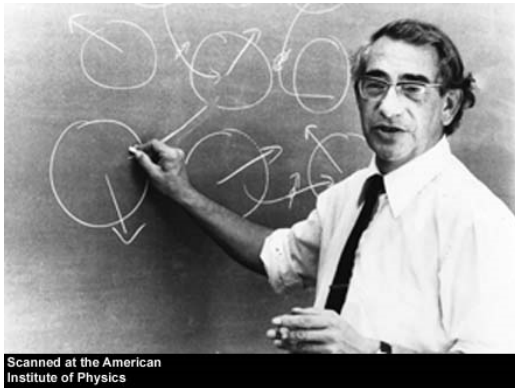
« Conventional » Mott transition



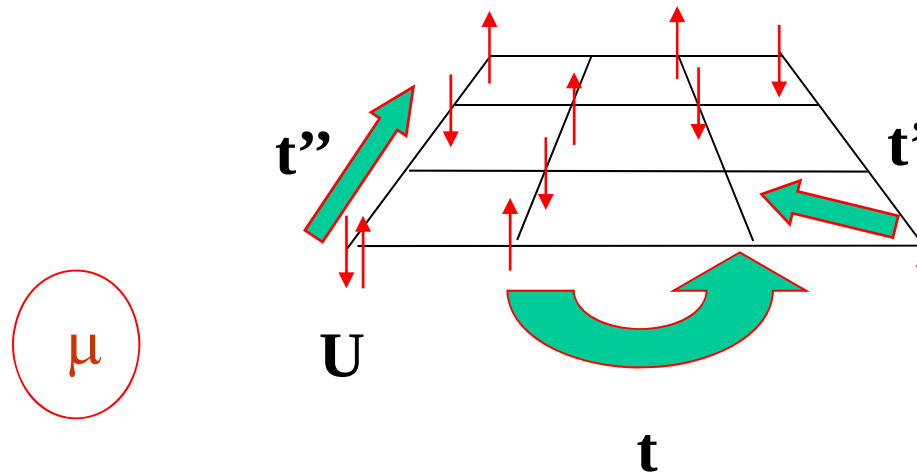
Understood from Hubbard model
and dynamical mean field theory

Figure: McWhan, PRB 1970; Limelette, Science 2003

Hubbard model

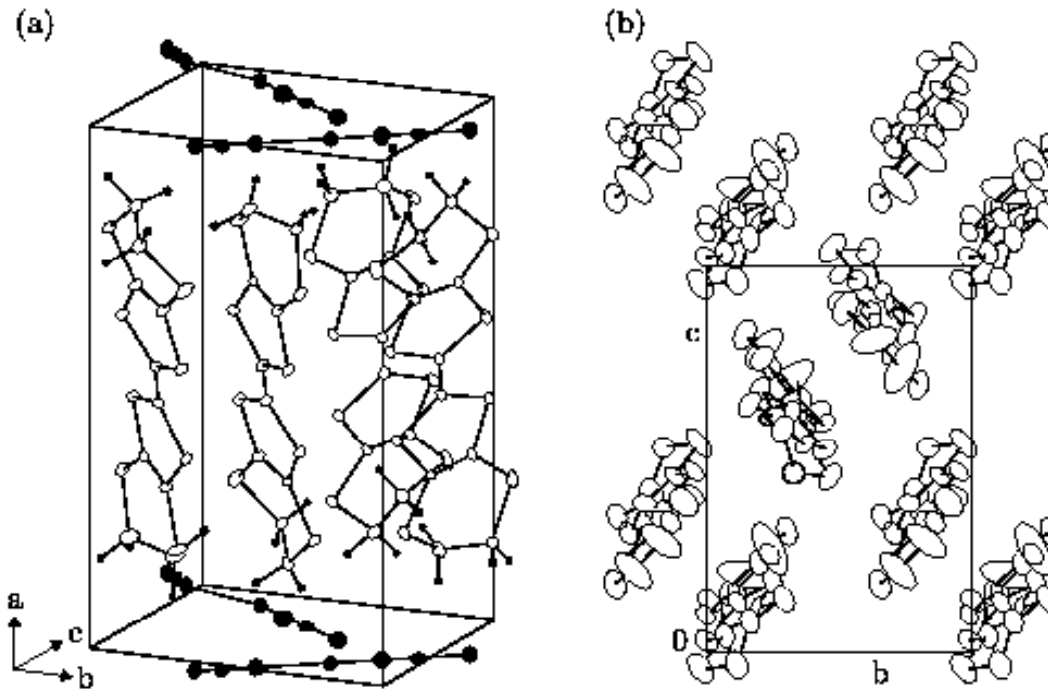


1931-1980



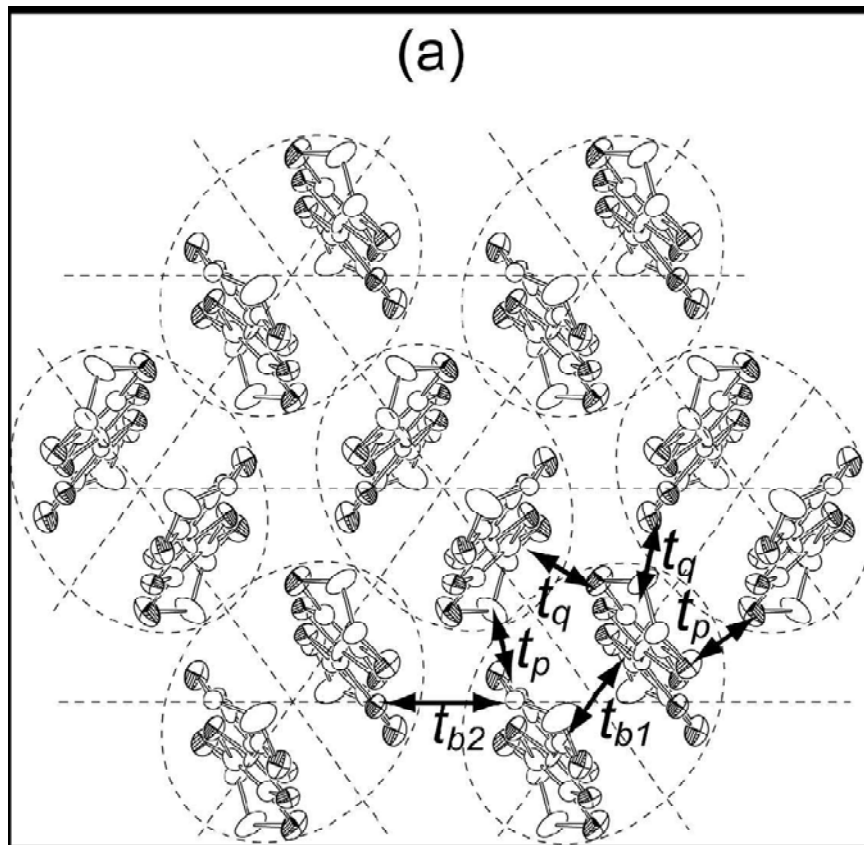
$$H = -\sum_{\langle ij \rangle \sigma} t_{i,j} (c_{i\sigma}^\dagger c_{j\sigma} + c_{j\sigma}^\dagger c_{i\sigma}) + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

Layered organic superconductors (κ -BEDT-X)

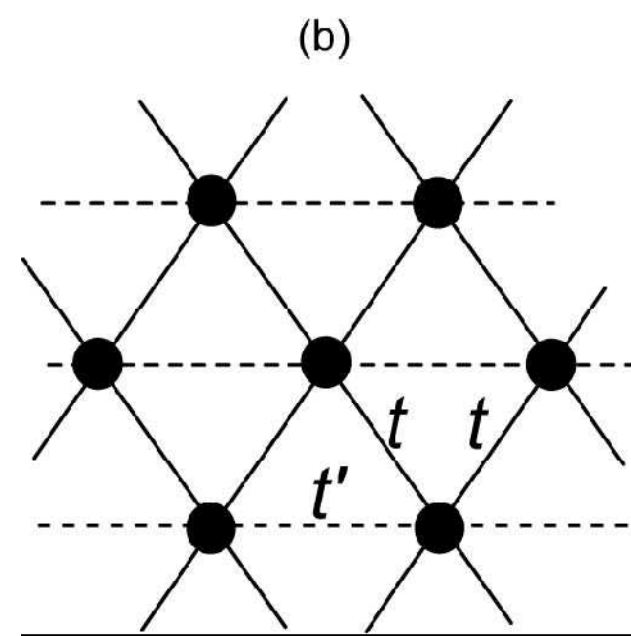


One-band Hubbard model of BEDT organics

H. Kino + H. Fukuyama, J. Phys. Soc. Jpn **65** 2158 (1996),
R.H. McKenzie, Comments Condens Mat Phys. **18**, 309 (1998)



Y. Shimizu, et al. Phys. Rev. Lett. **91**,
107001(2003)



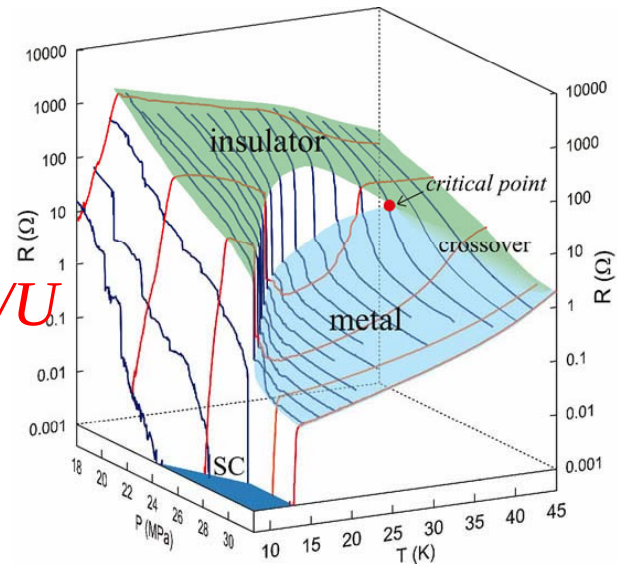
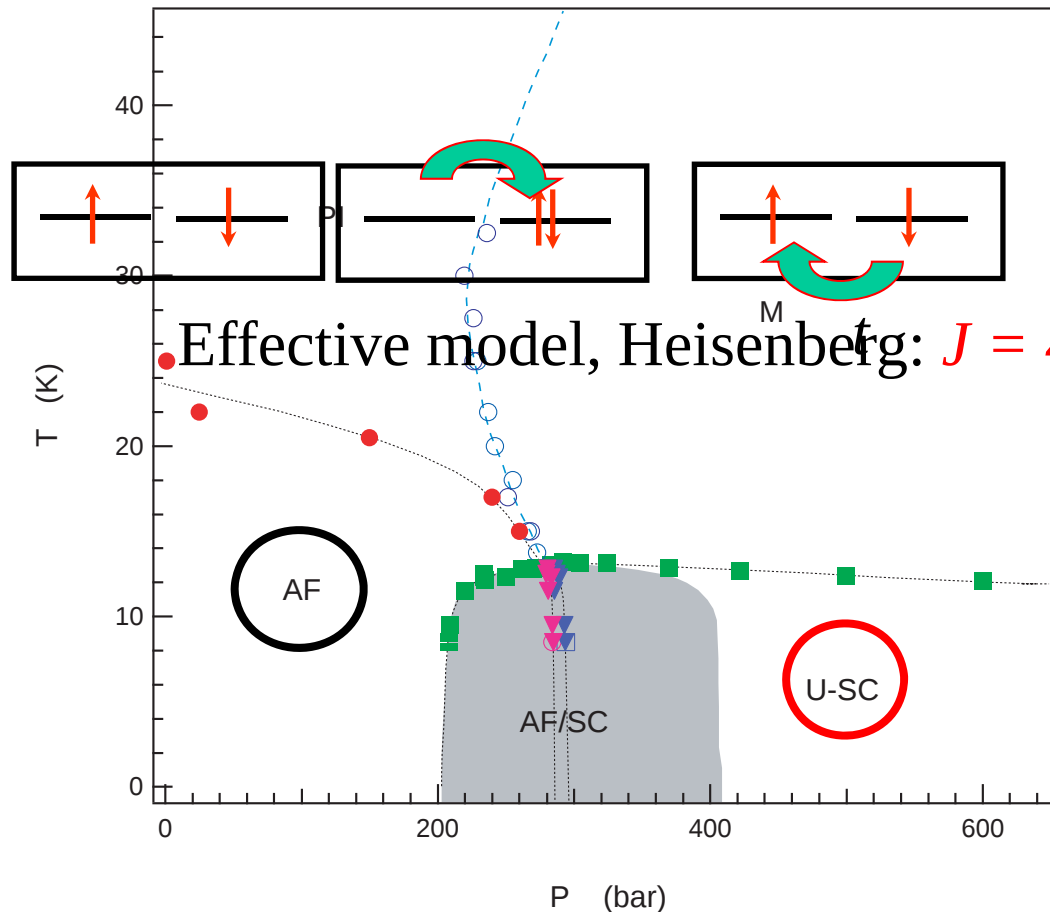
$$t \approx 50 \text{ meV}$$

$$\Rightarrow U \approx 400 \text{ meV}$$

$$t'/t \sim 0.6 - 1.1$$



Bare Mott critical point in organics

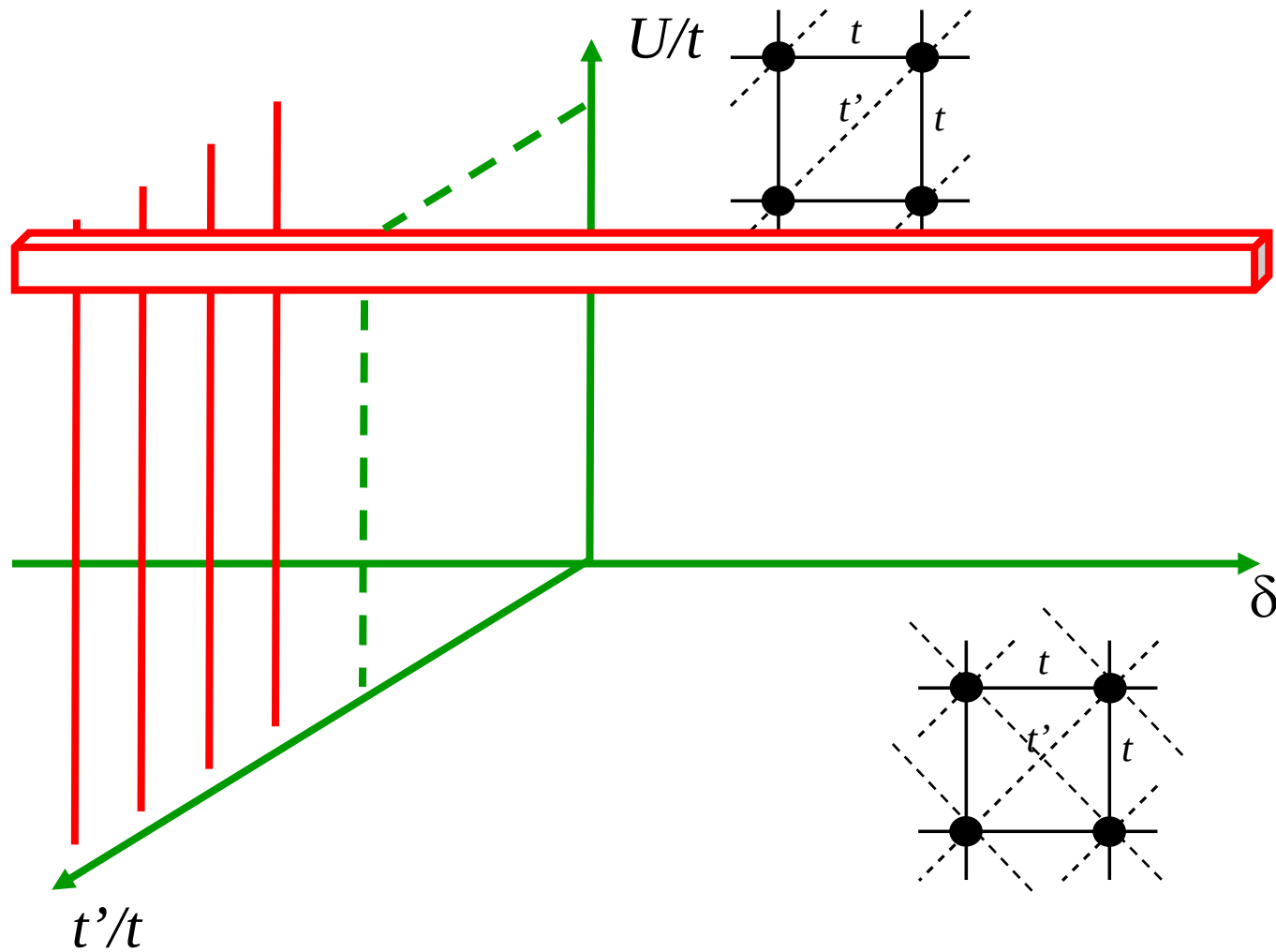


F. Kagawa, K. Miyagawa, + K. Kanoda
PRB **69** (2004) + Nature **436** (2005)

Phase diagram ($X=\text{Cu}[\text{N}(\text{CN})_2]\text{Cl}$)

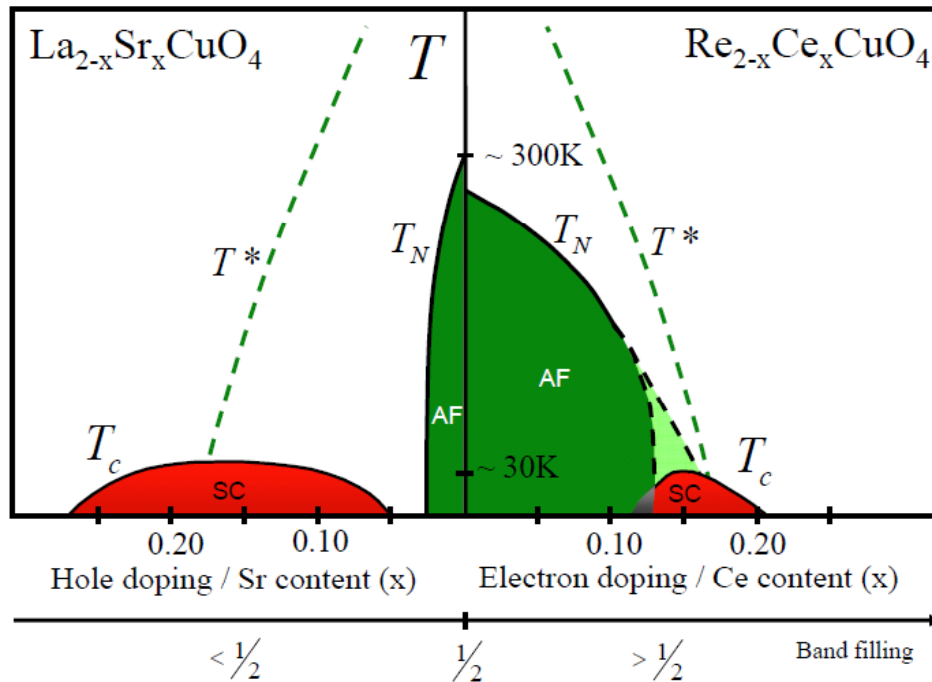
S. Lefebvre et al. PRL **85**, 5420 (2000), P. Limelette, et al. PRL **91** (2003)

Perspective



Normal state of high-temperature superconductors

Armitage, Fournier, Greene, RMP (2009)



What is under the dome?
Mott Physics away from $n = 1$

- Competing order

- Current loops: Varma, PRB **81**, 064515 (2010)
- Stripes or nematic: Kivelson et al. RMP 75 1201(2003); J.C.Davis
- d-density wave : Chakravarty, Nayak, Phys. Rev. B **63**, 094503 (2001); Affleck et al. flux phase
- SDW: Sachdev PRB **80**, 155129 (2009) ...

- Or Mott Physics?

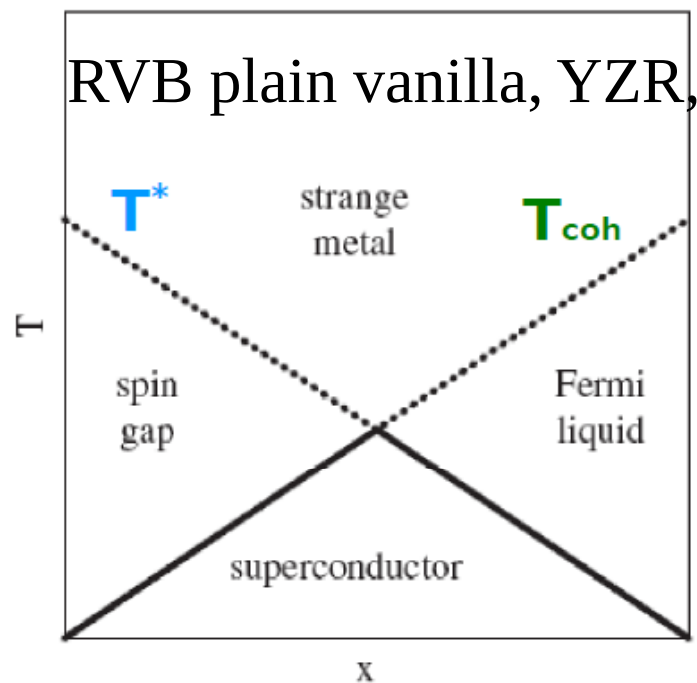
- RVB: P.A. Lee Rep. Prog. Phys. **71**, 012501 (2008)

Two views (caricature)

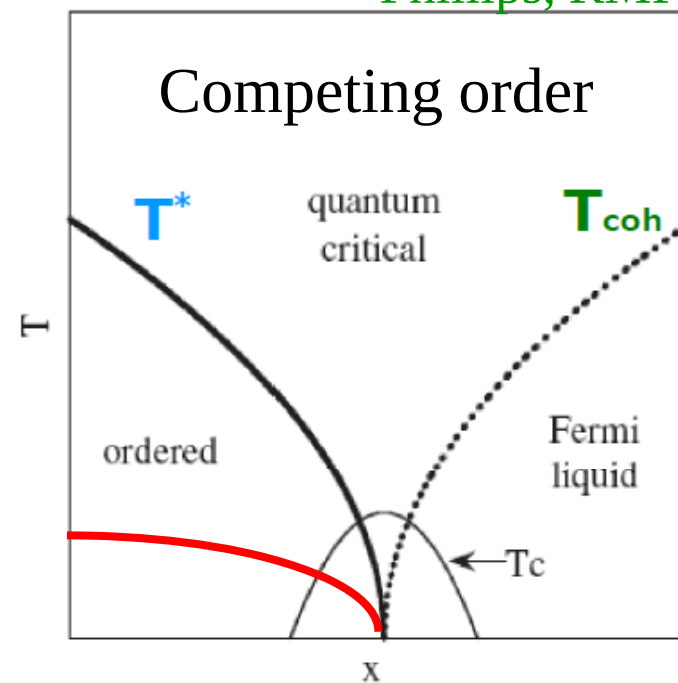
Norman, Adv. Phys. (2005)

Broun, Nat. Phys. (2006)

Phillips, RMP (2010)

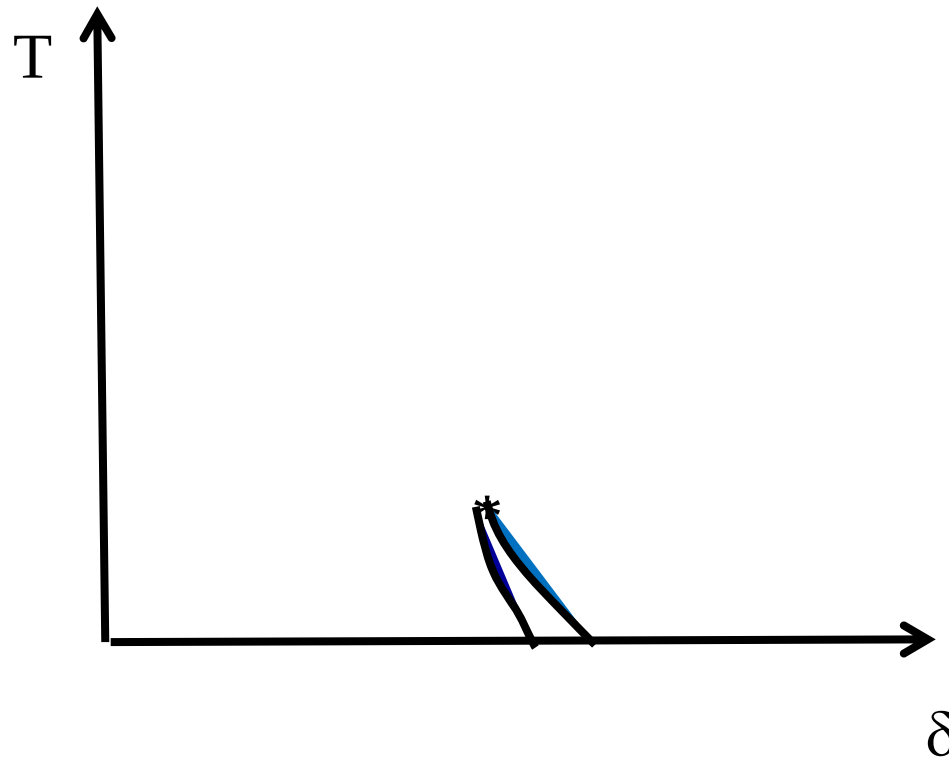


Why T_c decreases?
What is the origin of T^* ?
What is the strange metal?



Long correlation length or not.
What lies beneath the dome.
Mott Physics away from $n = 1$

An alternate view (a bit of both)



G. Sordi, K. Haule, A.-M.S.T
PRL, **104**, 226402 (2010)

Outline

- Method
- Normal state
- Superconducting state

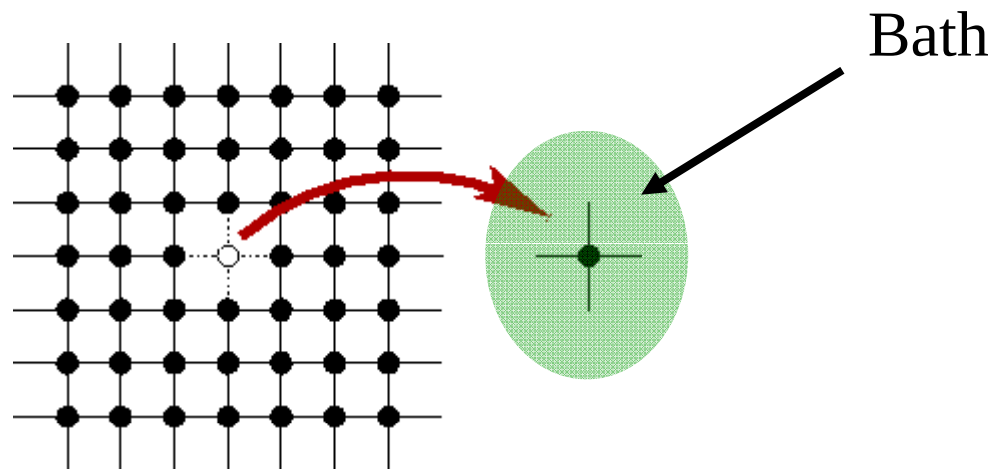
Method



Mott transition and Dynamical Mean-Field Theory.

The beginnings in $d = \text{infinity}$

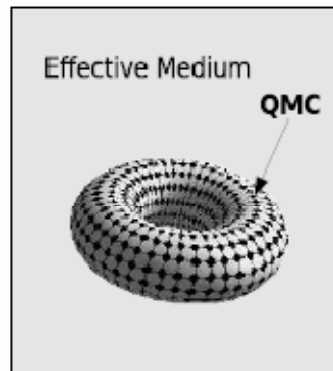
- Compute scattering rate (self-energy) of impurity problem.
- Use that self-energy (ω dependent) for lattice.
- Project lattice on single-site and adjust bath so that single-site DOS obtained both ways be equal.



W. Metzner and D. Vollhardt, PRL (1989)
A. Georges and G. Kotliar, PRB (1992)
M. Jarrell PRB (1992)

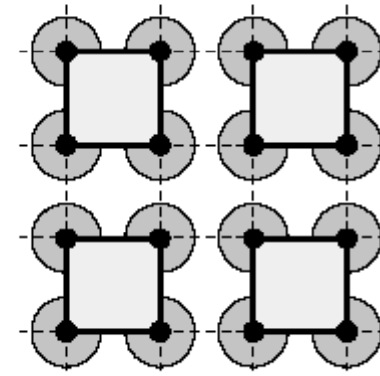
DMFT, ($d = 3$)

2d Hubbard: Quantum cluster method

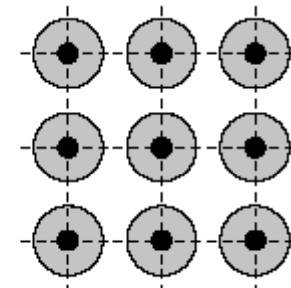


DCA

C-DMFT



DMFT



Hettler ...Jarrell...Krishnamurty PRB **58** (1998)
Kotliar et al. PRL **87** (2001)
M. Potthoff *et al.* PRL **91**, 206402 (2003).
Maier, Jarrell et al., Rev. Mod. Phys. **77**, 1027 (2005)

Understanding finite temperature phase from a *mean-field theory* down to $T = 0$

- Fermi liquid
 - Start from Fermi sea
 - Self-energy analytical
 - One to one correspondence of elementary excitations
 - Landau parameters
- Mott insulator
 - Hubbard model
 - Atomic limit
 - Self-energy singular
 - DMFT
 - How many sites in the cluster determines how low in temperature your description of the normal state is valid.

Mott insulator at finite T

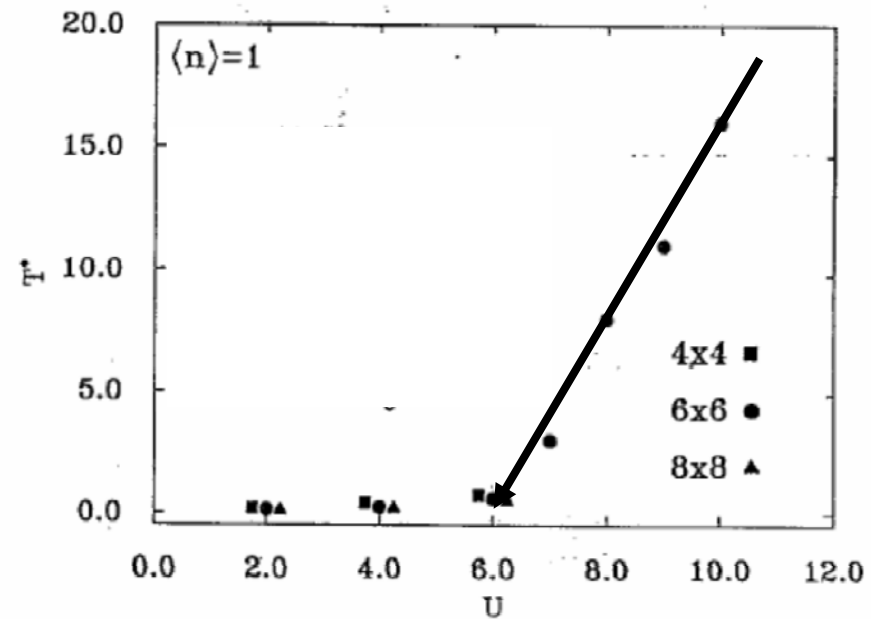
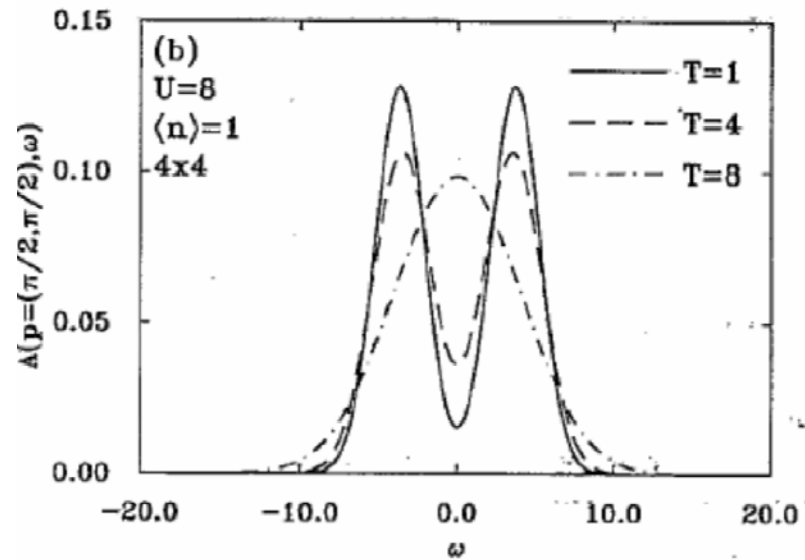


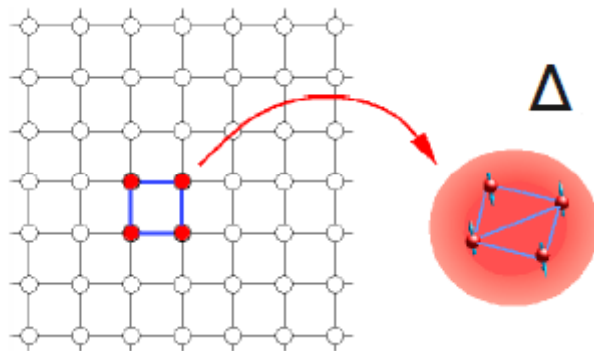
FIG. 5. The temperature T^* at which the gap develops vs U for 4×4 , 6×6 , and 8×8 lattices.

M. Vekic and S.R. White, PRB 47, 1160 (1993)

Interaction-induced Mott transition, $n = 1$

Method	U_{c1}	U_c	U_{c2}	Ref.
VCA+ED 2 x 2 + 8b	5.25	5.5	6.37	Balzer et al. EPL (2009)
CDMFT+CTQMC+H 2 x 2	5.3		5.7	Park et al. PRL (2008)
DCA+CTQMC+H 8	5.7		6.4	Gull et al. cond-mat (2009)
DCA+CTQMC+H 4	!	~4.2	!	Gull et al. EPL (2008)
Dual fermions	!	~6.5	!	Hafermann et al. (2008)
CDMFT+ED 2 x 2 + 8b 15 parameters	?	~5.6	?	Liebsch, Merino... (2008)
CDMFT+ED 2,3,4		~4		Zhang et al. PRB (2007) (3d also)
QMC 6 x 6		6		Vekic et al. (1993)

C-DMFT

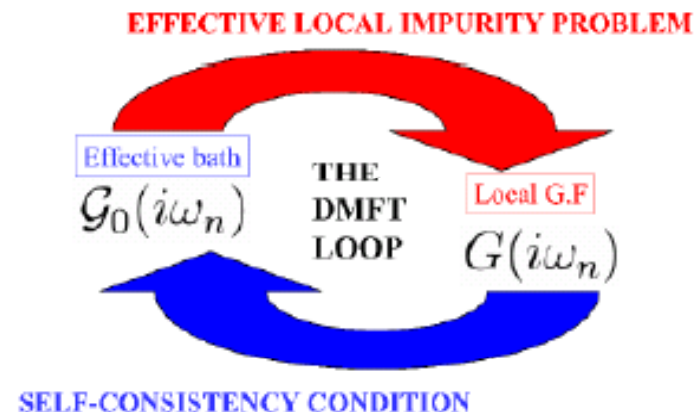


Mean-field is not a trivial problem! Many impurity solvers.

Here: continuous time QMC

P. Werner, PRL 2006
P. Werner, PRB 2007
K. Haule, PRB 2007

$$Z = \int \mathcal{D}[\psi^\dagger, \psi] e^{-S_c - \int_0^\beta d\tau \int_0^\beta d\tau' \sum_{\mathbf{k}} \psi_{\mathbf{k}}^\dagger(\tau) \Delta_{\mathbf{k}}(\tau, \tau') \psi_{\mathbf{k}}(\tau')}$$



$$\Delta(i\omega_n) = i\omega_n + \mu - \Sigma_c(i\omega_n)$$

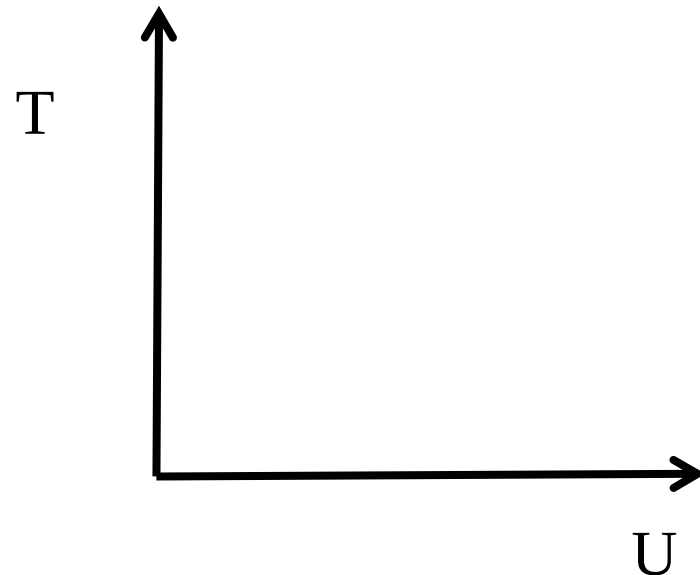
$$- \left[\sum_{\tilde{\mathbf{k}}} \frac{1}{i\omega_n + \mu - t_c(\tilde{\mathbf{k}}) - \Sigma_c(i\omega_n)} \right]^{-1}$$

Solving cluster in a bath problem

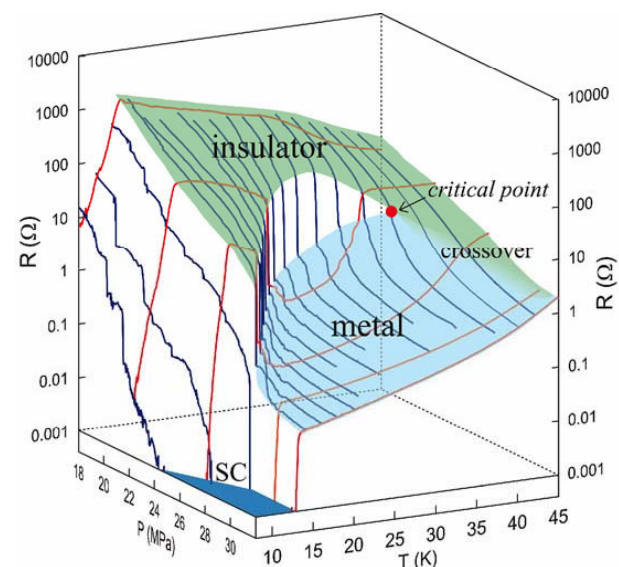
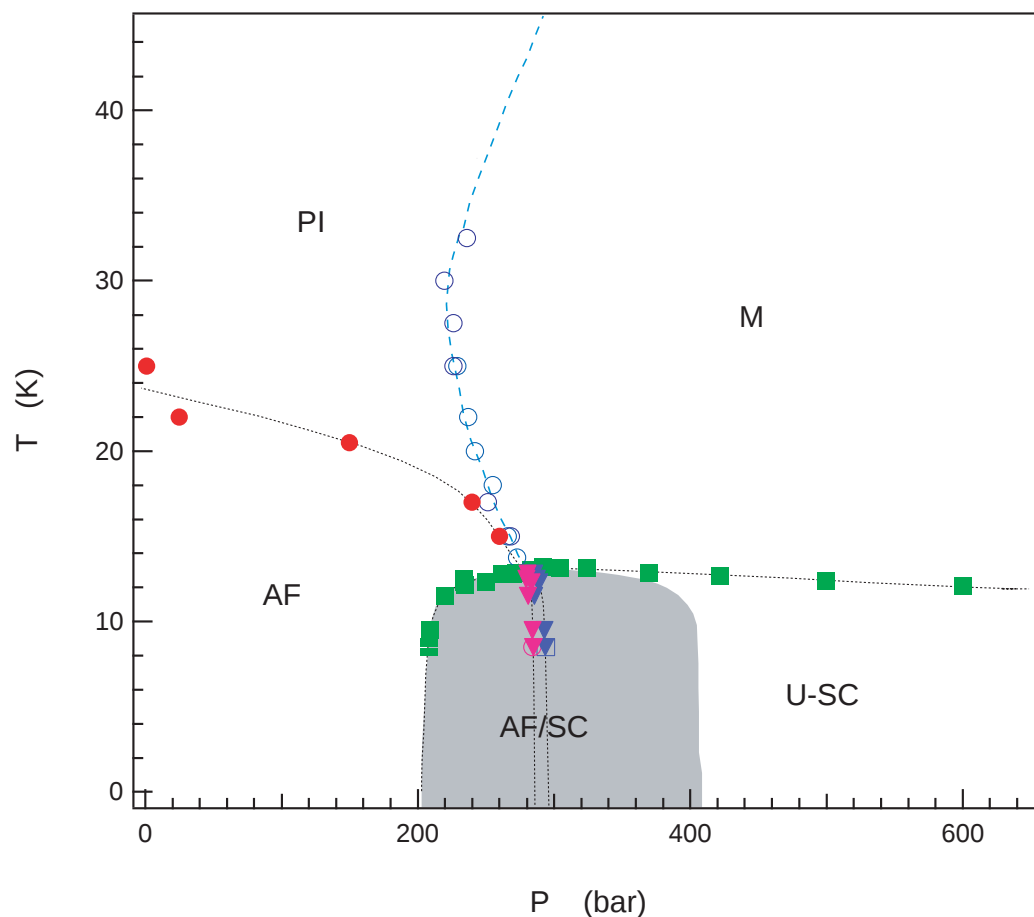
- Continuous-time Quantum Monte Carlo calculations to sum all diagrams generated from expansion in powers of hybridization.
 - P. Werner, A. Comanac, L. de' Medici, M. Troyer, and A. J. Millis, Phys. Rev. Lett. **97**, 076405 (2006).
 - K. Haule, Phys. Rev. B **75**, 155113 (2007).



Interaction-induced Mott transition $d = 2$



Bare Mott critical point in organics

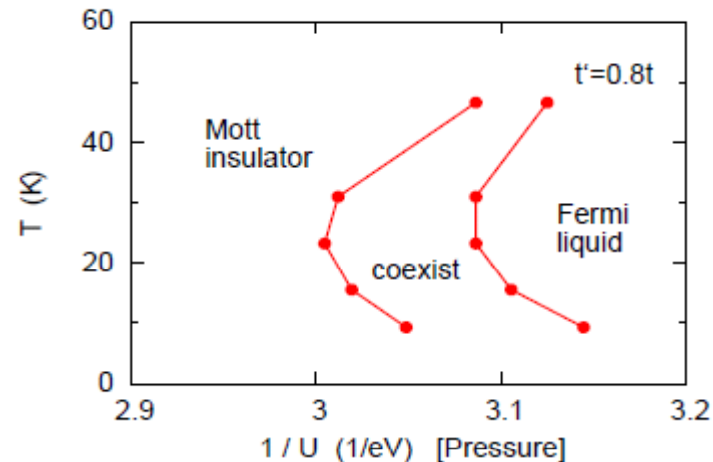


F. Kagawa, K. Miyagawa, + K. Kanoda
PRB **69** (2004) + Nature **436** (2005)

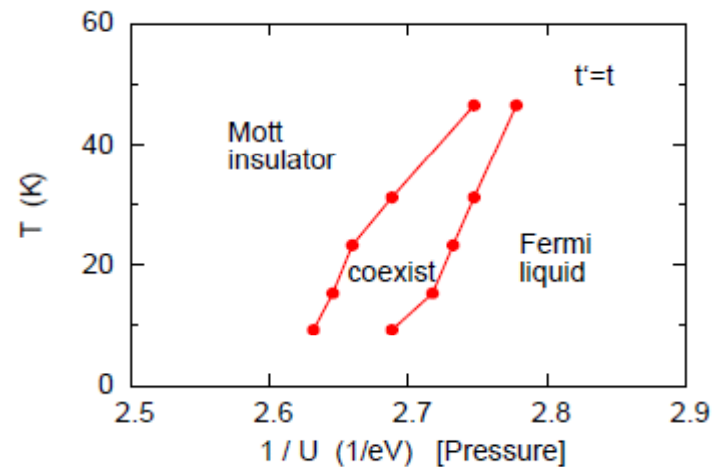
Phase diagram ($X=\text{Cu}[\text{N}(\text{CN})_2]\text{Cl}$)

S. Lefebvre et al. PRL **85**, 5420 (2000), P. Limelette, et al. PRL **91** (2003)

Interaction-induced Mott transition theory



κ -BEDT-Cl



κ -BEDT-CN

Liebsch Phys. Rev. B **79**, 195108 (2009)

See also: Ohashi et al. PRL **100**, 076402 (2008)

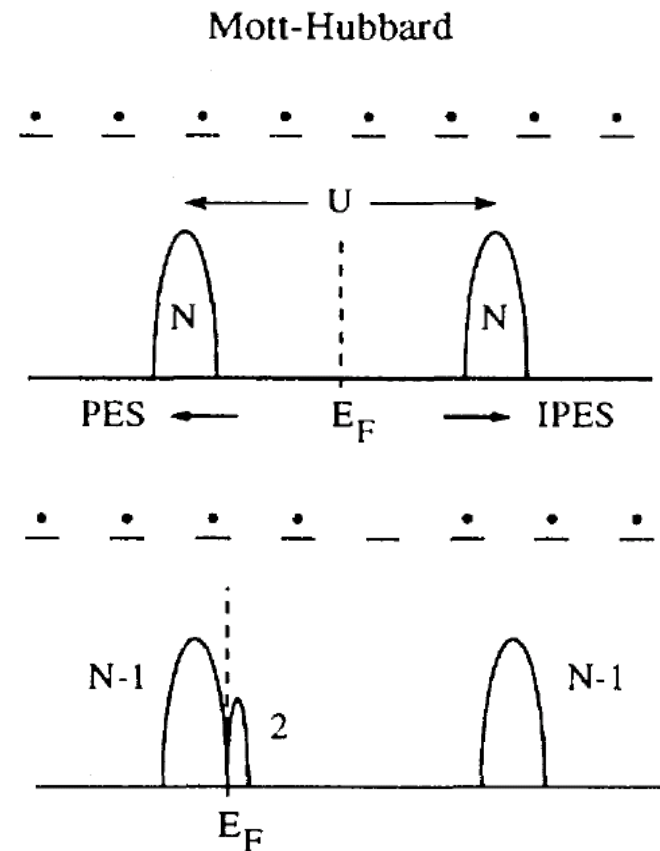
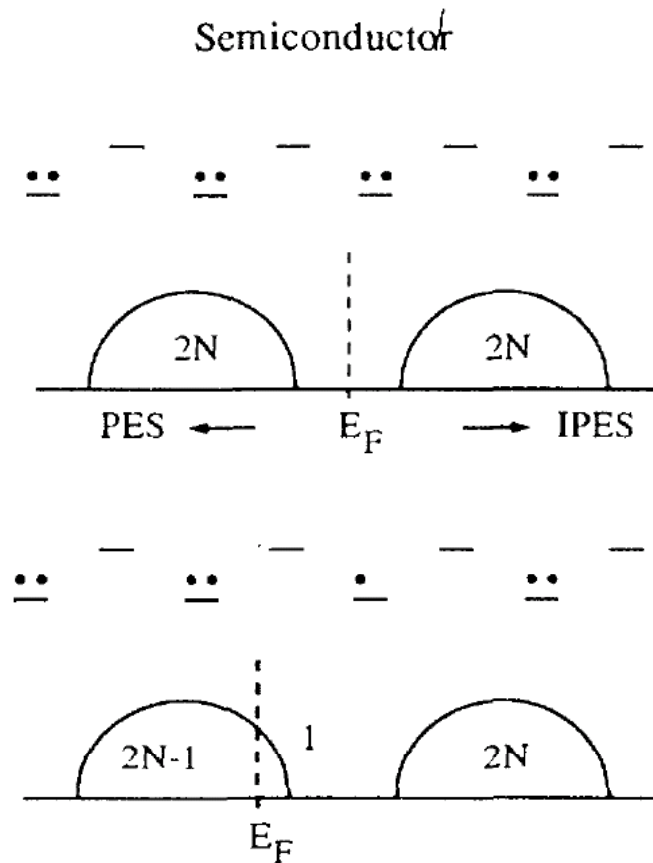


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Cuprates as doped Mott insulators

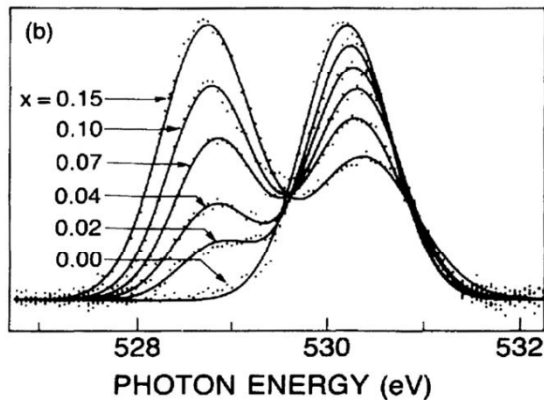


Spectral weight transfer

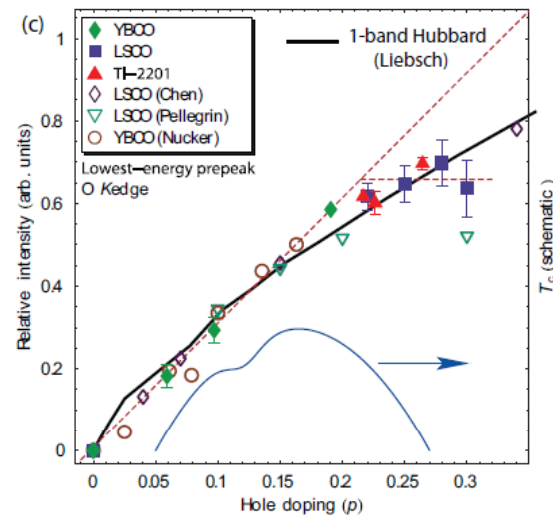


Meinders *et al.* PRB **48**, 3916 (1993)

Experiment: X-Ray absorption



Chen et al. PRL **66**, 104 (1991)



Peets et al. PRL **103**, (2009), Phillips, Jarrell arXiv

Number of low energy states above $\omega = 0$ scales as $2x +$
Not as $1+x$ as in Fermi liquid

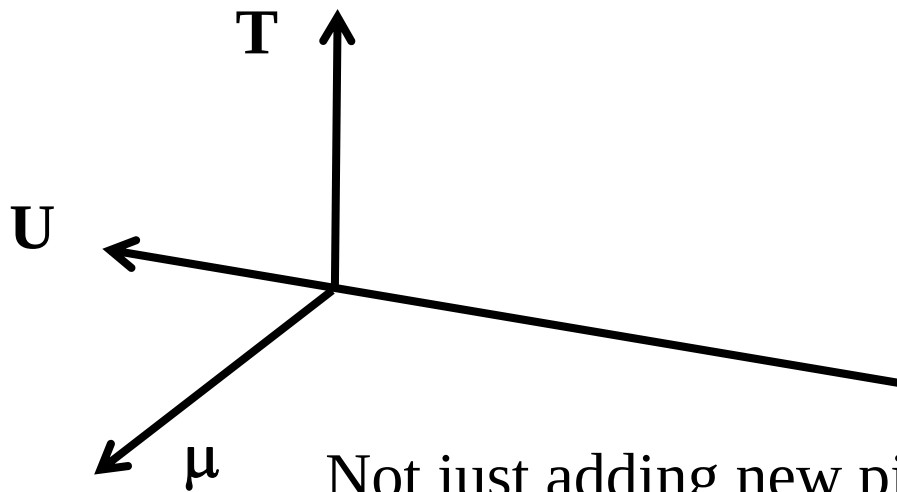
Meinders *et al.* PRB **48**, 3916 (1993)



Giovanni Sordi

G. Sordi, K. Haule, A.-M.S.T
PRL, **104**, 226402 (2010)

Doping-induced Mott transition ($t'=0$)



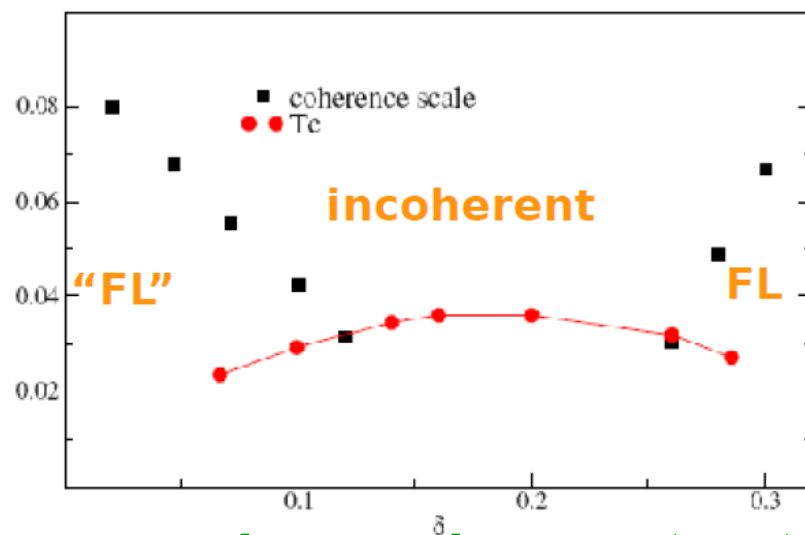
Kristjan Haule

Not just adding new piece:
Lesson from DMFT, first order transition + critical
point governs phase diagram

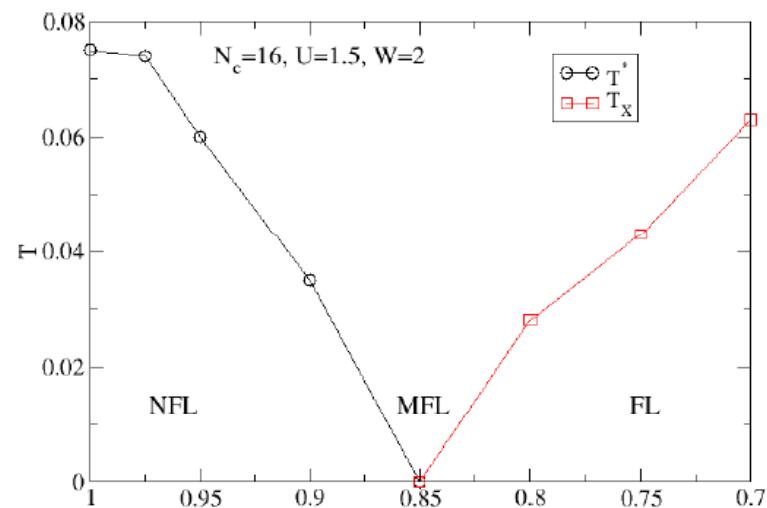


Doping driven Mott transition, $t' = 0$

Method	t'	Orbital selective	U	Critical point	Ref.
D+C+H 8			7		Werner et al. cond-mat (2009)
D+C+H 4					Gull et al. EPL (2008)
	-0.3		10,6		Liebsch, Merino... (2008)
					Ferrero et al. PRB (2009)
D+C+H 8			7		Gull, et al. PRB (2009)

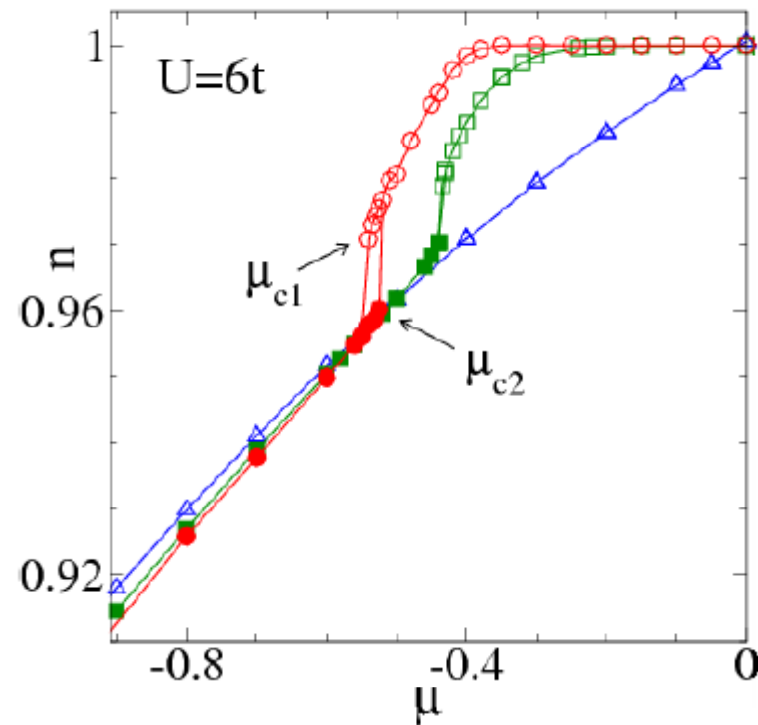


K. Haule, G. Kotliar, PRB (2008)

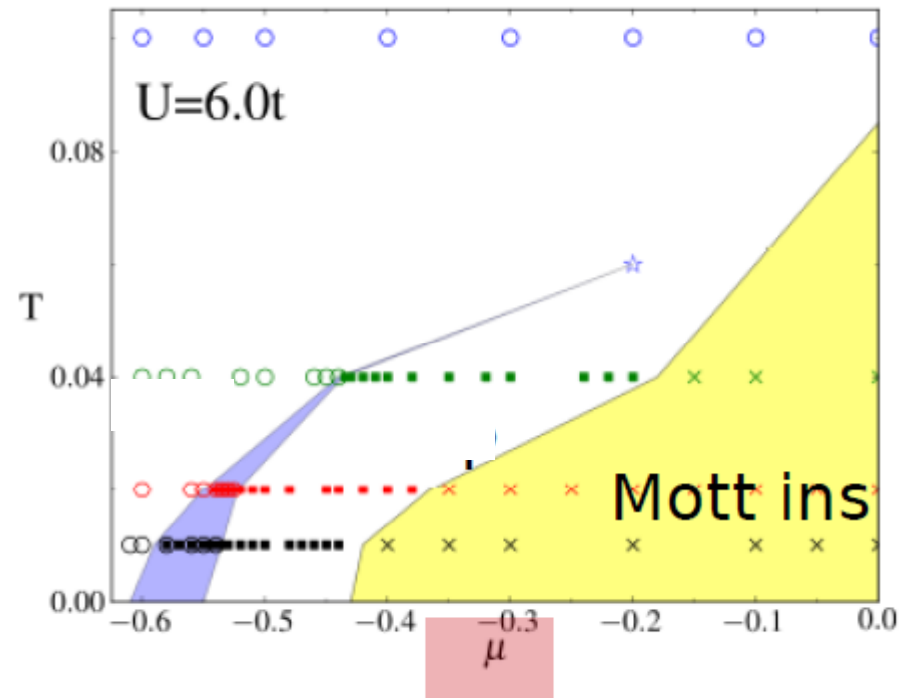


Vildhyadhiraja, PRL (2009)

First order transition at finite doping

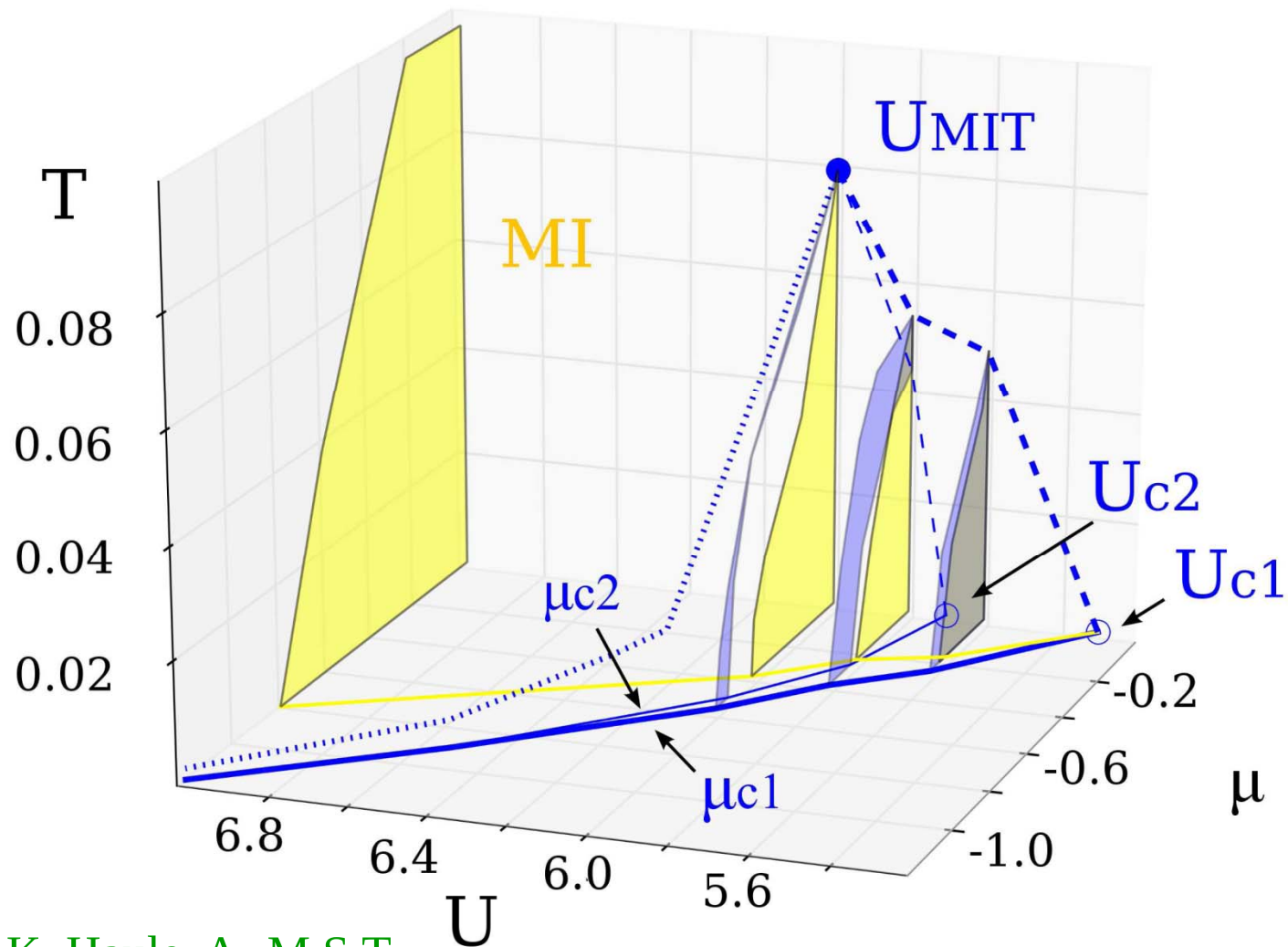


$n(\mu)$ for several temperatures:
 $T/t = 1/10$, $1/25$, $1/50$



Hysteretic behavior:
 fingerprint first order
 transition!

Normal state phase diagram



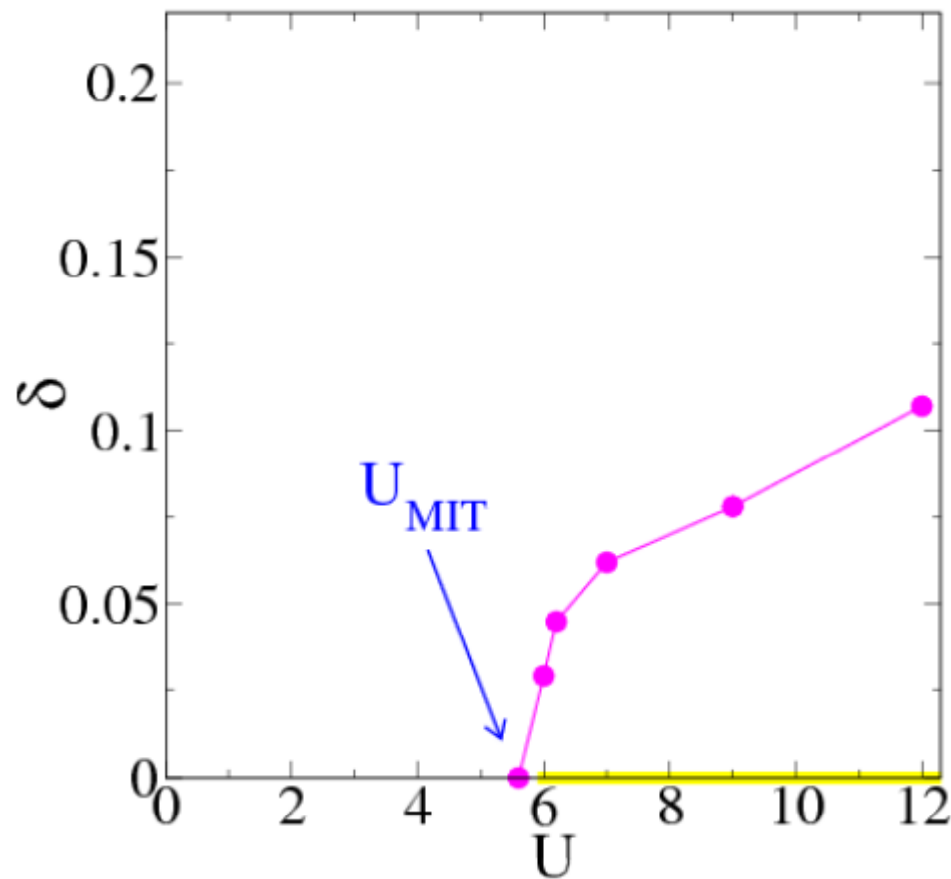
G. Sordi, K. Haule, A.-M.S.T
PRL, **104**, 226402 (2010)



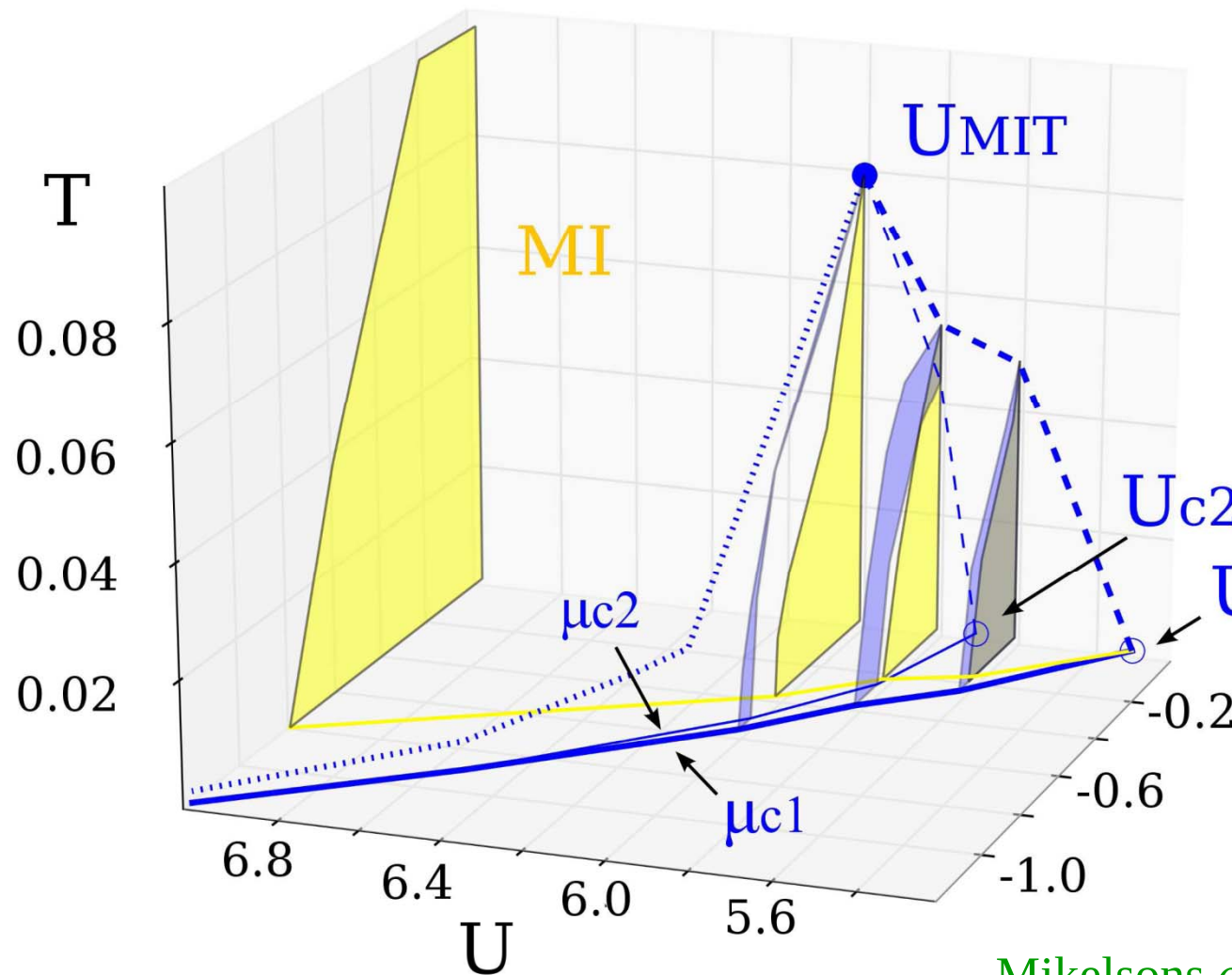
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Link to Mott transition up to optimal doping

Doping dependence of critical point as a function of U



Thermodynamic properties



Clausius-Clapeyron

$$\left(\frac{dT_c}{d\mu_c} \right)_U = \frac{(n_1 - n_2)}{(S_2 - S_1)}$$

$$\left(\frac{dU_c}{d\mu_c} \right)_T = \frac{(n_1 - n_2)}{(D_1 - D_2)}$$

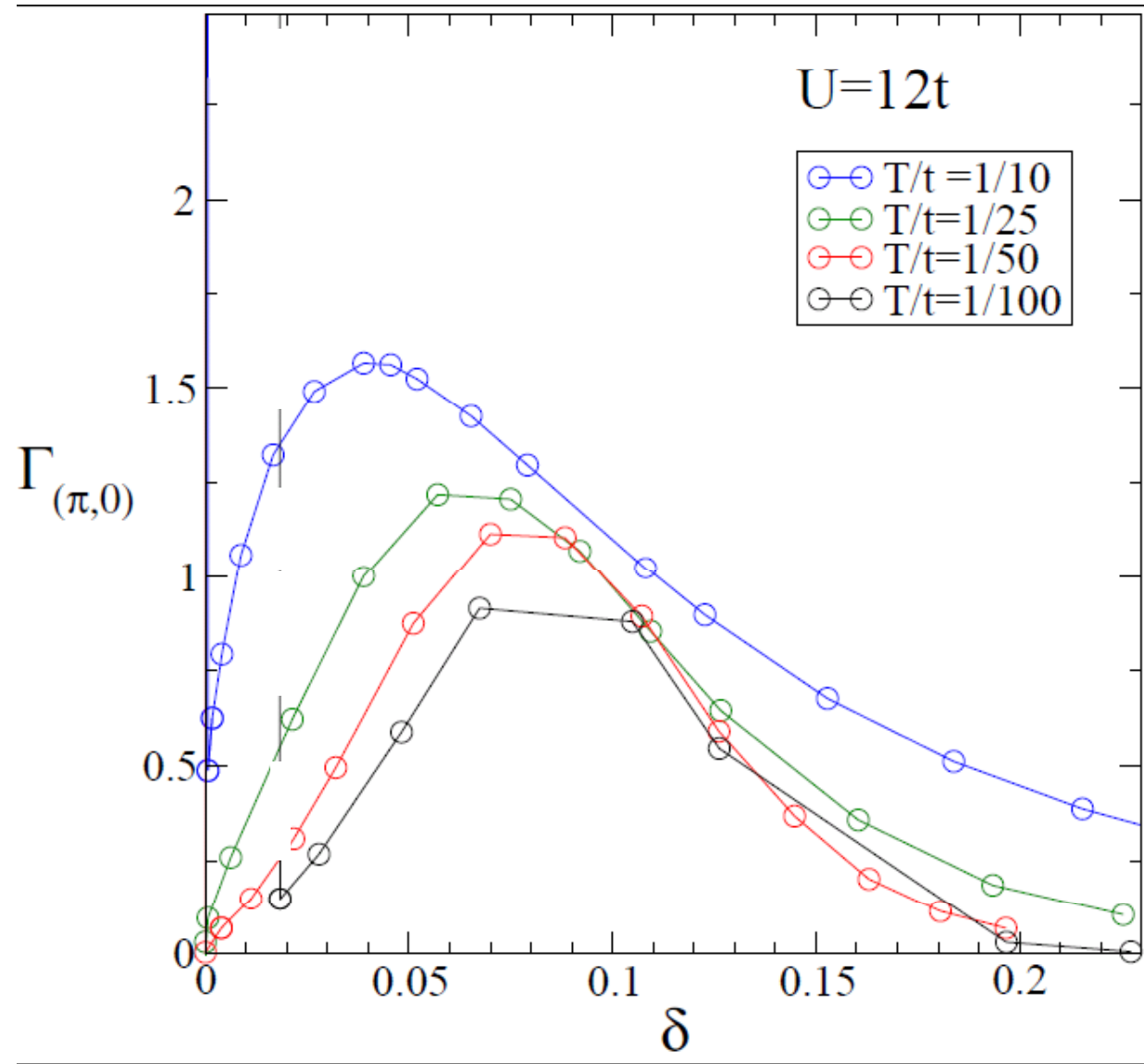
Underdoped phase
lower D and S
More compressible

$$\mu \quad \left(\frac{\partial S}{\partial n} \right)_T = 0$$

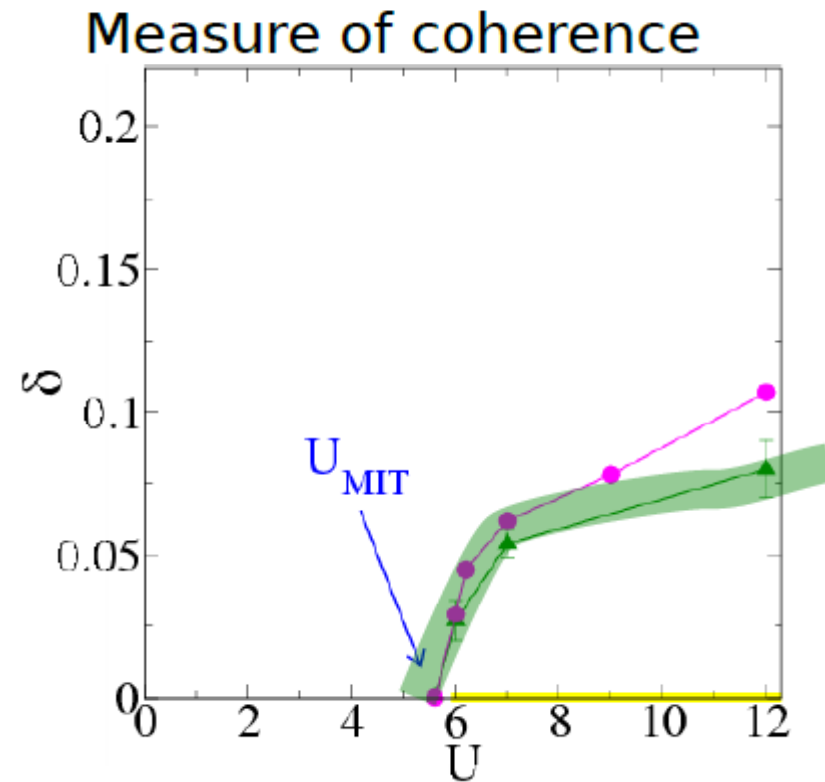
Mikelsons *et al.* PRB **80**, (2009)



Maximum scattering rate



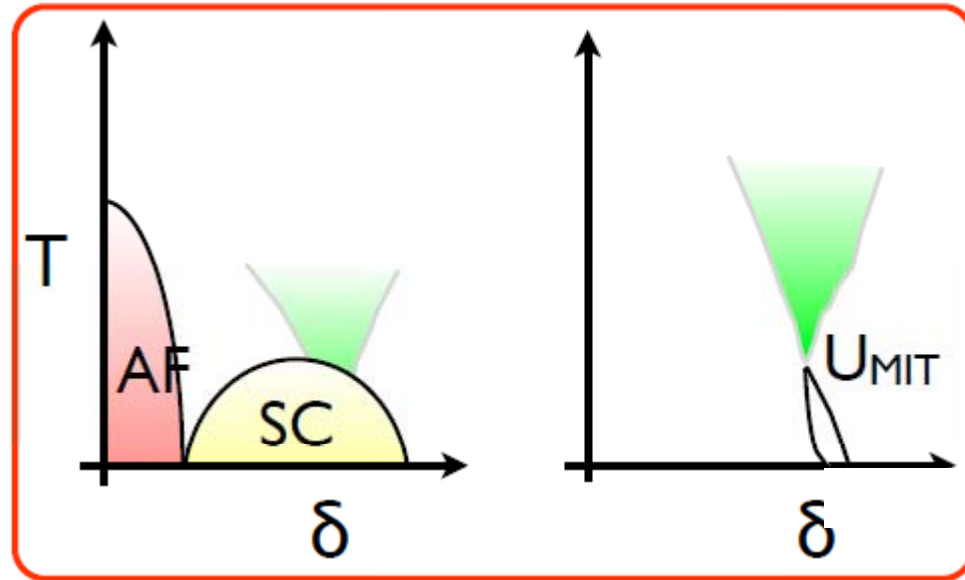
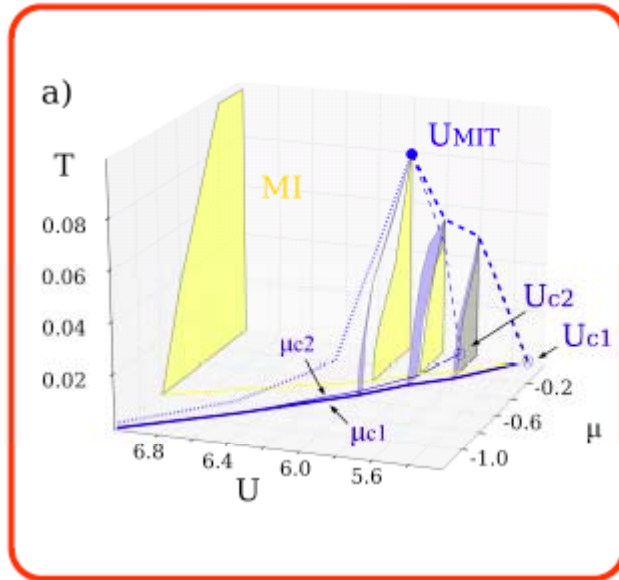
Maximum scattering rate from Mott Physics



δ

Max [$\text{Im}\Sigma_K(w \rightarrow 0)$] originate at U_{MIT}

What is under the dome



Giovanni Sordi



Kristjan Haule

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Underdoped metal very sensitive to anisotropy

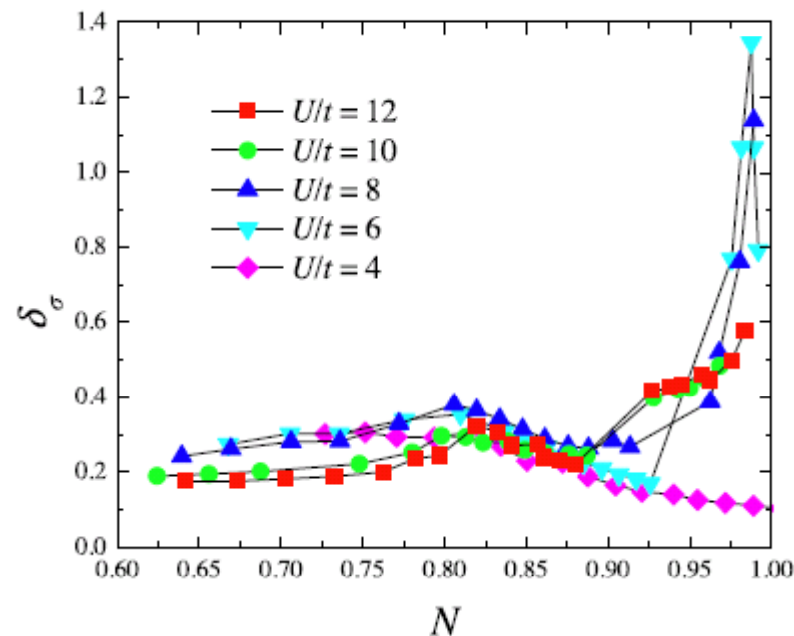
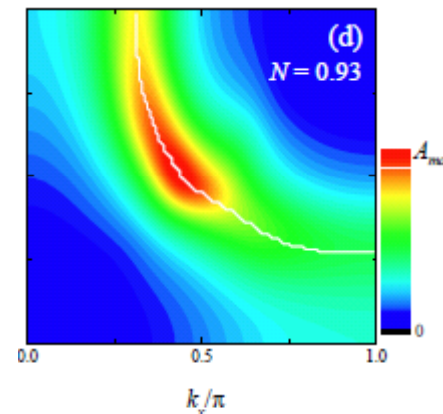
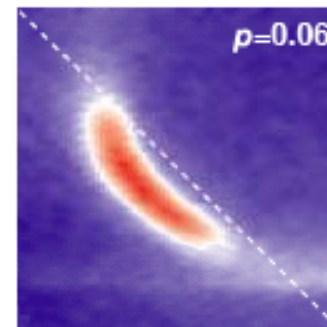


FIG. 3: (Color online) Anisotropy in the CDMFT conductivity $\delta_\sigma = 2 [\sigma_x(0) - \sigma_y(0)] / [\sigma_x(0) + \sigma_y(0)]$ as a function of filling N for various values of U and $\eta = 0.1$, $\delta_0 = 0.04$.

Okamoto, Sénéchal, Civelli, AMST
PRB (R) in press,



g



Satoshi Okamoto



David Sénéchal



D. Fournier *et al.* Nature Physics (Marcello Civelli)

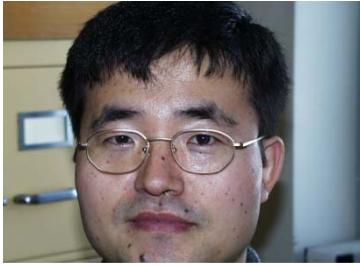
Superconductivity

Phase diagram

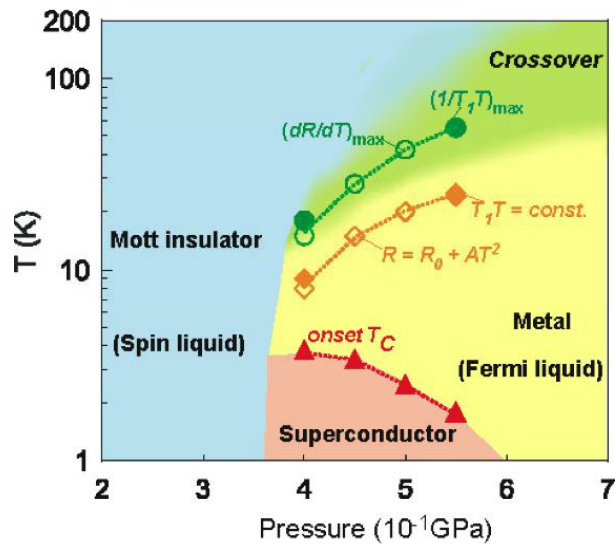
Exact diagonalization as impurity
solver ($T=0$).



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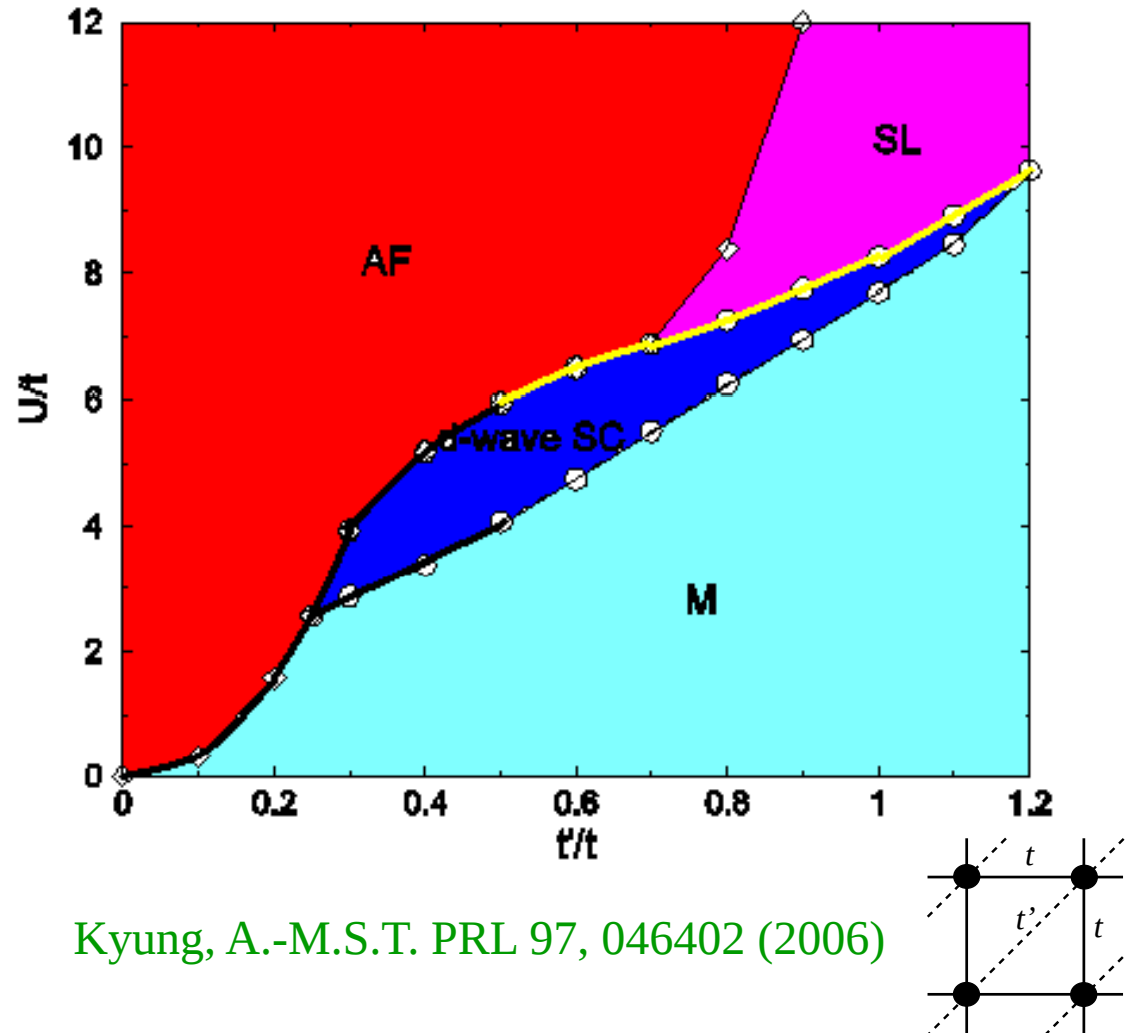


Theoretical phase diagram BEDT



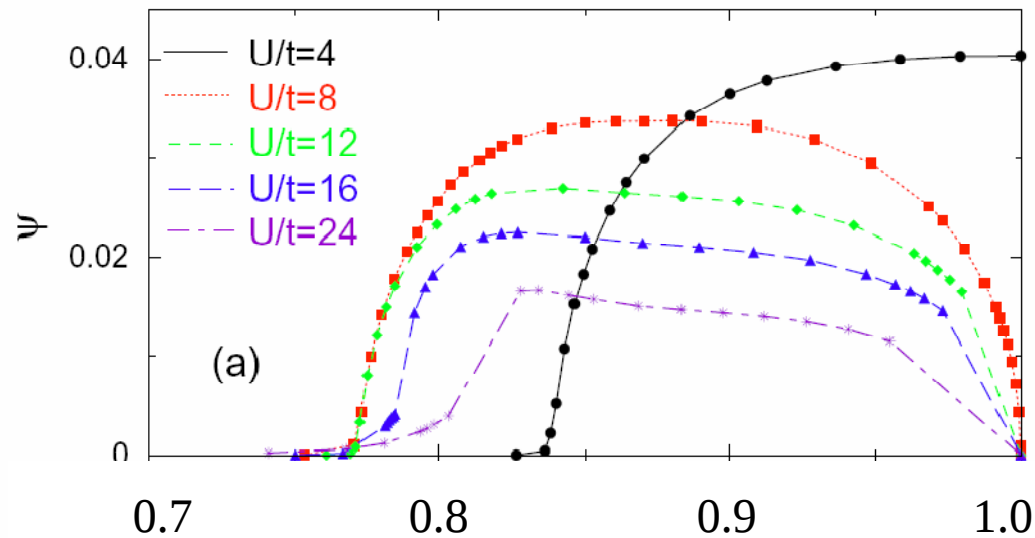
Y. Kurisaki, et al.

Phys. Rev. Lett. **95**, 177001(2005) Y. Shimizu, et al. Phys. Rev. Lett. **91**, (2003)



Kyung, A.-M.S.T. PRL 97, 046402 (2006)

Dome vs Mott (CDMFT)

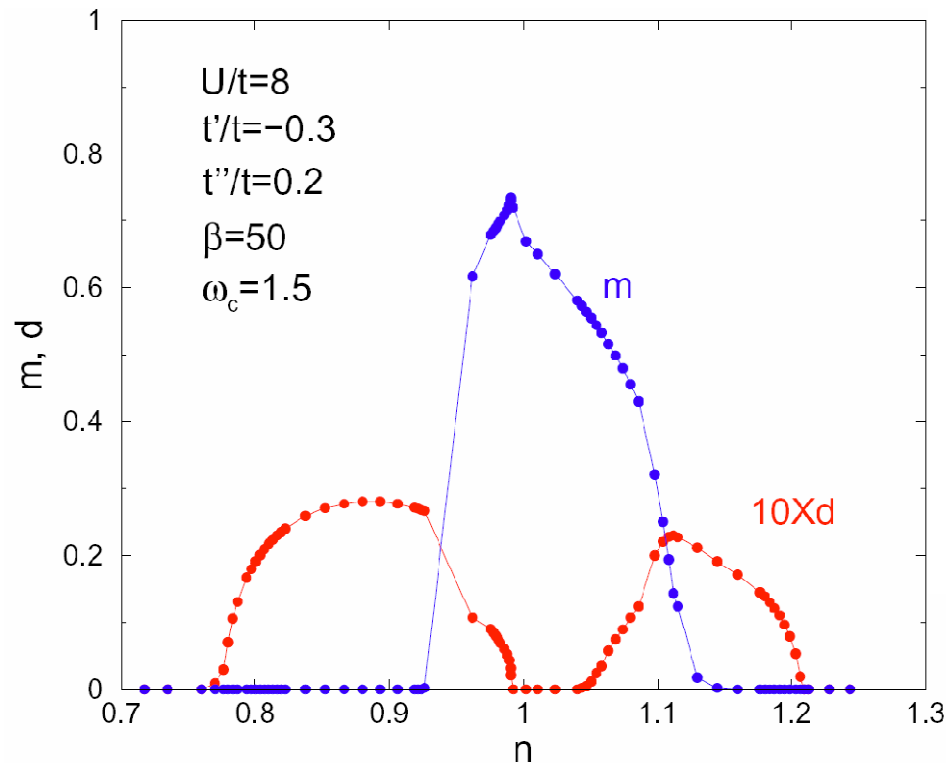


Kancharla, Kyung, Civelli,
Sénéchal, Kotliar AMST
Phys. Rev. B (2008)

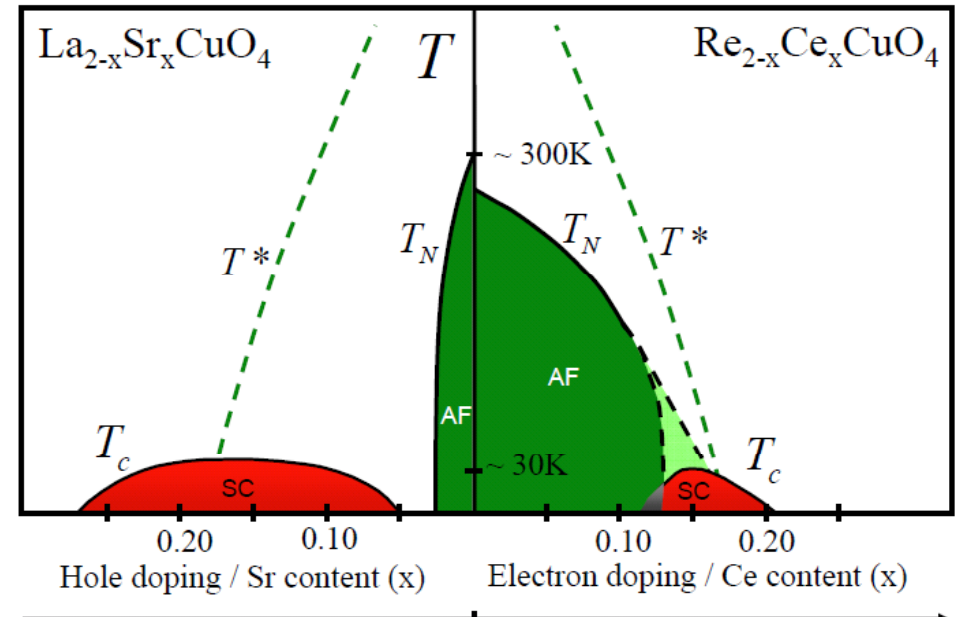


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CDMFT global phase diagram



Kancharla, Kyung, Civelli,
 S  n  chal, Kotliar AMST
 Phys. Rev. B (2008)

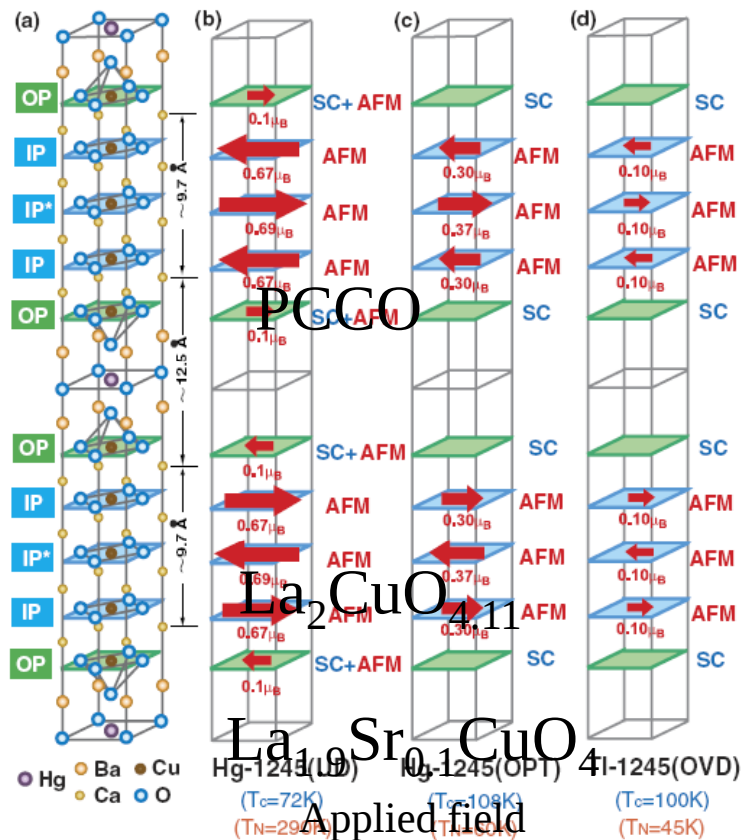


Armitage, Fournier, Greene, RMP (2009)



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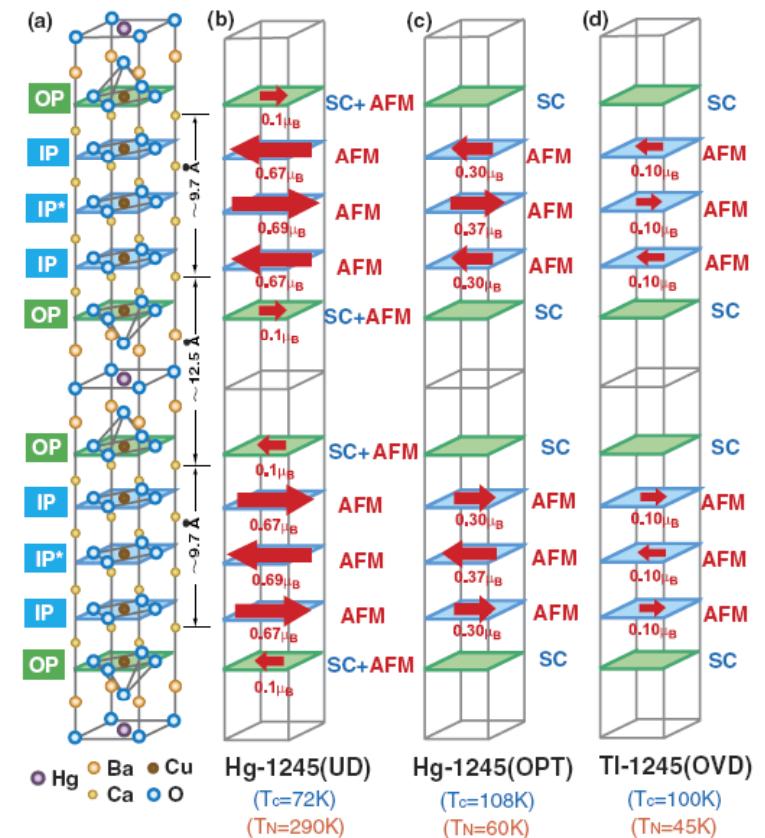
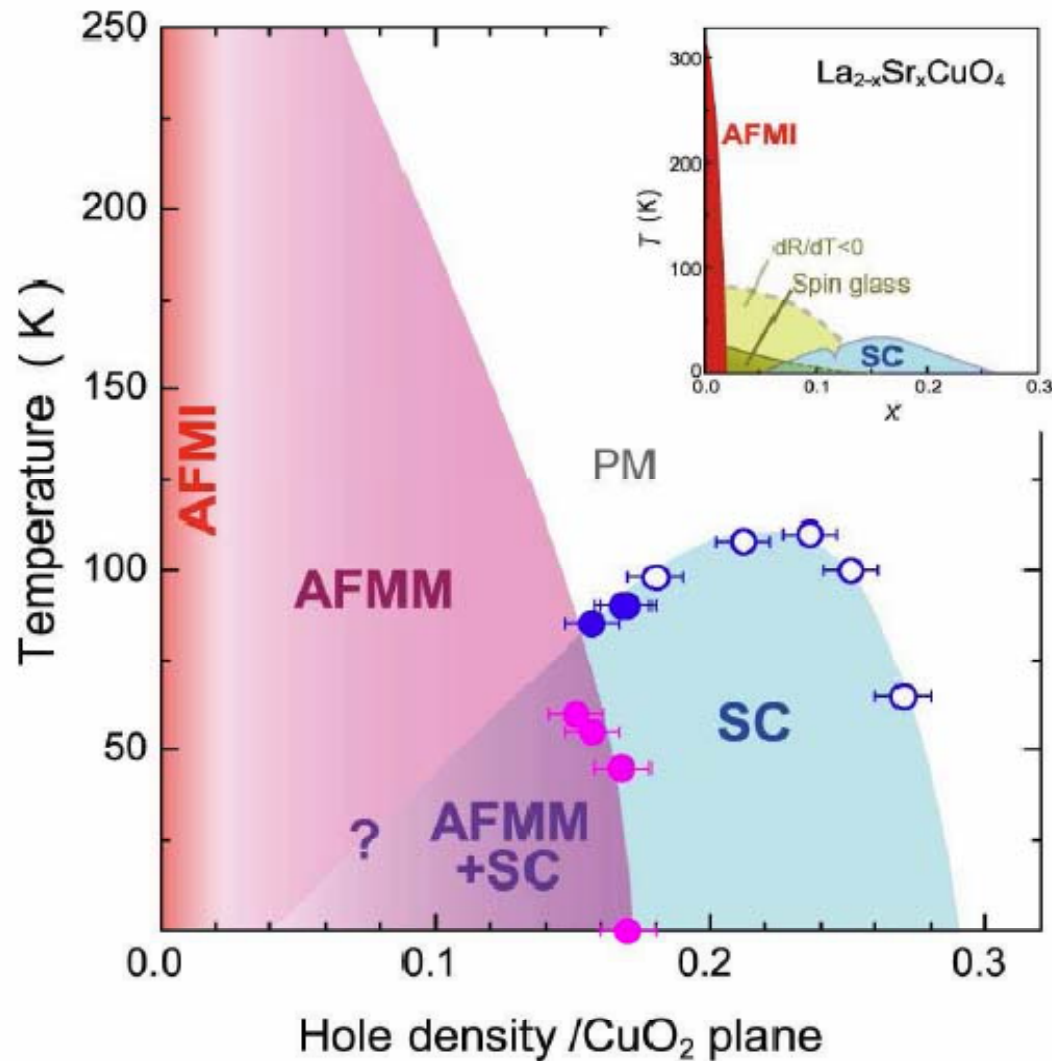
Homogeneous coexistence (experimental)



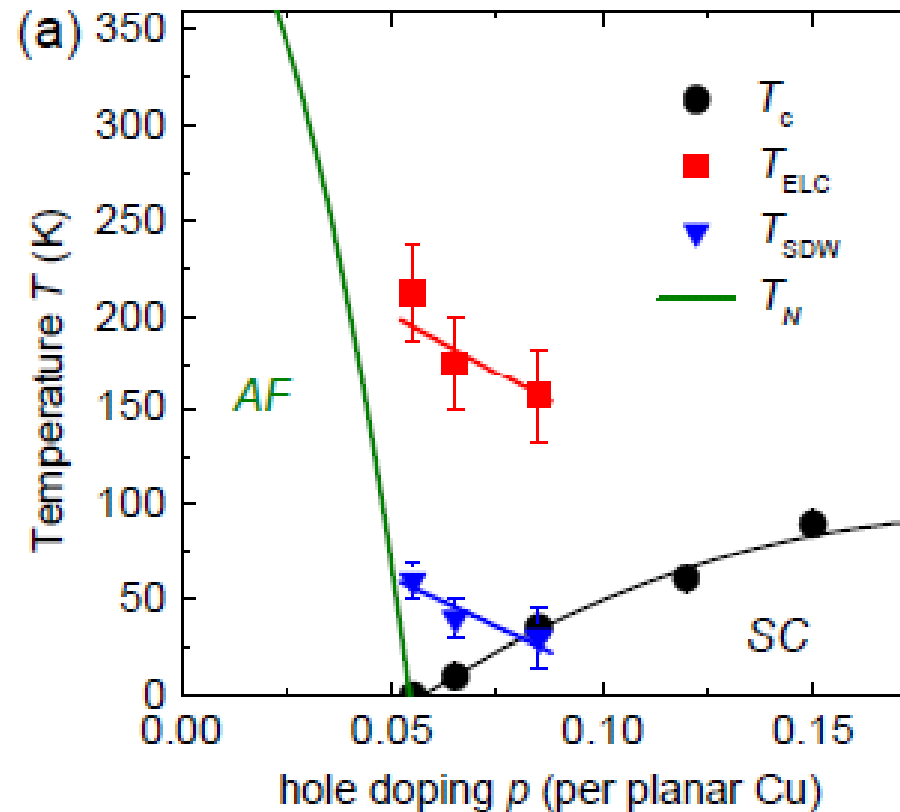
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Consistent with following experiment

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Magnetic phase diagram of YBCO



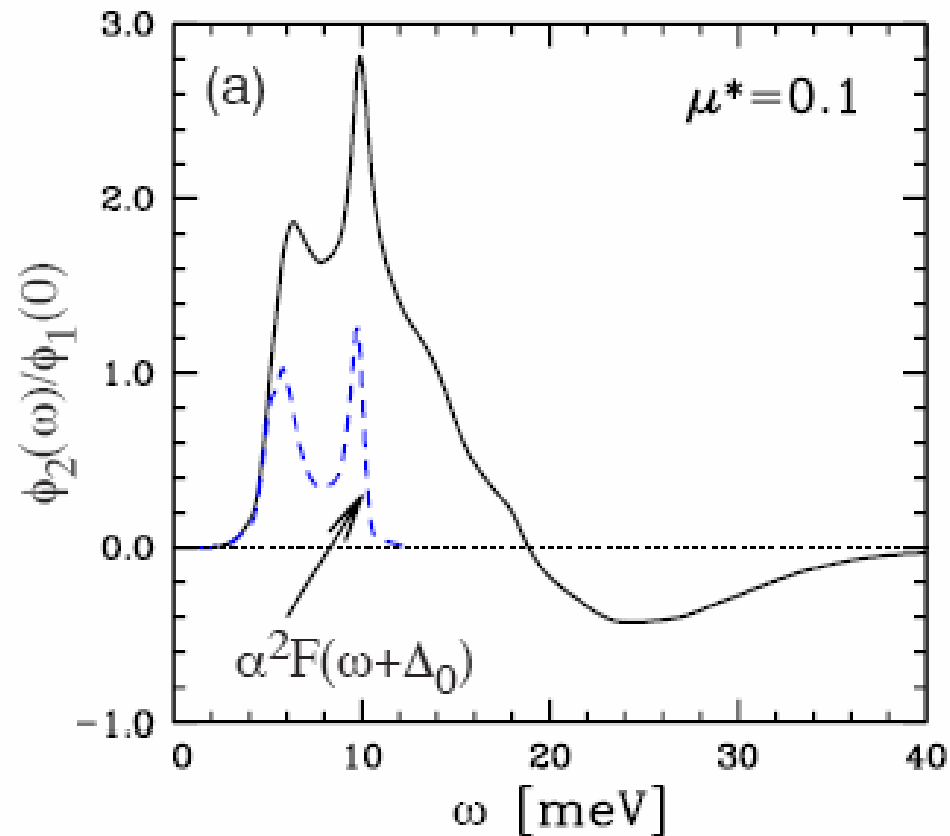
Haug, ... Keimer, arXiv:1008.4298

The glue



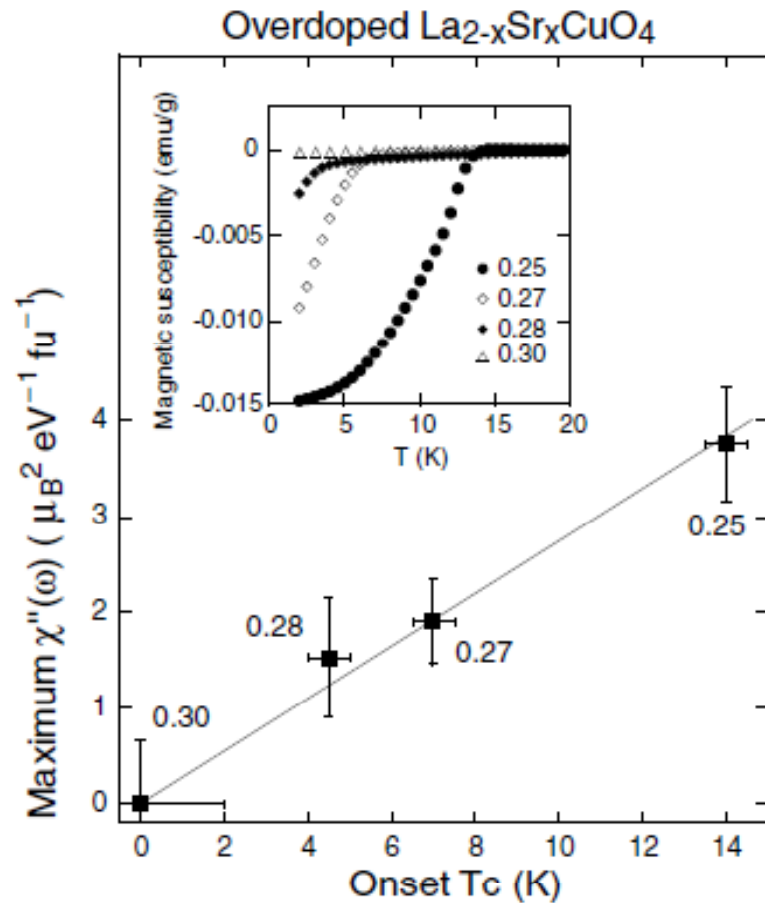
$\text{Im } \Sigma_{\text{an}}$ and electron-phonon in Pb

Maier, Poilblanc, Scalapino, PRL (2008)

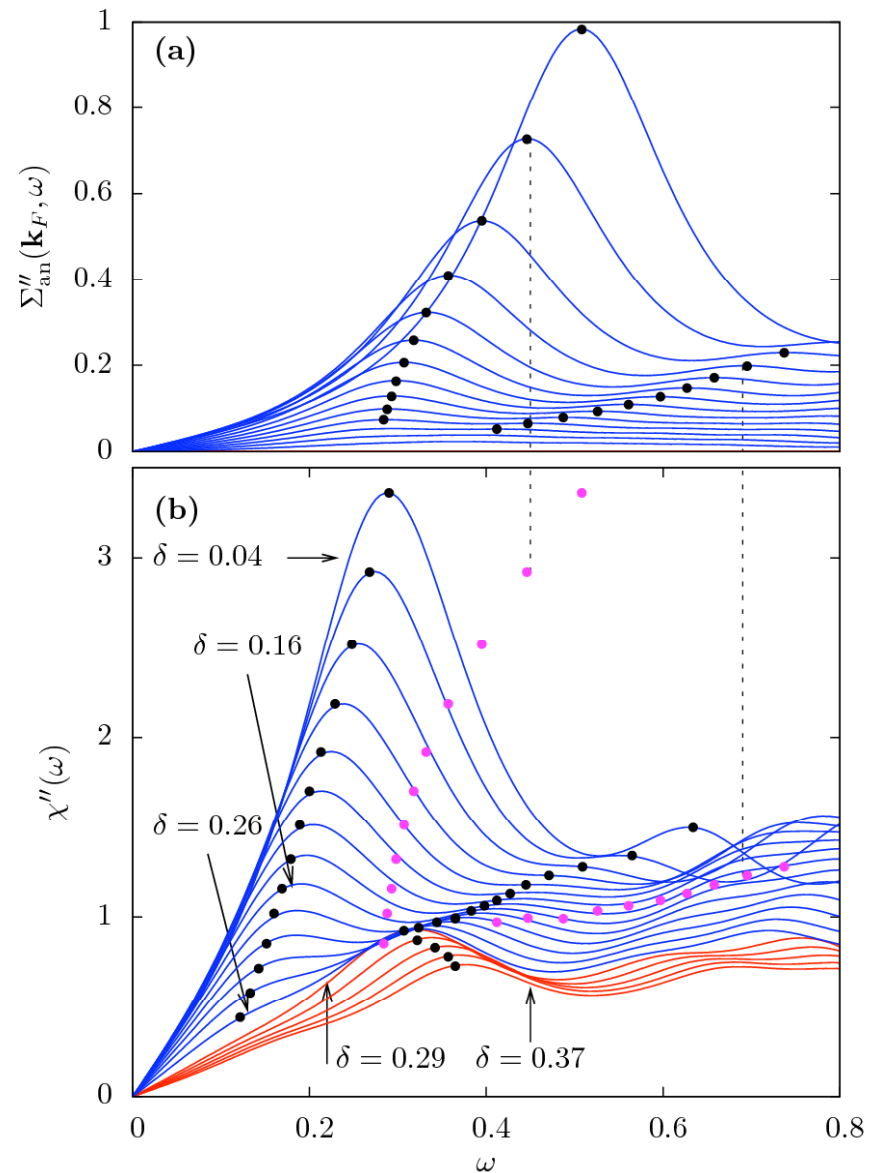


The glue

Kyung, S  n  chal, Tremblay, Phys. Rev. B
80, 205109 (2009)



Wakimoto ... Birgeneau
 PRL (2004)



The glue and neutrons

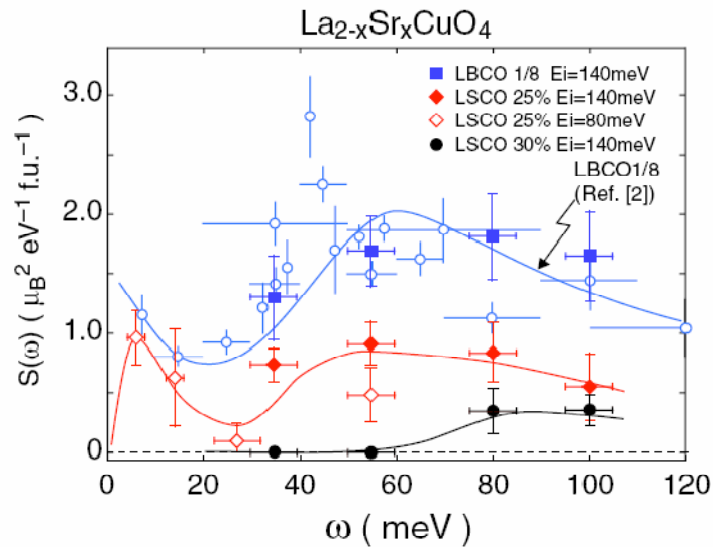
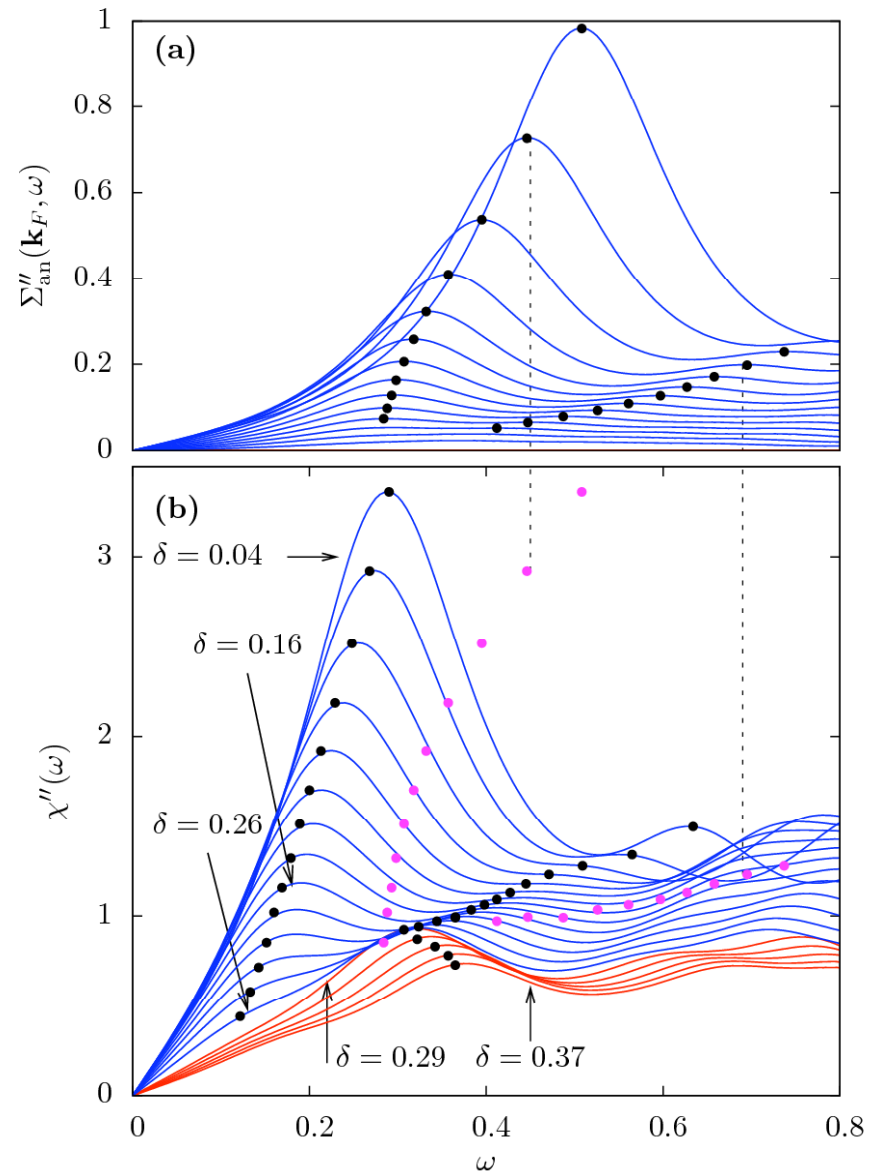


FIG. 3 (color online). **Q**-integrated dynamic structure factor $S(\omega)$ which is derived from the wide- H integrated profiles for LBCO 1/8 (squares), LSCO $x = 0.25$ (diamonds; filled for $E_i = 140$ meV, open for $E_i = 80$ meV), and $x = 0.30$ (filled circles) plotted over $S(\omega)$ for LBCO 1/8 (open circles) from [2]. The solid lines following data of LSCO $x = 0.25$ and 0.30 are guides to the eyes.

Wakimoto ... Birgeneau PRL (2007);
PRL (2004)



Main collaborators



Giovanni Sordi



David Sénéchal



Kristjan Haule



Bumsoo Kyung



Marcello Civelli



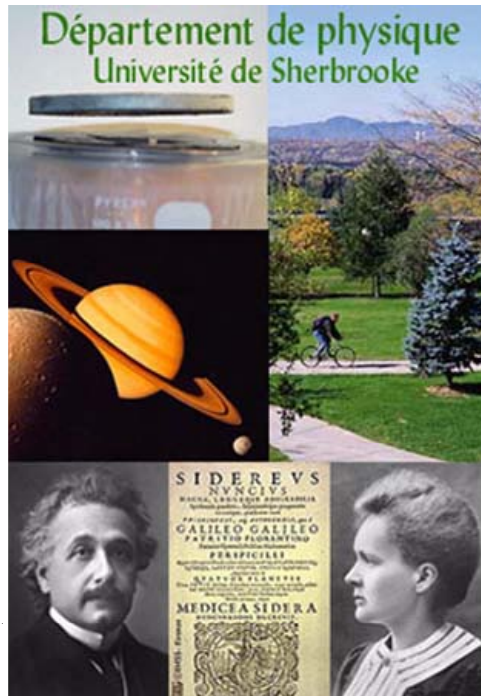
Satoshi Okamoto



Conclusions

- The influence of Mott Physics extends way beyond half-filling
- Conjecture that quantum-critical like behavior is constant U cut of our phase diagram, i.e. very low T critical point.
- Superconductivity follows naturally and retardation effects in pairing come from spin fluctuations.

André-Marie Tremblay



Le regroupement québécois sur les matériaux de pointe



Sponsors:



Merci

Arigato

Thank you