

What is special about strongly correlated superconductivity

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Giovanni Sordi, Patrick Sémon, Kristjan Haule,
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**ISS22, ENSE³-2, Exotic Superconductivity,
SCES2014 Grenoble, 9 July 2014**



Weakly and strongly correlated antiferromagnets

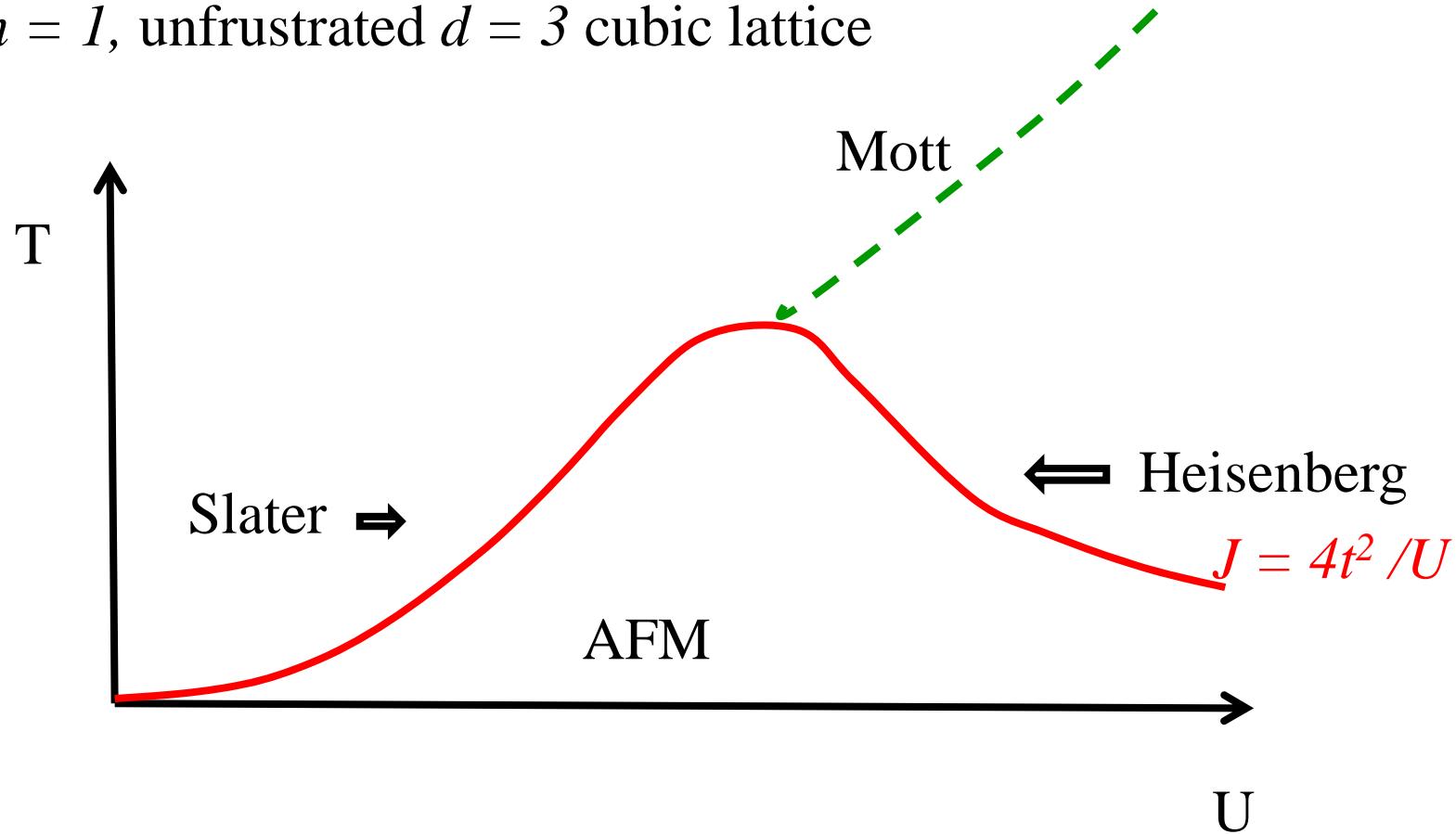
What is a phase?



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Weak vs Strong correlations

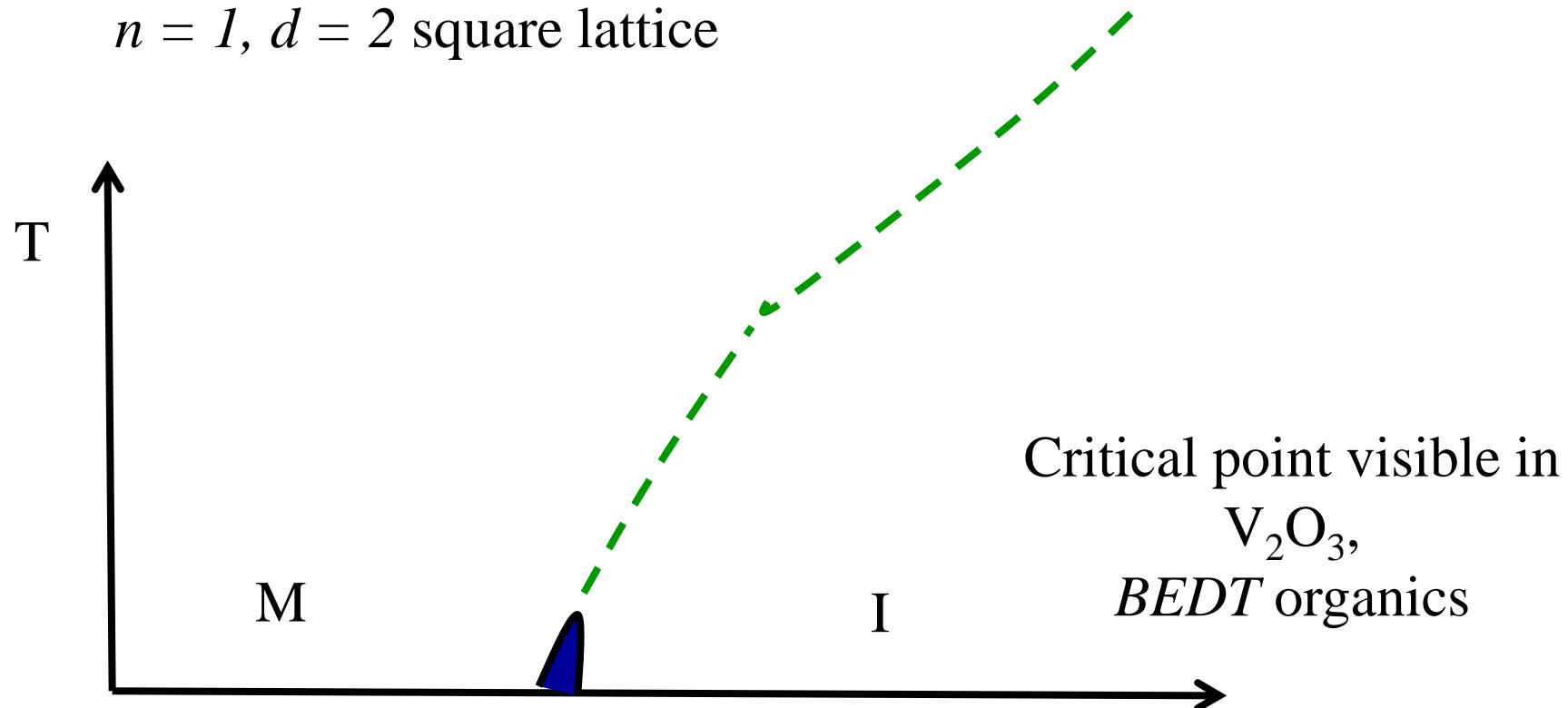
$n = 1$, unfrustrated $d = 3$ cubic lattice



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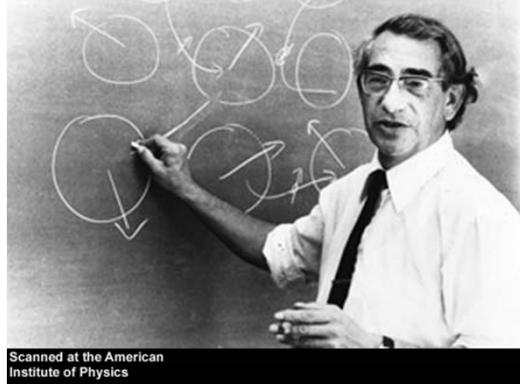
Local moment and Mott transition

$n = 1, d = 2$ square lattice



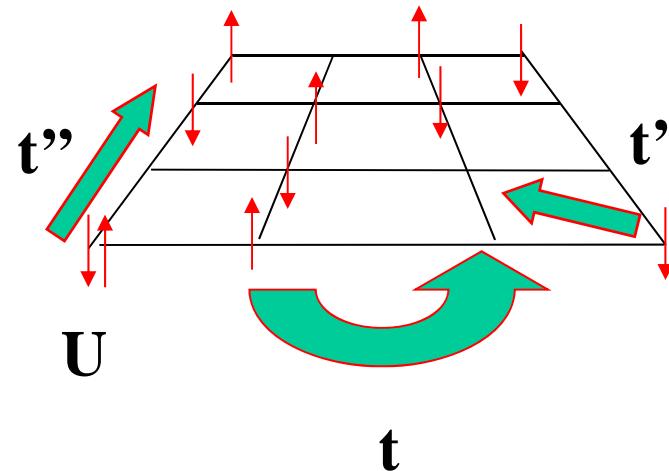
Understanding finite temperature phase from a *mean-field theory* down to $T = 0$

Hubbard model

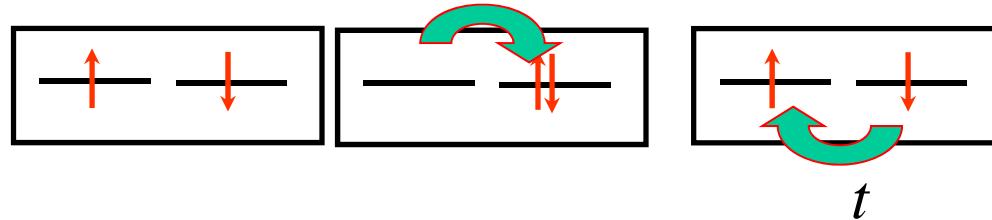


Scanned at the American
Institute of Physics

1931-1980



$$H = -\sum_{\langle ij \rangle \sigma} t_{i,j} (c_{i\sigma}^\dagger c_{j\sigma} + c_{j\sigma}^\dagger c_{i\sigma}) + U \sum_i n_{i\uparrow} n_{i\downarrow}$$



$t = 1$

Effective model, Heisenberg: $J = 4t^2 / U$



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Overview

- **An appetizer:**
 - Two gaps
 - Superfluid density
- A cartoon picture of pairing at weak and strong correlations
- The view from dynamical mean-field theory
 - The glue
 - A paradox: The goo remover (V) ?



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Two gaps in underdoped regime of cuprates

Le Tacon *et al.* Nature Physics 2, 537 - 543 (2006)

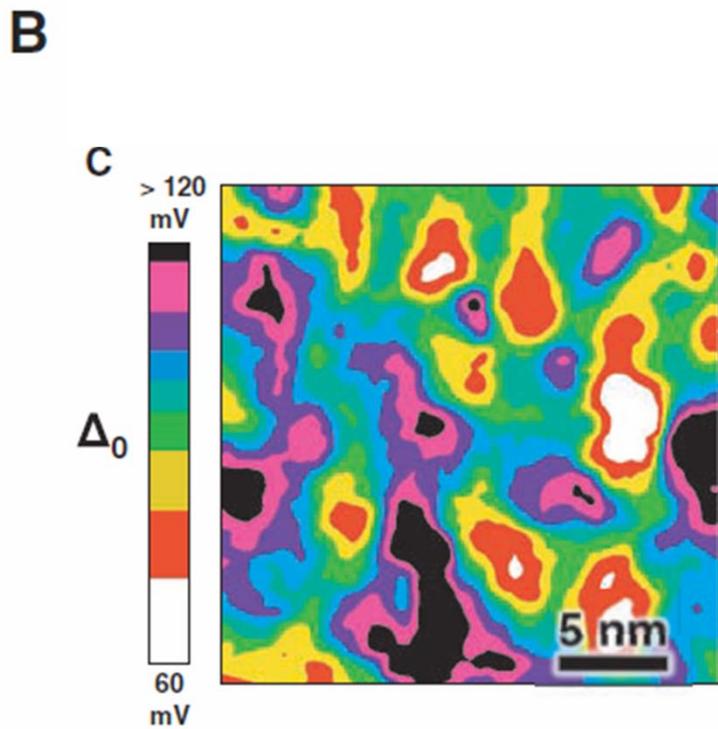
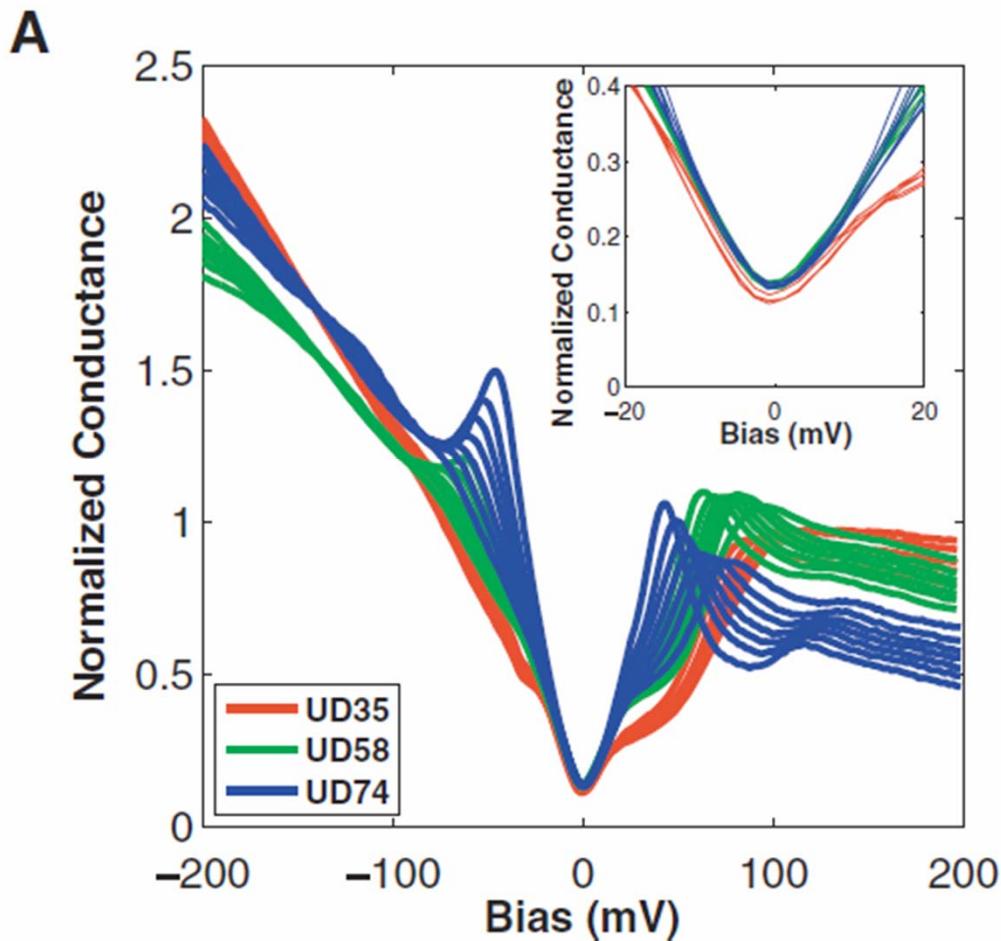
....

Sakai et al. PRL 111, 107001 (2013)



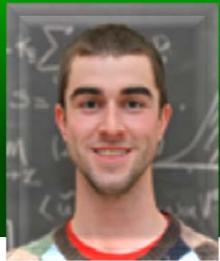
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Superconducting gap in STM



A. Pushp, Parker, ... A. Yazdani,
Science **364**, 1689 (2009)





Simon Verret

Experiment vs Theory, STM



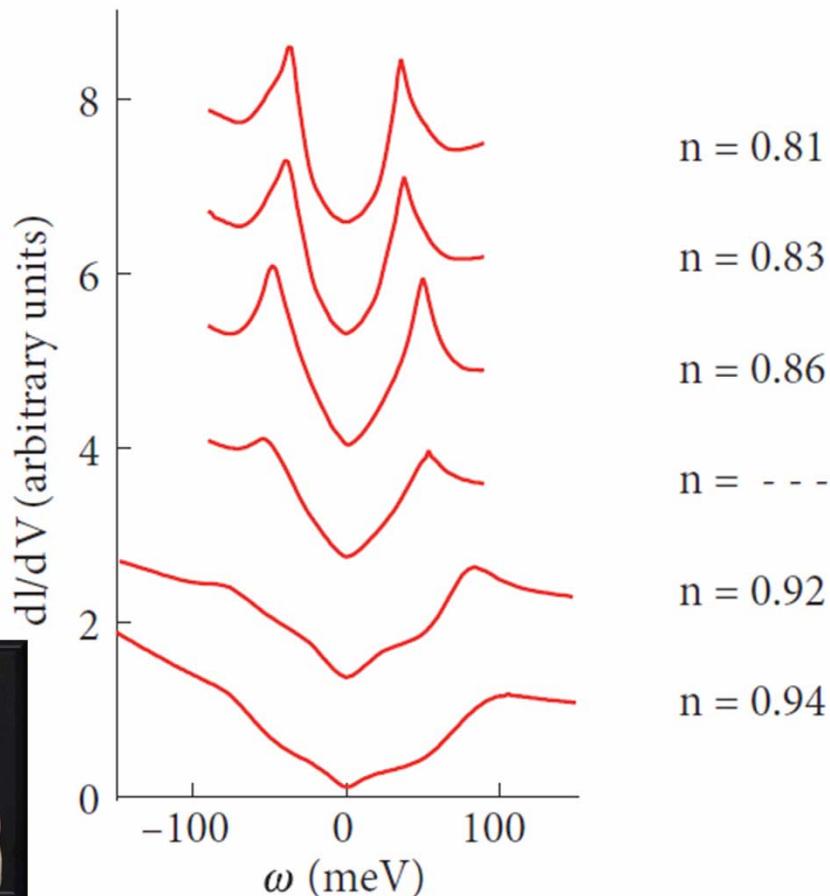
Jyotirmoy Roy

STM data

Kohsaka *et al.*, Nature **454** 1072 (2008)

CDMFT

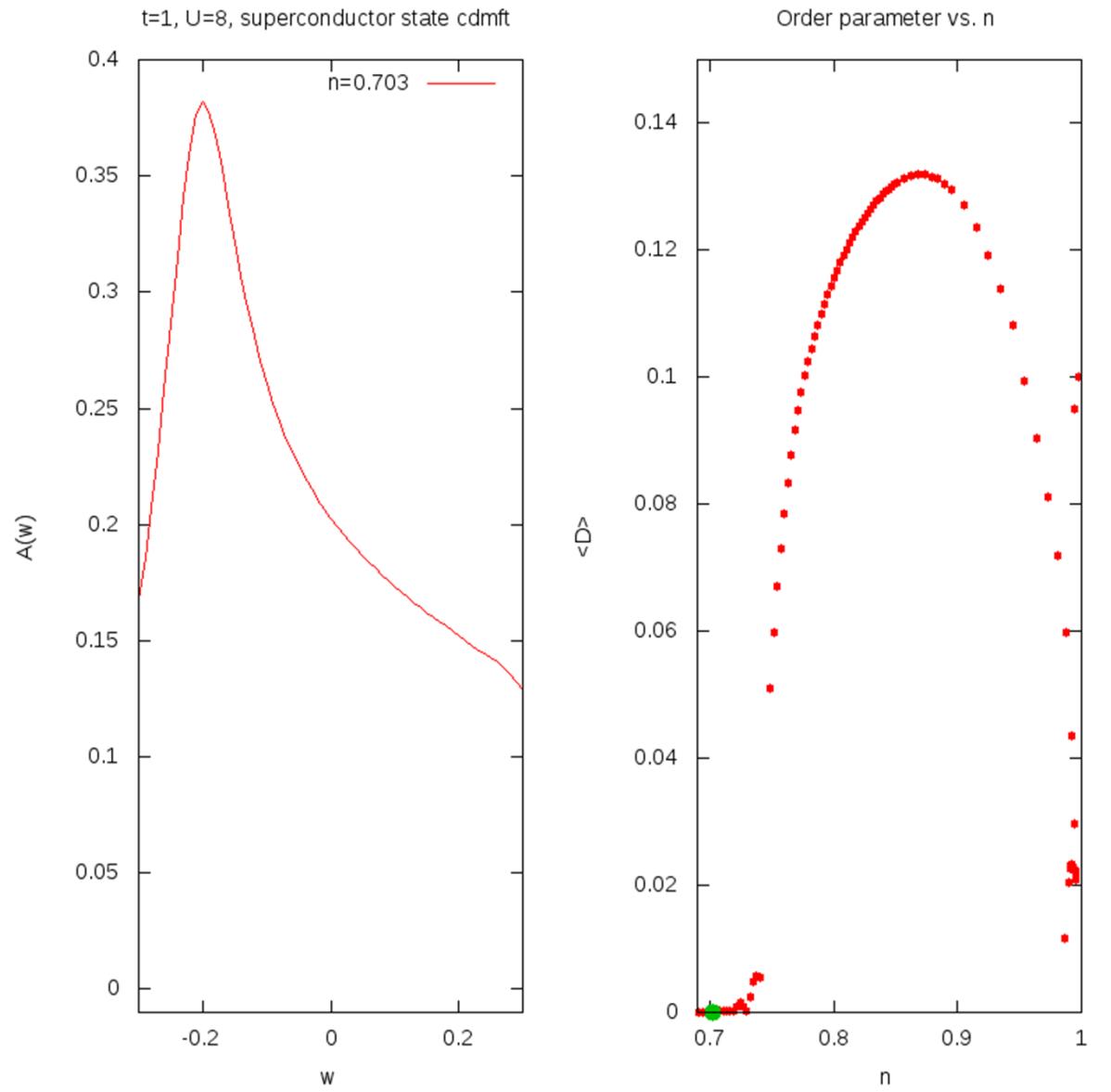
Unpublished



David Sénéchal

Evolution of SC gap and pseudogap with n

$t' = -0.3 t$
 $t'' = 0.2 t$
 $U=8t$

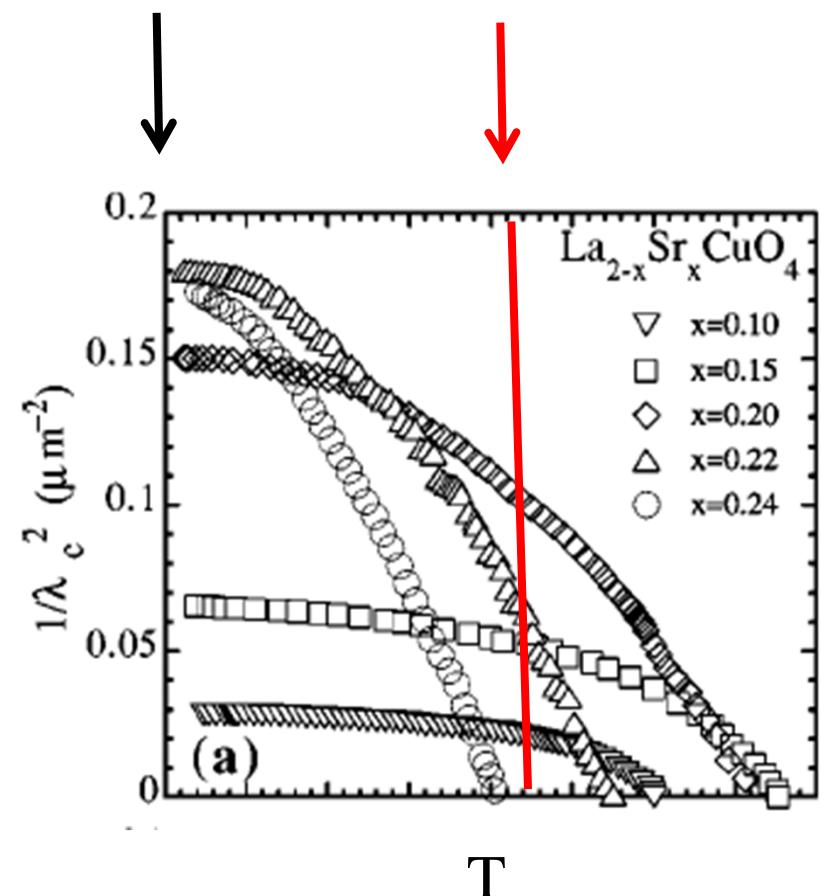
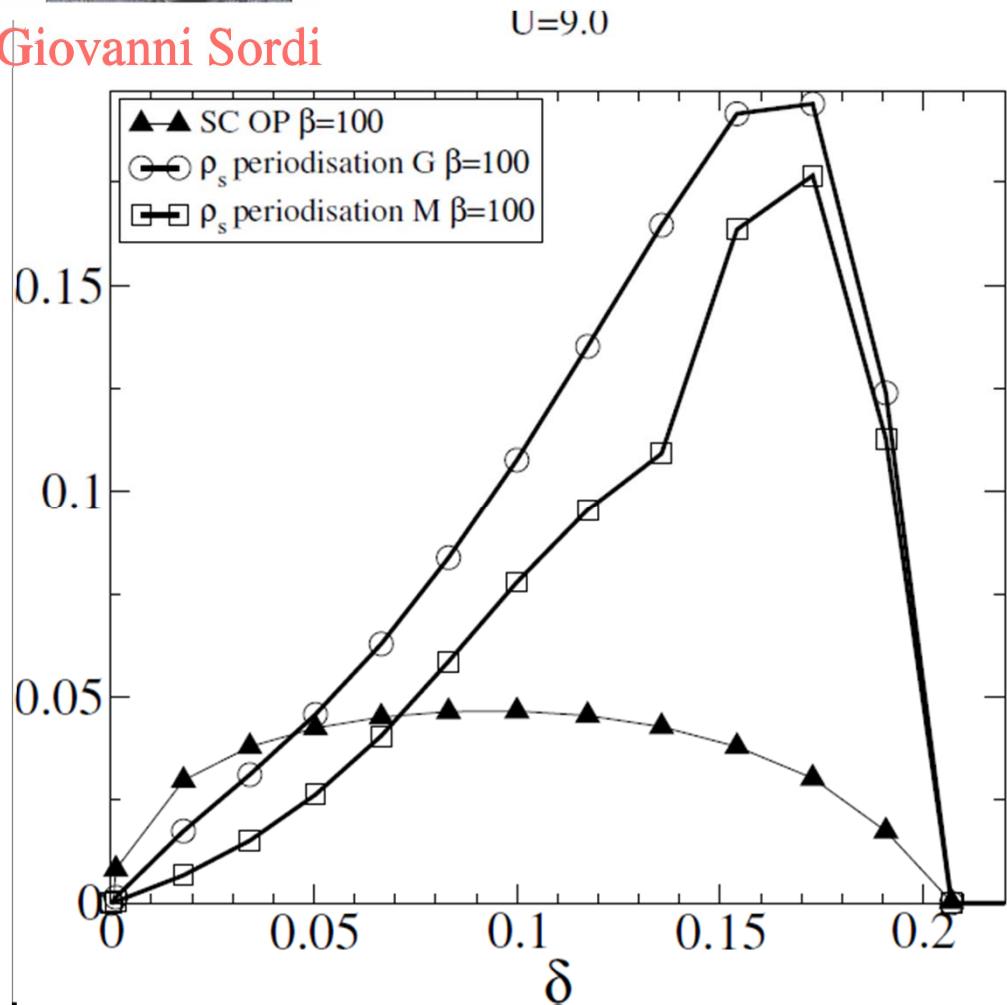




c-axis Superfluid density

$$U = 9t, T = 1/100t$$

Giovanni Sordi



Panagopoulos et al. PRB 2000

See also, Gull Millis, PRB, 2013

Overview

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- The view from dynamical mean-field theory
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Weakly and strongly correlated superconductivity

Analog to weakly and strongly
correlated antiferromagnets



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Weakly and strongly correlated superconductivity

Weakly correlated case



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Weak coupling methods

- Functional renormalization group

$$(a) \partial_\ell \text{ (square loop)} = \text{(square loop with dashed line)} + \text{(square loop with dashed line)}$$

$$+ \text{(square loop with dashed line)} + \dots$$

$$\partial_\ell \text{ (square loop)} = \text{(square loop with dashed line)} + \dots$$

D. Zanchi and H.J. Schulz, PRB 61, 13609 (2000)

C. Honerkamp, et al. PRB 63, 035109 (2001)

R. Shankar, Rev. Mod. Phys. 66, 129 (1994)

C. Bourbonnais Sedeki PRB 2012

$$(b) \partial_\ell \rightarrow \Sigma_+ = \text{(dashed line with cross)} + \text{(dashed line with cross)} + \dots$$

- Other weak coupling methods

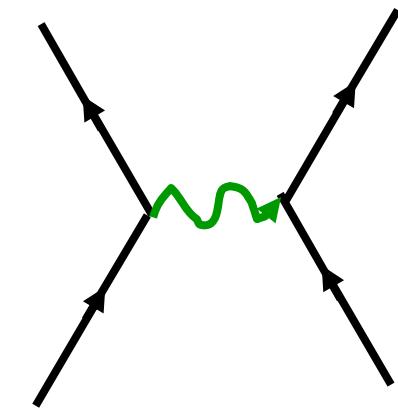
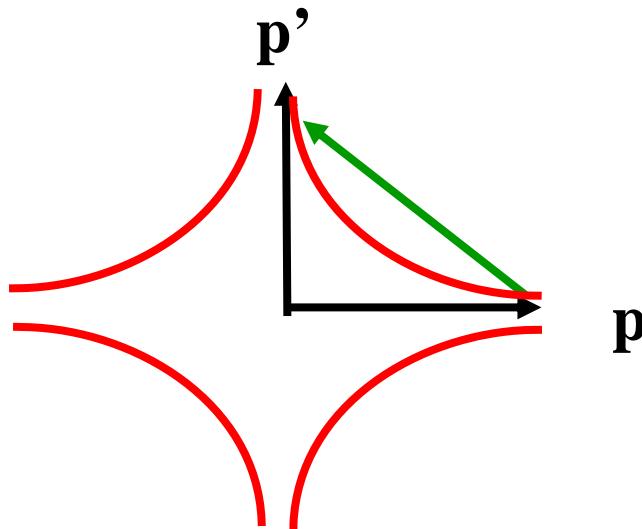
- N.E. Bickers, et al. Phys. Rev. Lett. 62, 961 (1989) FLEX
- B. Kyung, J.-S. Landry, and A.-M.S.T., PRB 68, 174502 (2003)



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Cartoon « BCS » weak-coupling picture

$$\Delta_{\mathbf{p}} = -\frac{1}{2V} \sum_{\mathbf{p}'} U(\mathbf{p} - \mathbf{p}') \frac{\Delta_{\mathbf{p}'}}{E_{\mathbf{p}'}} (1 - 2n(E_{\mathbf{p}'}))$$



Béal–Monod, Bourbonnais, Emery
P.R. B. **34**, 7716 (1986).

Exchange of spin waves?
Kohn-Luttinger
D. J. Scalapino, E. Loh, Jr., and J. E. Hirsch
P.R. B **34**, 8190-8192 (1986).

T_c with pressure

Kohn, Luttinger, P.R.L. **15**, 524 (1965).

P.W. Anderson Science 317, 1705 (2007)

Weakly and strongly correlated superconductivity

Strongly correlated point of view

Many methods



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A cartoon strong coupling picture

$$J \sum_{\langle i,j \rangle} \mathbf{S}_i \cdot \mathbf{S}_j = J \sum_{\langle i,j \rangle} \left(\frac{1}{2} c_i^\dagger \vec{\sigma} c_i \right) \cdot \left(\frac{1}{2} c_j^\dagger \vec{\sigma} c_j \right)$$

$$d = \langle \hat{d} \rangle = 1/N \sum_{\vec{k}} (\cos k_x - \cos k_y) \langle c_{\vec{k},\uparrow}^\dagger c_{-\vec{k},\downarrow} \rangle$$

$$H_{MF} = \sum_{\vec{k},\sigma} \varepsilon(\vec{k}) c_{\vec{k},\sigma}^\dagger c_{\vec{k},\sigma} - 4Jm\hat{m} - Jd(\hat{d} + \hat{d}^\dagger) + F_0$$

Pitaevskii Brückner:

Pair state orthogonal to repulsive core of Coulomb interaction

P.W. Anderson Science
317, 1705 (2007)

Miyake, Schmitt–Rink, and Varma
P.R. B 34, 6554-6556 (1986)

More sophisticated Slave Boson: Kotliar Liu PRB 1988



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A paradox (Scalapino)

Resilience to near-neighbor repulsion V

$$\hat{\mathcal{H}}_{Hubbard} = - \sum_{\langle i,j \rangle_{1,2,3}} \left(t_{ij} \hat{c}_{i\sigma}^\dagger \hat{c}_{j\sigma} + c.h \right) + U \sum_i \hat{n}_{i\uparrow} \hat{n}_{i\downarrow} + V \sum_{\langle i,j \rangle} \hat{n}_i \hat{n}_j - \mu \sum_i \hat{n}_{i\sigma}$$

YBa₂Cu₃O₇ : $t = 1$ $t' = -0.3$ $t'' = 0.2$

We expect superconductivity to disappear when:

$V > \frac{U^2}{W}$ **In weakly correlated case**
 $U/W < 1$

$V > J$ **In mean-field strongly correlated case**

In cuprates:

$$V = 400 \text{ meV}$$

$$J = 130 \text{ meV}$$

$$U_c = V_c / [1 + N(0)V_c \ln(E_F/\omega_c)]$$

S. Onari, R. Arita, K. Kuroki et H. Aoki, PRB **70**, 094523 (2004)
S. Raghu, E. Berg, A. V. Chubukov et S. A. Kivelson, PRB **85**,

024516 (2012)
S. Sorella, et al. Phys. Rev. Lett. **88**, 117002 (2002)

Overview

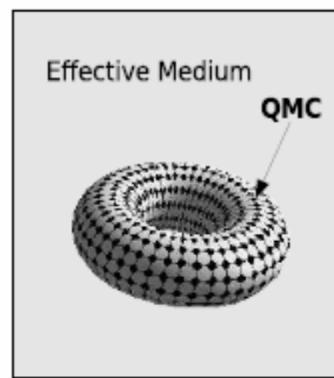
- An appetizer:
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Strongly correlated superconductivity: the view from dynamical mean-field theory

2d Hubbard: Quantum cluster method

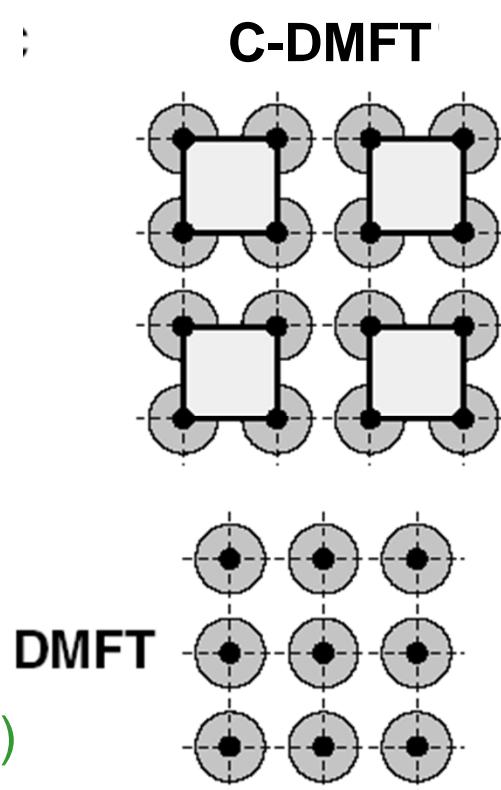


DCA

Hettler ... Jarrell ... Krishnamurty PRB **58** (1998)

Kotliar et al. PRL **87** (2001)

M. Potthoff et al. PRL **91**, 206402 (2003).



REVIEWS

Maier, Jarrell et al., RMP. (2005)

Kotliar et al. RMP (2006)

AMST et al. LTP (2006)



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+ and -

- Long range order:
 - Allow symmetry breaking in the bath (mean-field)
- Included:
 - Short-range dynamical and spatial correlations
- Missing:
 - Long wavelength p-h and p-p fluctuations



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$T = 0$ results

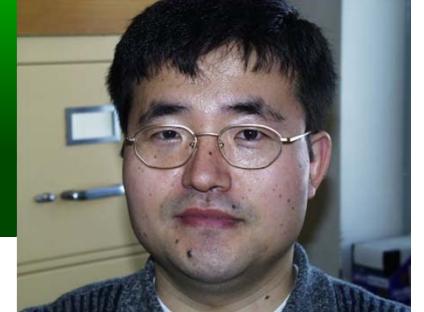


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Competition AFM-dSC



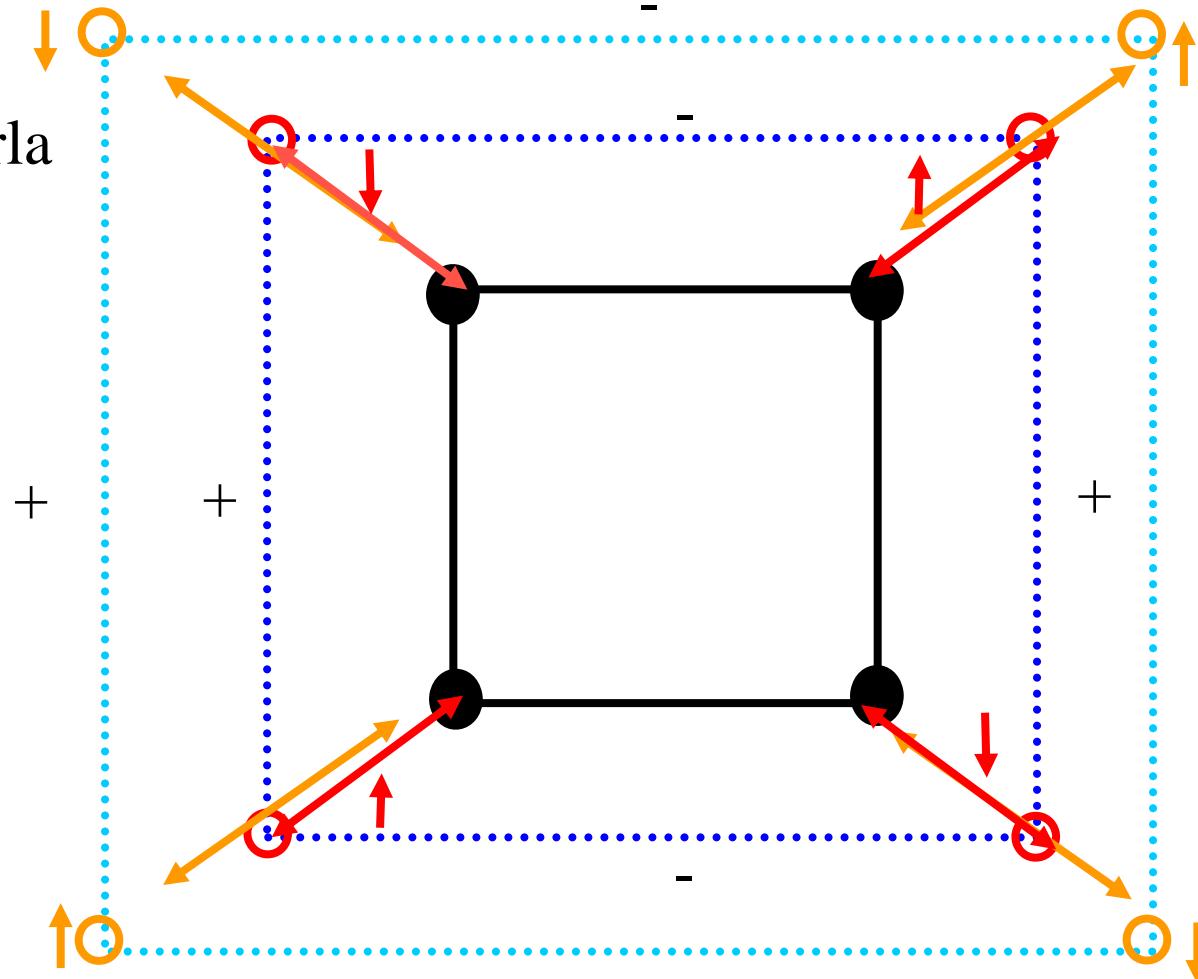
S. Kancharla



B. Kyung



David Sénéchal



See also, Capone and Kotliar, Phys. Rev. B 74, 054513 (2006),
Macridin, Maier, Jarrell, Sawatzky, Phys. Rev. B 71, 134527 (2005)



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Pairing, glue and glue remover

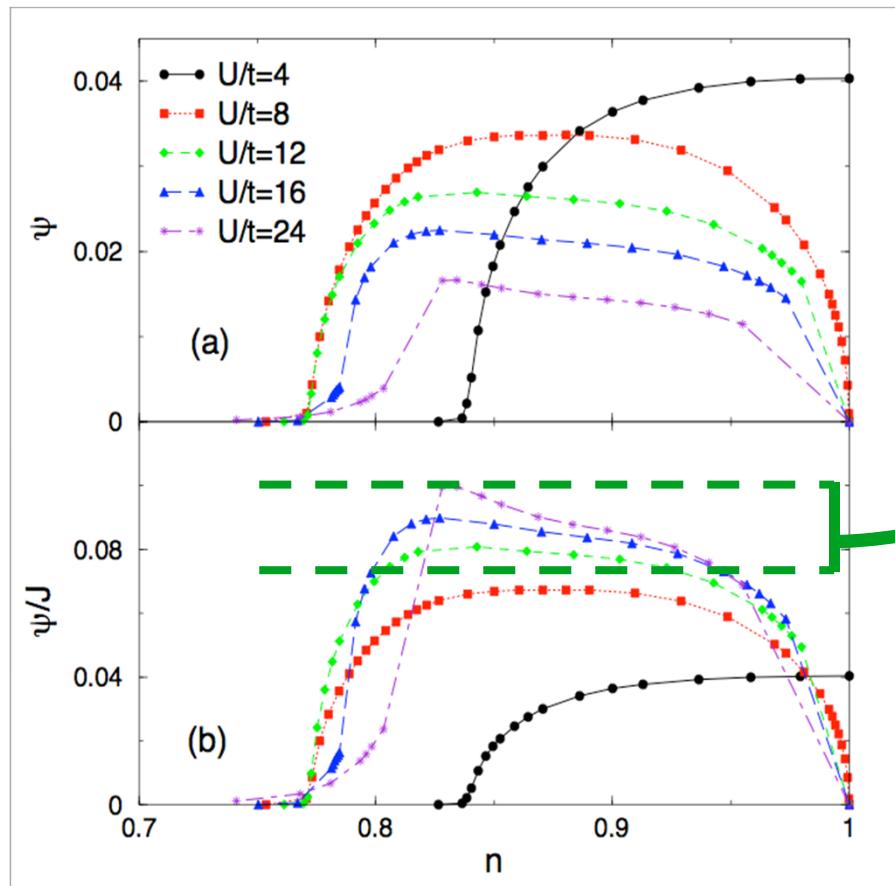
Strength of pairing

Kancharla, Kyung, Civelli,
Sénéchal, Kotliar AMST PRB (2008)

In cuprates:

$V = 400 \text{ meV}$

$J = 130 \text{ meV}$



The
superconducting
order parameter
scales like J



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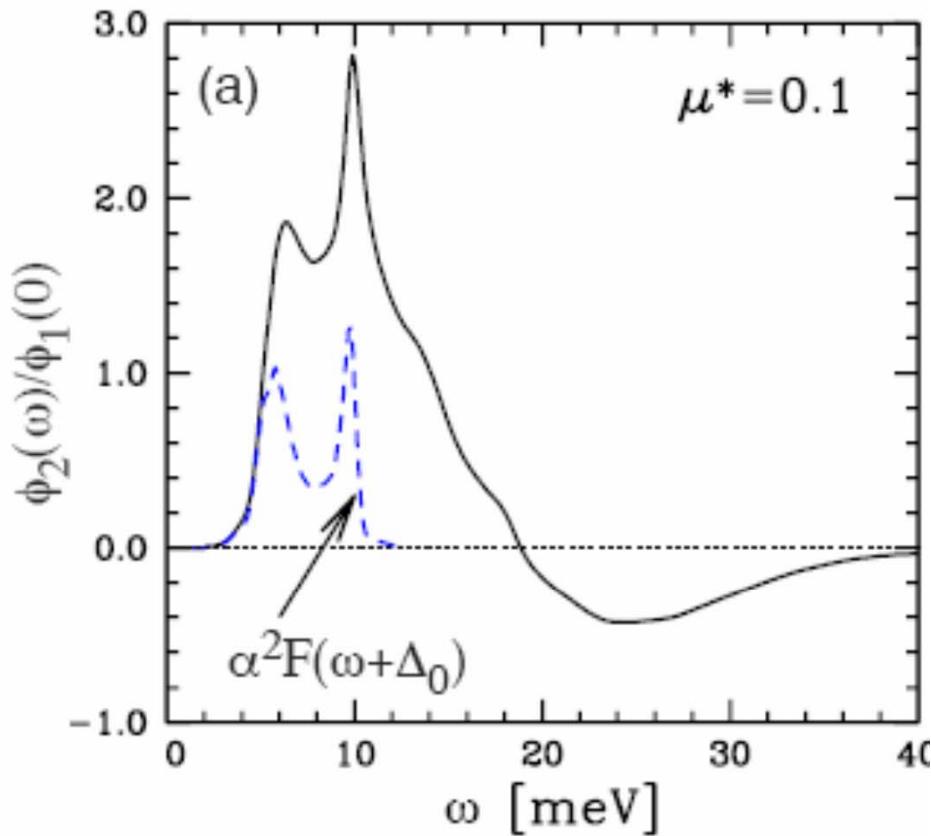
The glue



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$\text{Im } \Sigma_{\text{an}}$ and electron-phonon in Pb

Maier, Poilblanc, Scalapino, PRL (2008)



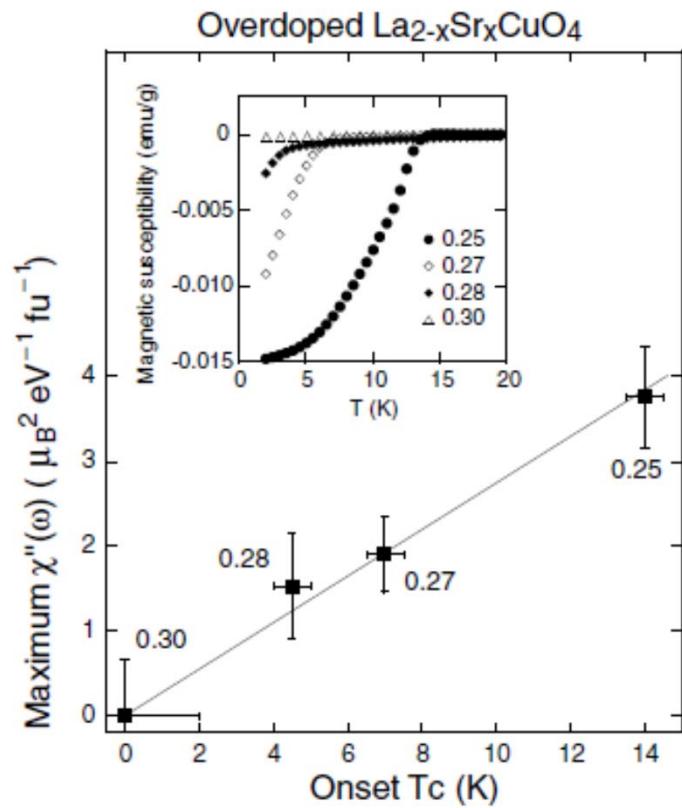
Retardation

$$V_{el-ph}^{eff}(\vec{q}, \omega) = \frac{e^2}{4\pi\epsilon_0(q^2 + k_{TF}^2)} \left[1 + \frac{\omega_{ph}^2(\vec{q})}{\omega^2 - \omega_{ph}^2(\vec{q})} \right]$$

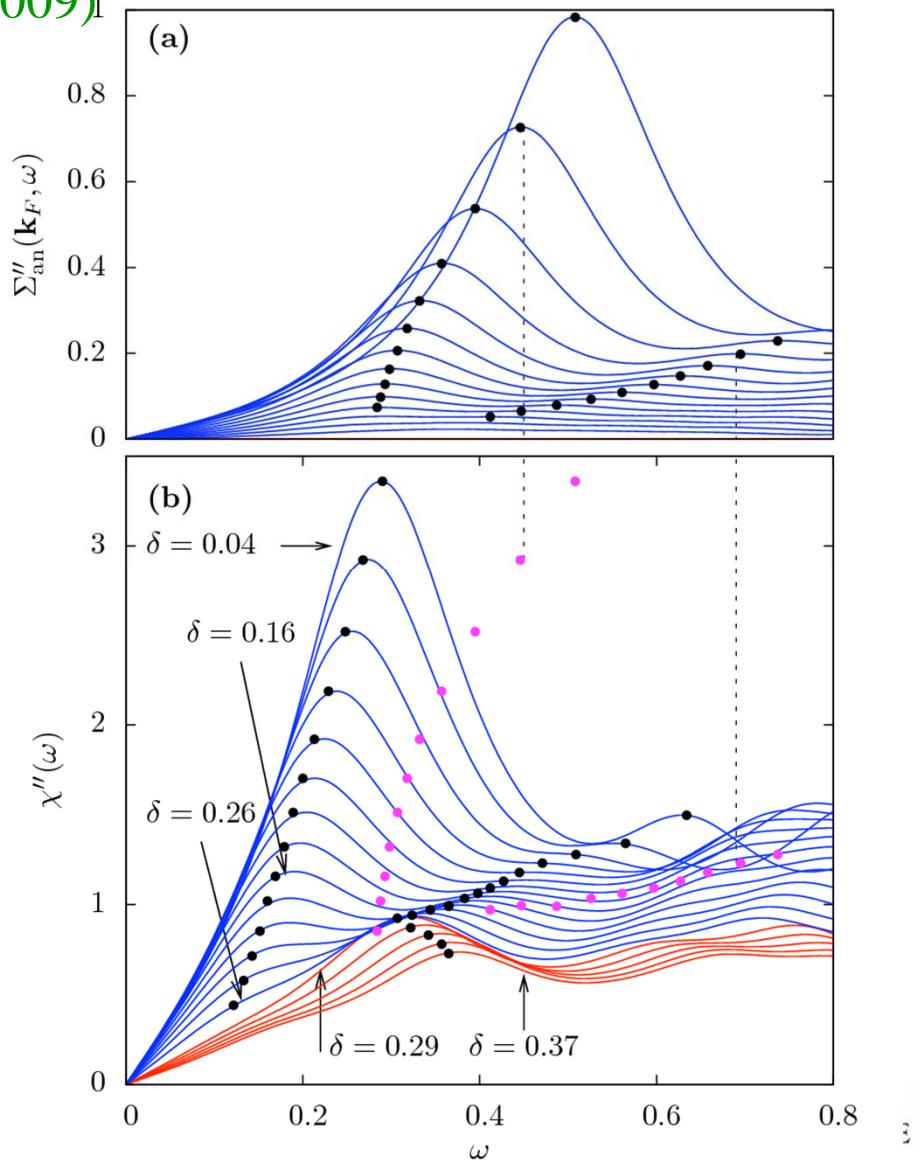


The glue for strong corrections

Kyung, Sénéchal, AMST PRB **80**, 205109 (2009)
Civelli, PRL + PRB **79** 195113 (2009)
Gull and Millis arXiv:1407.0704



Wakimoto ... Birgeneau
PRL (2004)



The glue and neutrons

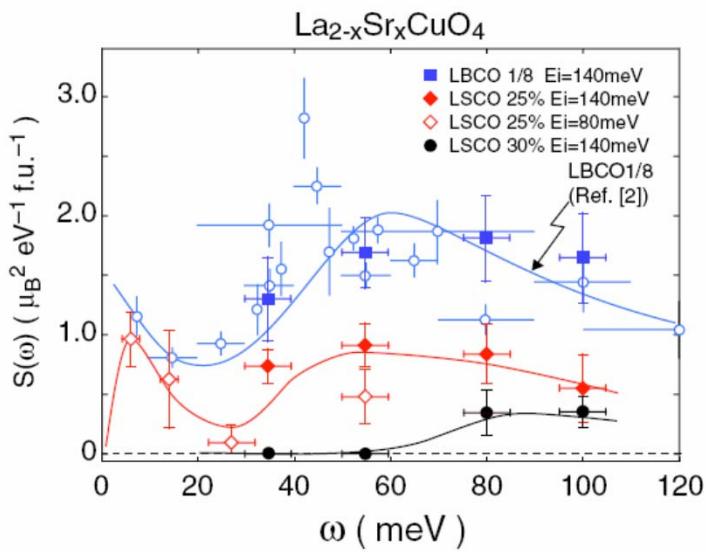
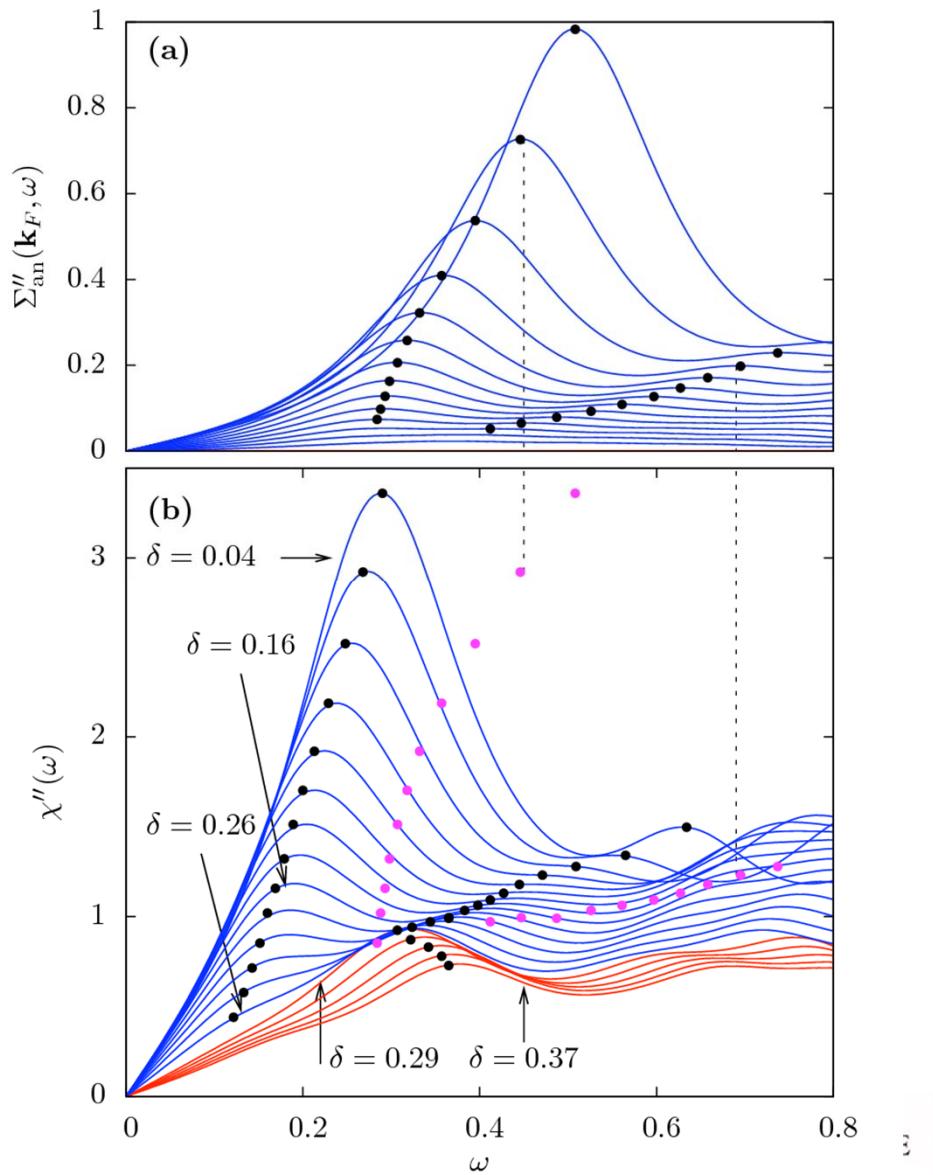


FIG. 3 (color online). \mathbf{Q} -integrated dynamic structure factor $S(\omega)$ which is derived from the wide- H integrated profiles for LBCO 1/8 (squares), LSCO $x = 0.25$ (diamonds; filled for $E_i = 140$ meV, open for $E_i = 80$ meV), and $x = 0.30$ (filled circles) plotted over $S(\omega)$ for LBCO 1/8 (open circles) from [2]. The solid lines following data of LSCO $x = 0.25$ and 0.30 are guides to the eyes.

Wakimoto ... Birgeneau PRL (2007);
 PRL (2004)

Kyung, Sénéchal, AMST PRB **80**, 205109
 (2009); Haule Kotliar (2007)





Frequencies important for pairing



Bumsoo Kyung

Anomalous Green function

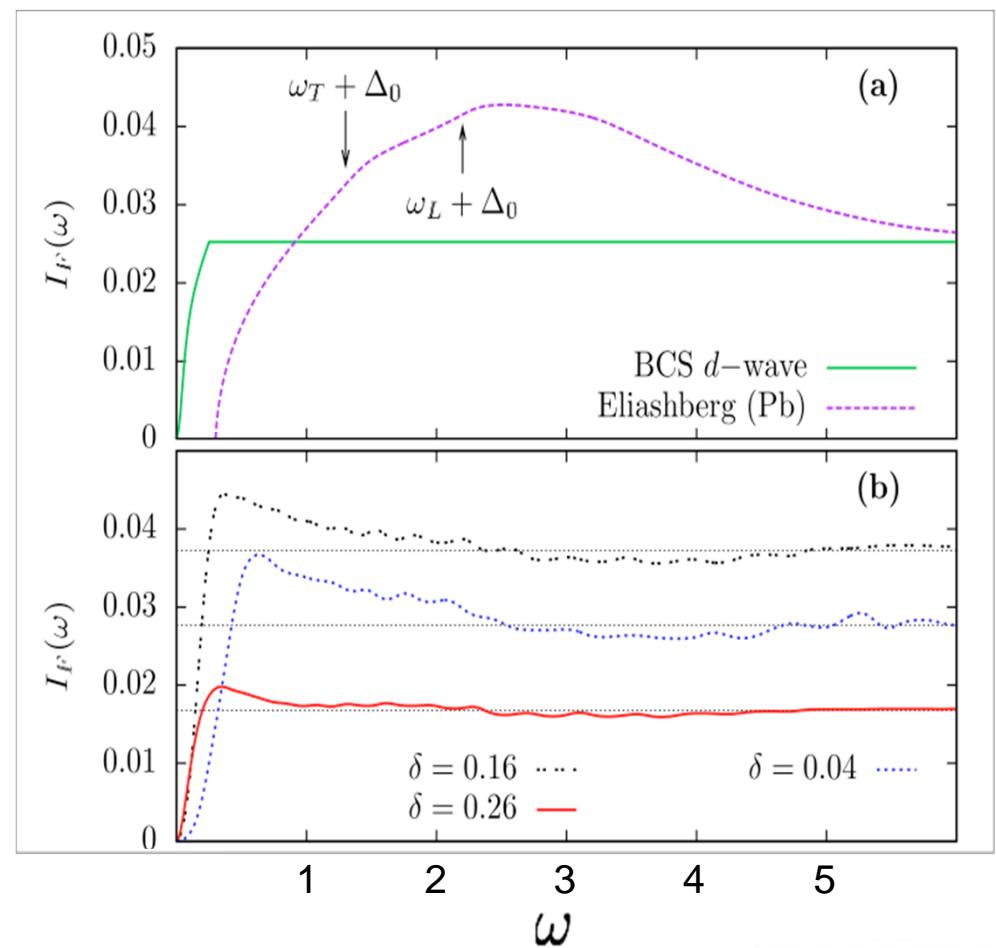
$$[\mathcal{F}_{an}(t)]_{lm} = -i\theta(t) \langle \{\hat{c}_{l\uparrow}(t), \hat{c}_{m\downarrow}(0)\} \rangle_{\mathcal{H}_{AIM}}$$

Anomalous spectral function

$$[\mathcal{A}_{an}(\omega)]_{lm} = -\frac{1}{\pi} \text{Im} [\mathcal{F}_{an}(\omega)]_{lm}$$

Cumulative order parameter:

$$I_{\mathcal{F}}(\omega) = - \int_0^{\omega} \frac{d\omega'}{\pi} \text{Im} [\mathcal{F}_{an}(\omega')]_{lm}$$
$$I_{\mathcal{F}}(\omega) \xrightarrow[\omega \rightarrow +\infty]{} \langle \hat{c}_{l\uparrow} \hat{c}_{m\downarrow} \rangle_{\mathcal{H}_{AIM}}$$



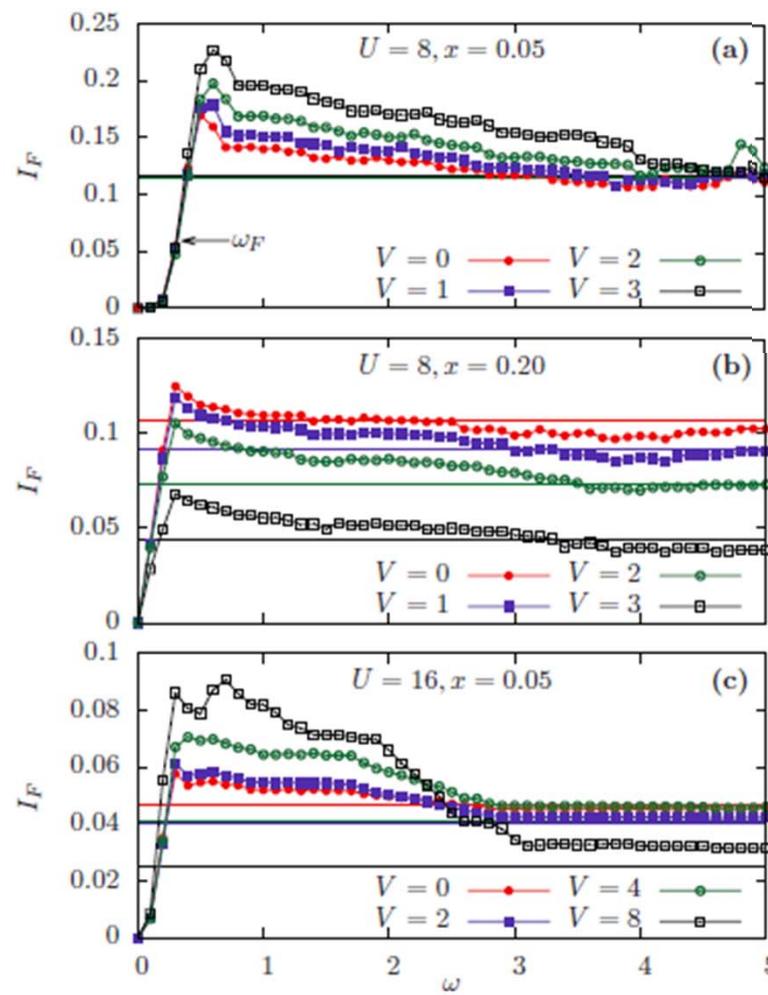
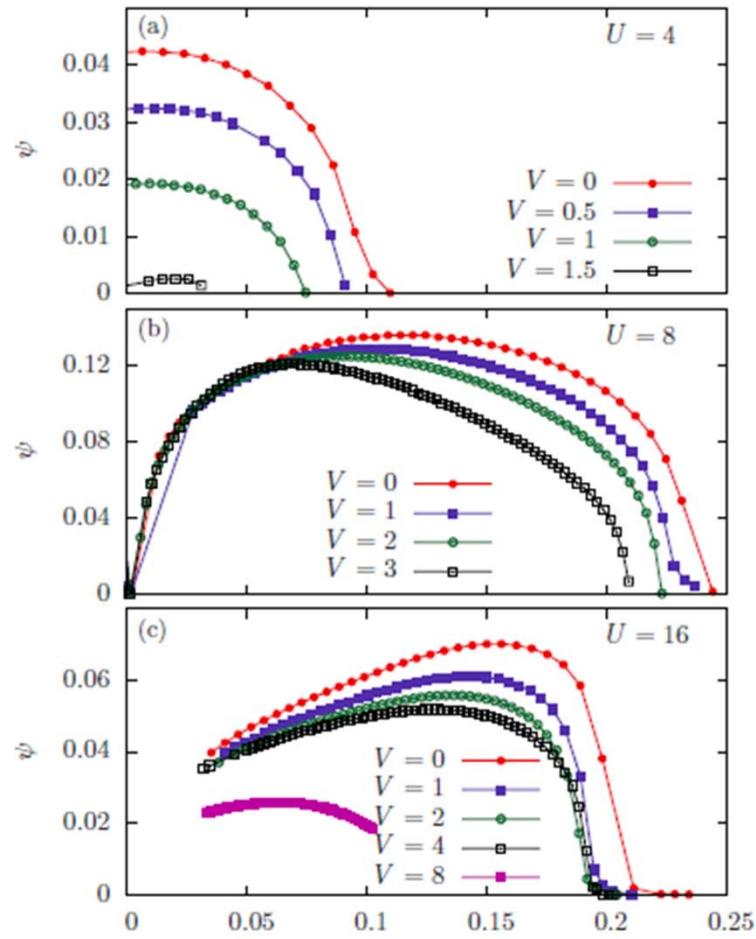


Resilience to near-neighbor repulsion



David Sénéchal

Alexandre Day



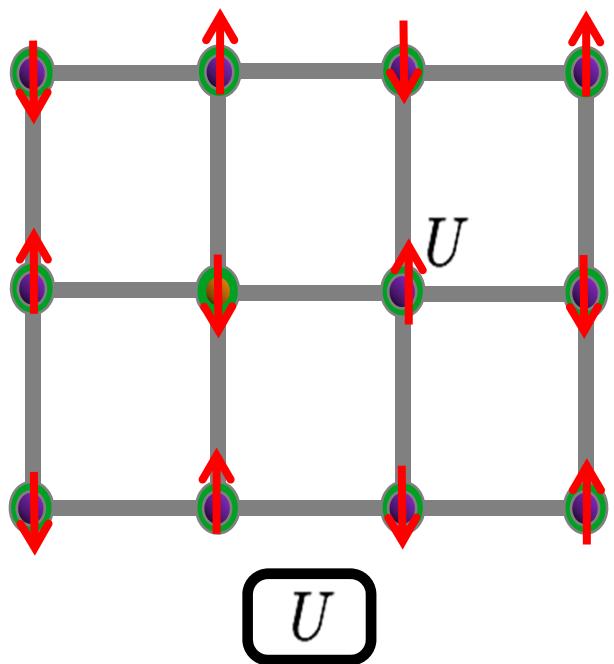
Vincent Bouliane

Sénéchal, Day, Bouliane, AMST PRB **87**, 075123 (2013)

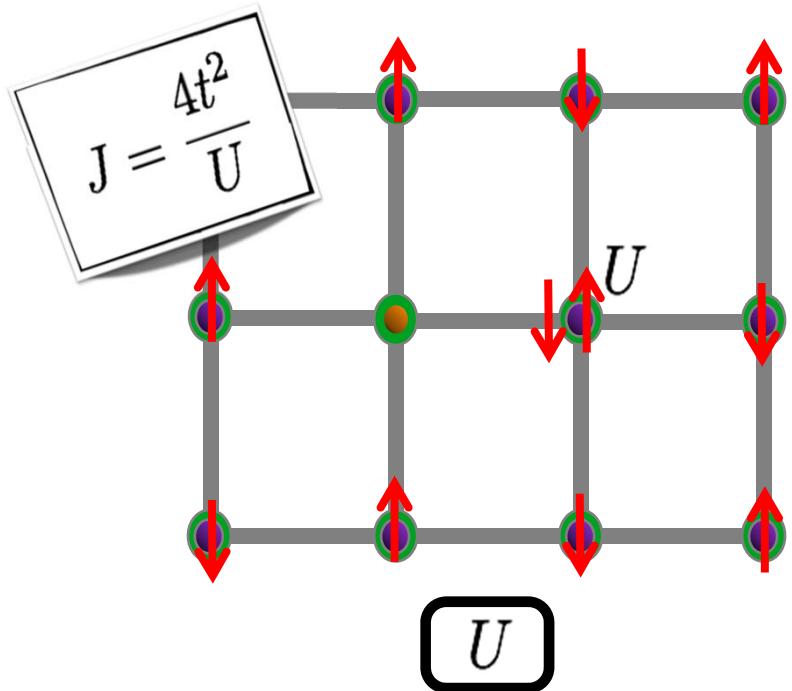


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V also increases J

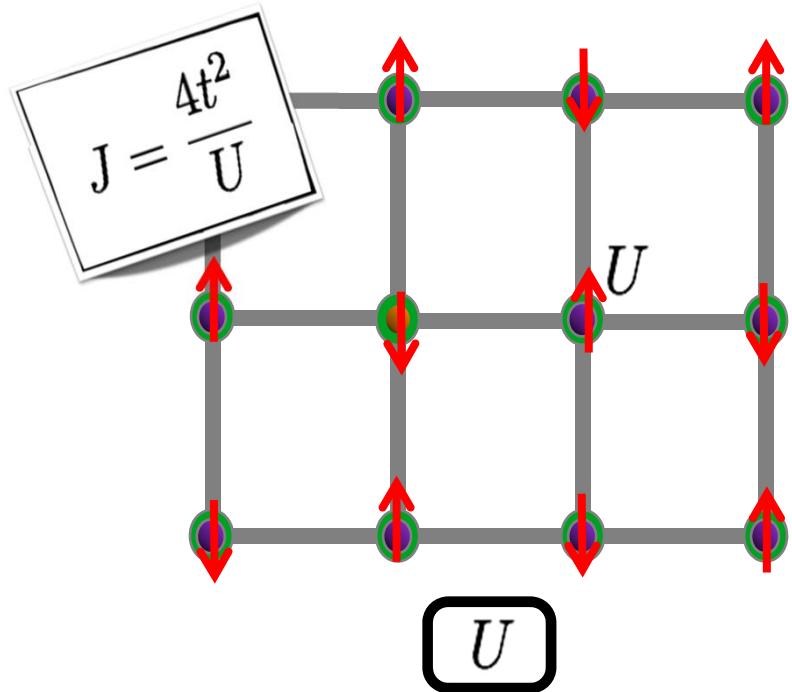


V also increases J

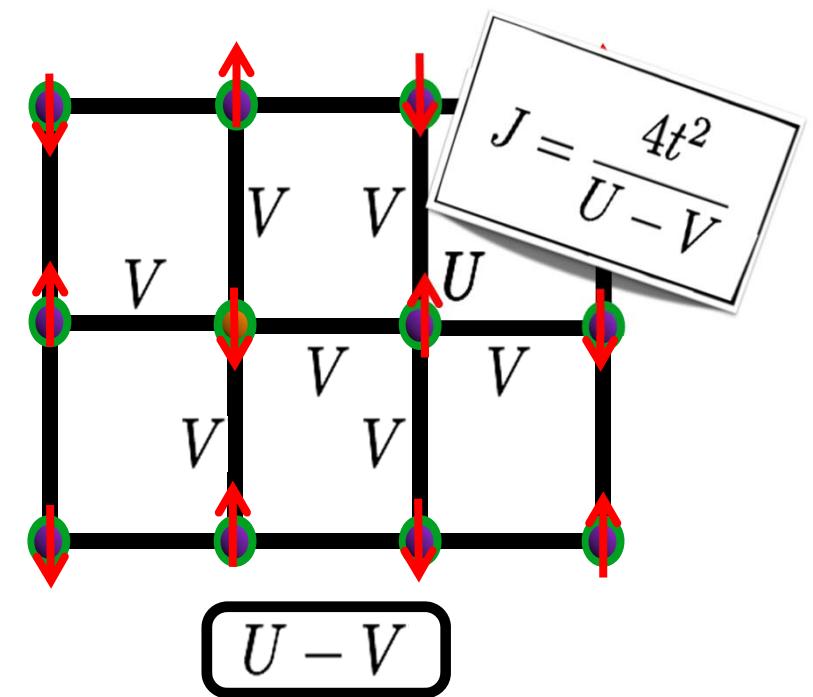
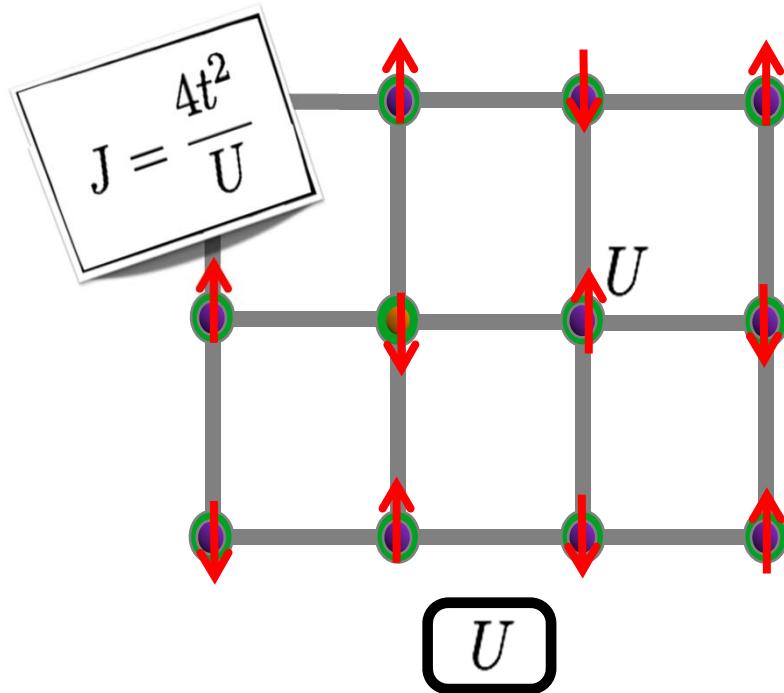


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V also increases J

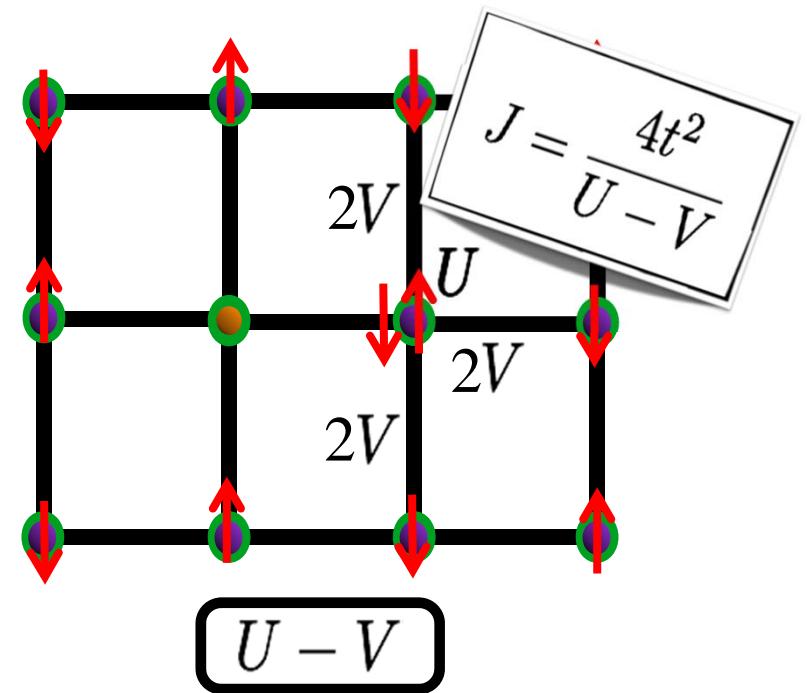
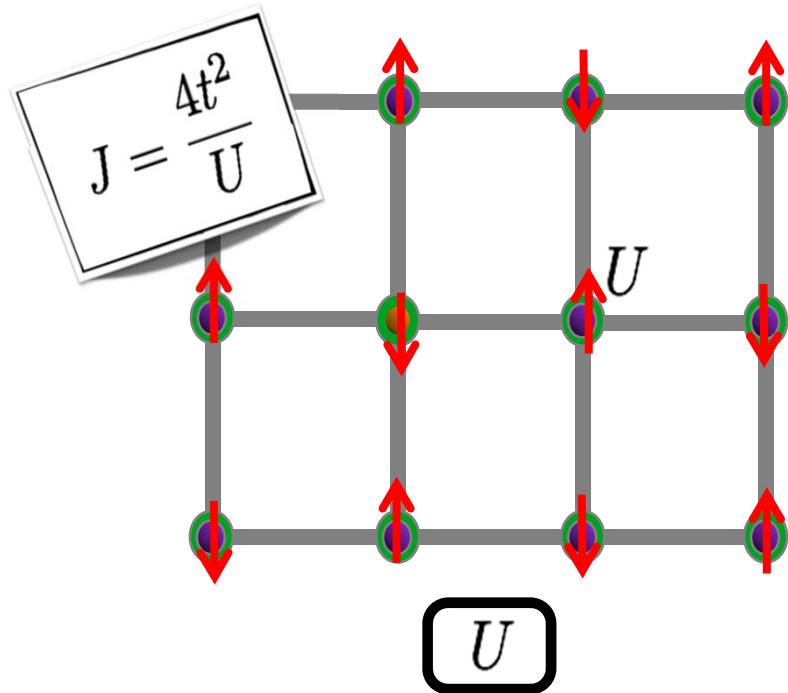


V also increases J



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V also increases J



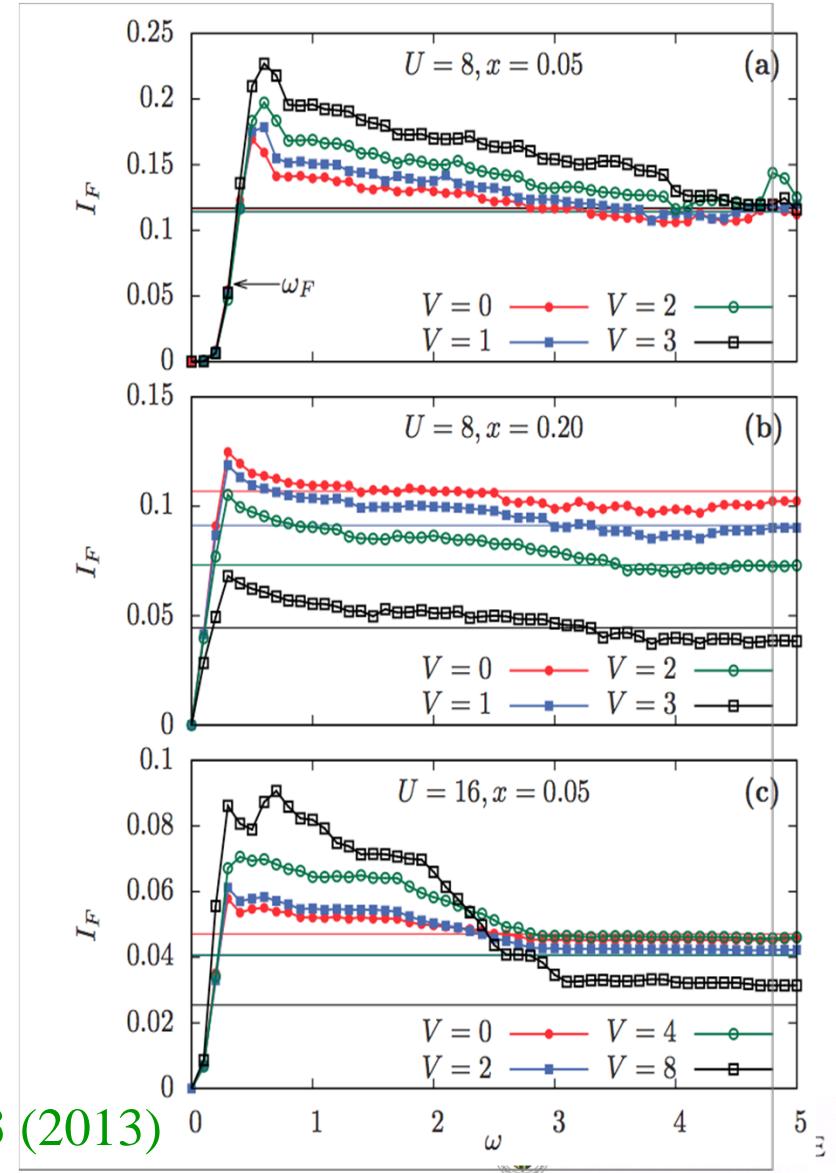
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Binding aspects of V

$$J = \frac{4t^2}{U - V}$$

**J increases with V
explaining better pairing at
low frequency**

**But V also induces more
repulsion at high frequency,
explaining the negative
impact at high frequency on
binding**



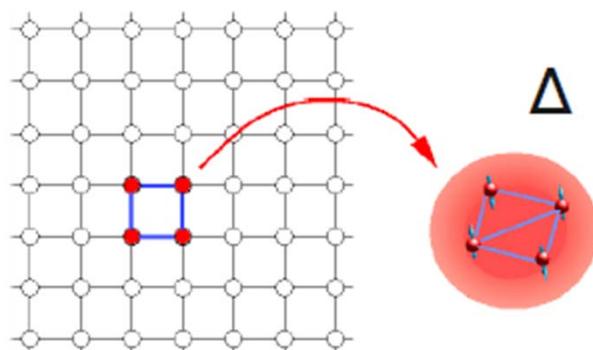
Finite T and more experiments

H. Alloul C.R. Physique, **15** (2014)

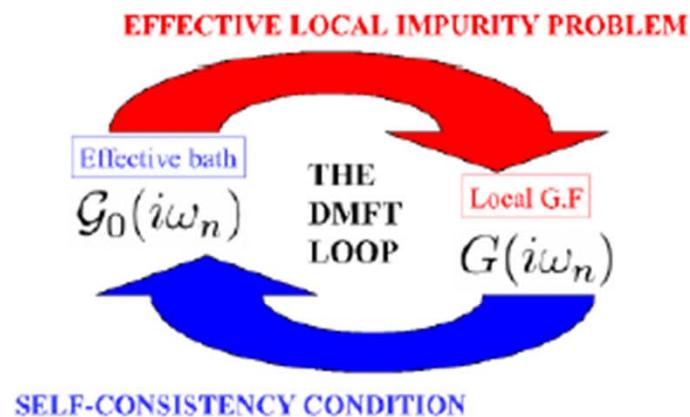


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C-DMFT



$$Z = \int \mathcal{D}[\psi^\dagger, \psi] e^{-S_c - \int_0^\beta d\tau \int_0^\beta d\tau' \sum_{\mathbf{k}} \psi_{\mathbf{k}}^\dagger(\tau) \Delta(\tau, \tau') \psi_{\mathbf{k}}(\tau')}$$



Mean-field is not a trivial problem! Many impurity solvers.

Here: continuous time QMC

-
- P. Werner, PRL 2006
 - P. Werner, PRB 2007
 - K. Haule, PRB 2007

$$\Delta(i\omega_n) = i\omega_n + \mu - \Sigma_c(i\omega_n)$$

$$- \left[\sum_{\tilde{k}} \frac{1}{i\omega_n + \mu - t_c(\tilde{k}) - \Sigma_c(i\omega_n)} \right]^{-1}$$



Giovanni Sordi



Patrick Sémon

Standard 2 operator updates in CTQMC
are not ergodic in SC state (fixed here)

Please ask in question period...

P. Sémon, G. Sordi, A.-M.S. Tremblay, Phys. Rev. B **89**, 165113/1-6 (2014)



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Superconductivity at $U > U_{MIT}$



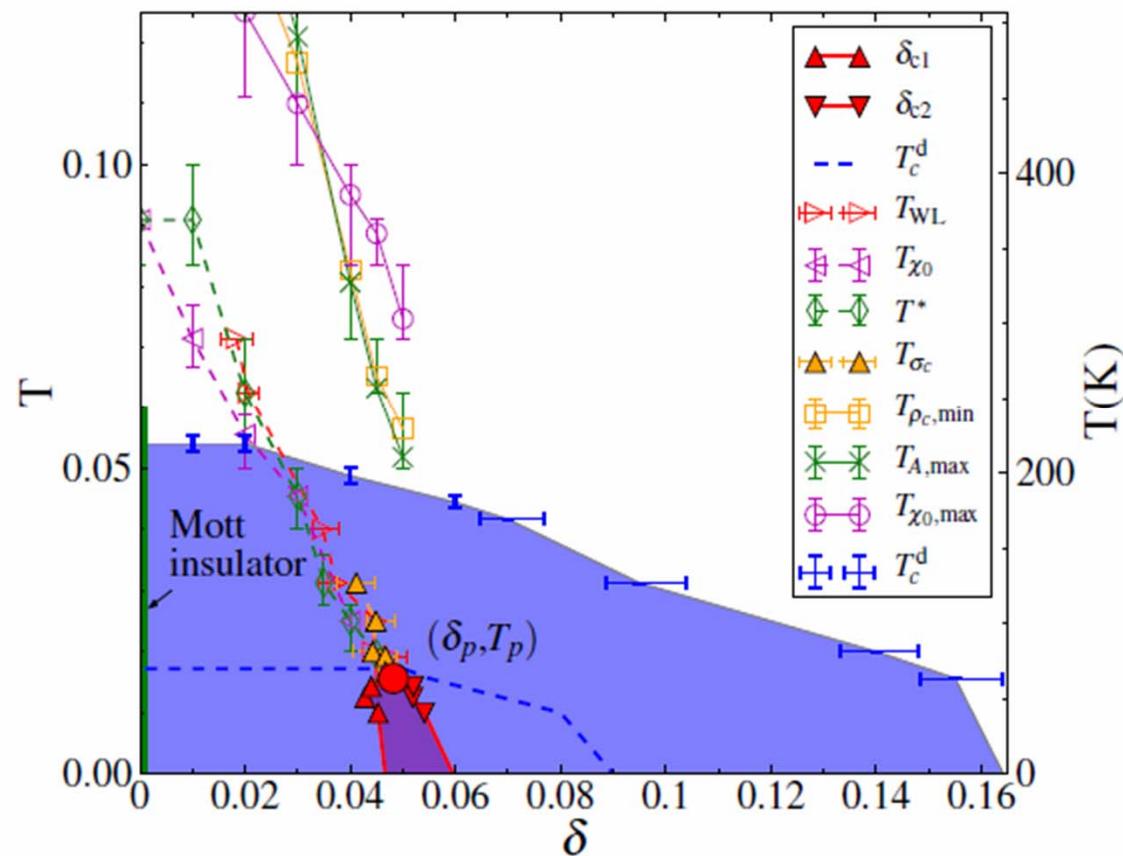
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Giovanni Sordi



Patrick Sémon



G. Sordi et al. Phys. Rev. Lett. 108, 216401/1-6 (2012)

P. Sémon, G. Sordi, A.-M.S.T., Phys. Rev. B **89**, 165113/1-6 (2014)

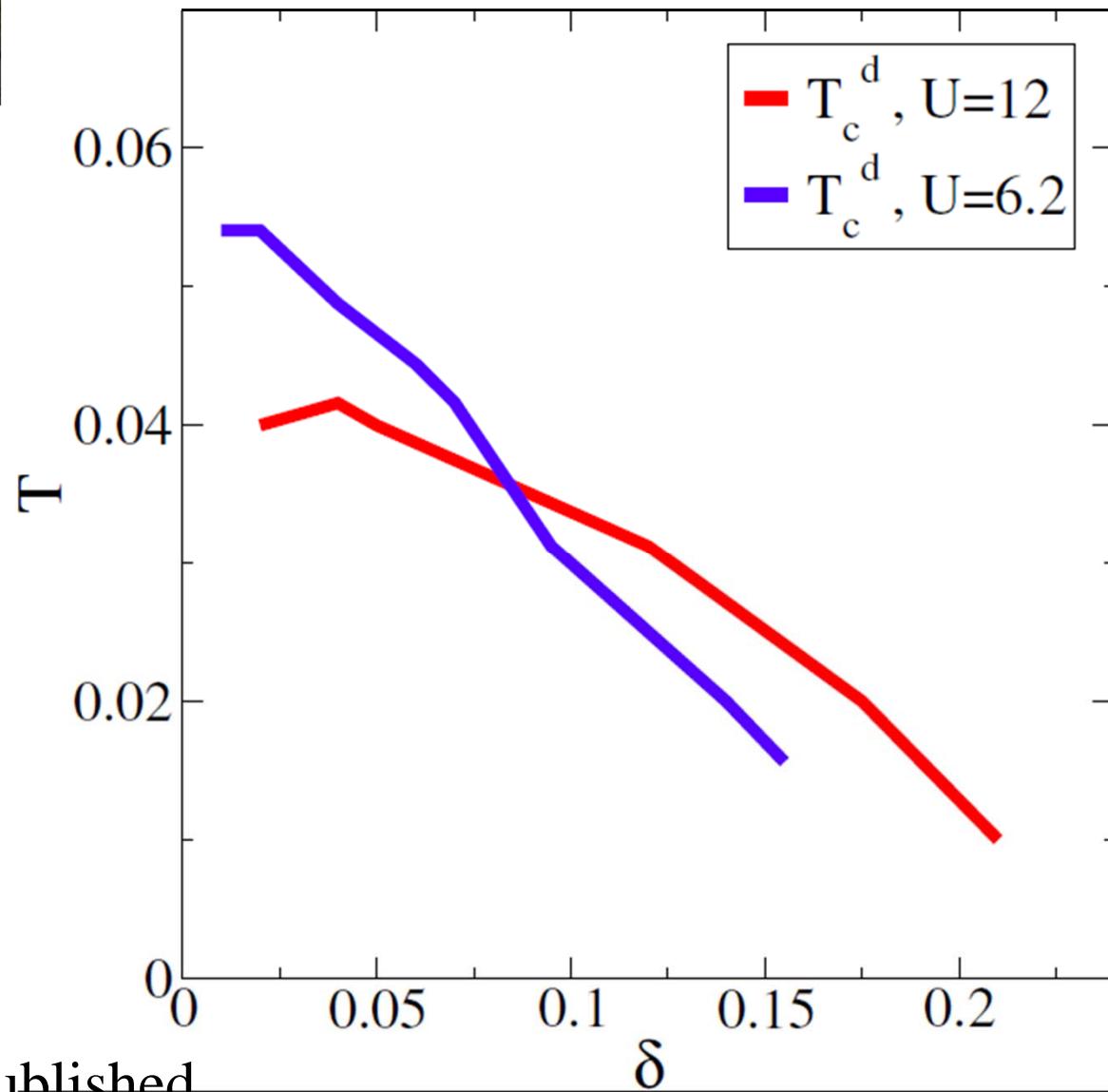


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T_c^d

Giovanni Sordi



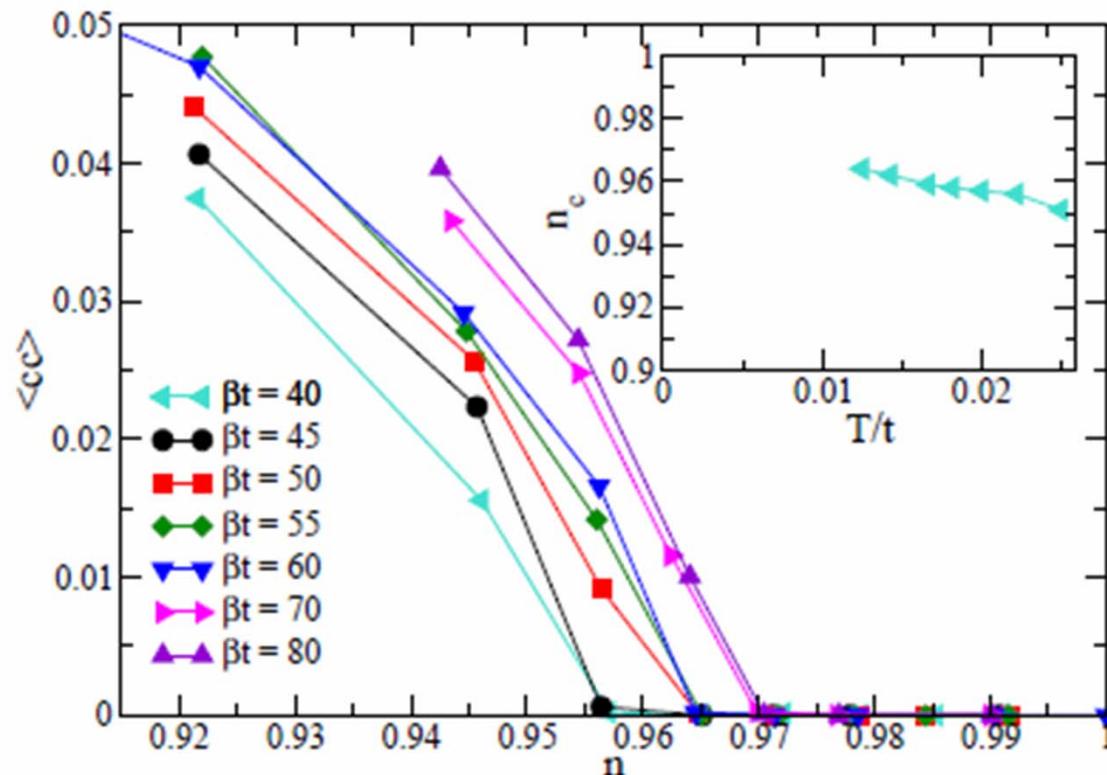
G. Sordi Unpublished



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8 site cluster

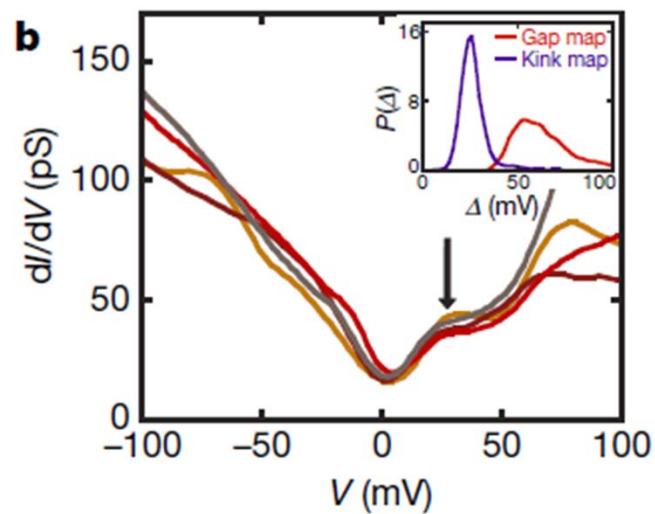
Ends at finite doping
Gull, Parcollet, Millis, PRL 110, 216405 (2013)



Gull, Parcollet, Millis PRL 110, 216405 (2013)

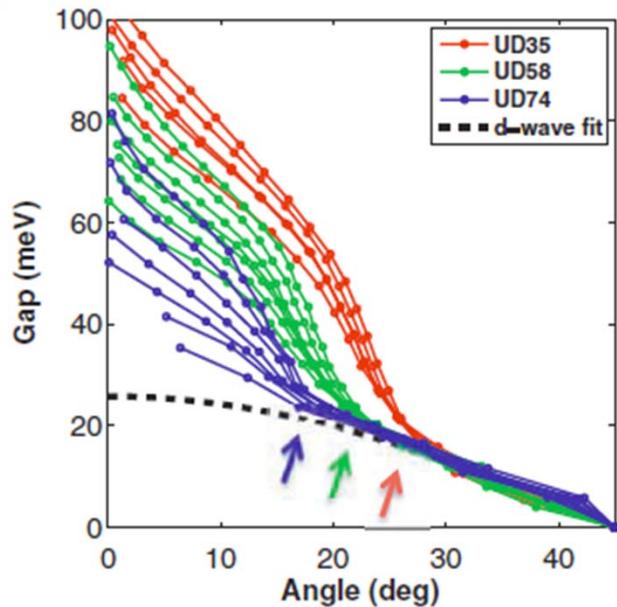
Meaning of T_c^d

- Local pair formation



K. K. Gomes, A. N. Pasupathy, A. Pushp,
S. Ono, Y. Ando, and A. Yazdani,
Nature **447**, 569 (2007)

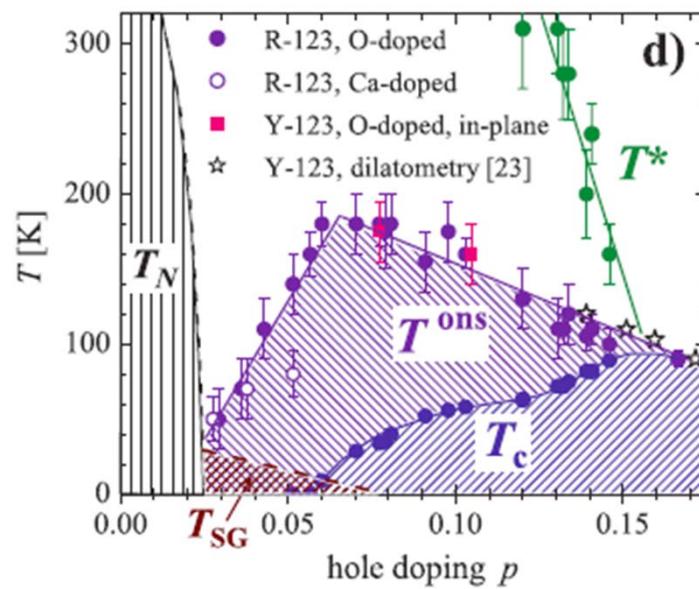
Meaning of T_c^d : Local pair formation



A. Pushp, Parker, ... A. Yazdani,
Science **364**, 1689 (2009)

However, our measurements demonstrate that the nodal gap does not change with reduced doping. The pairing strength does not get weaker or stronger as the Mott insulator is approached; rather, it saturates.

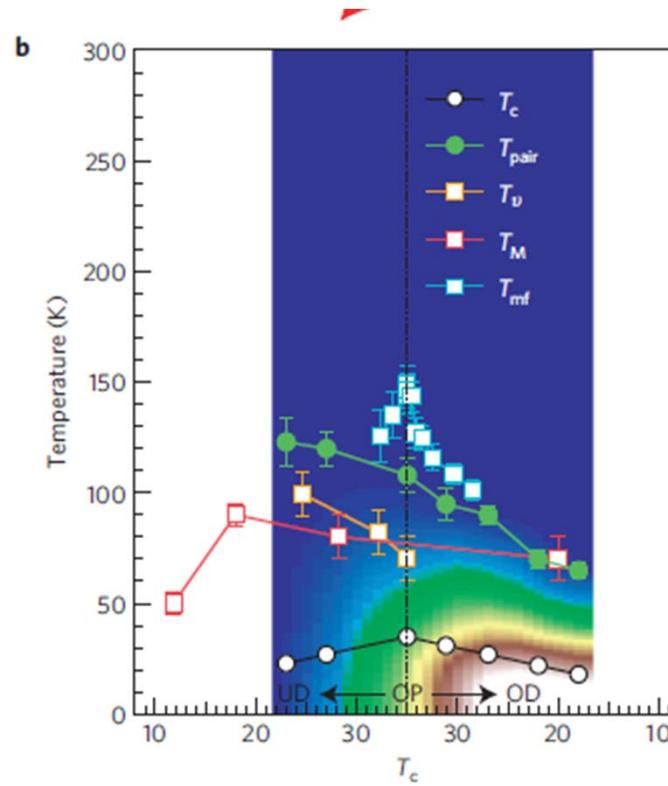
Fluctuating region



Infrared response

Dubroka et al. 106, 047006 (2011)

T_{pair}



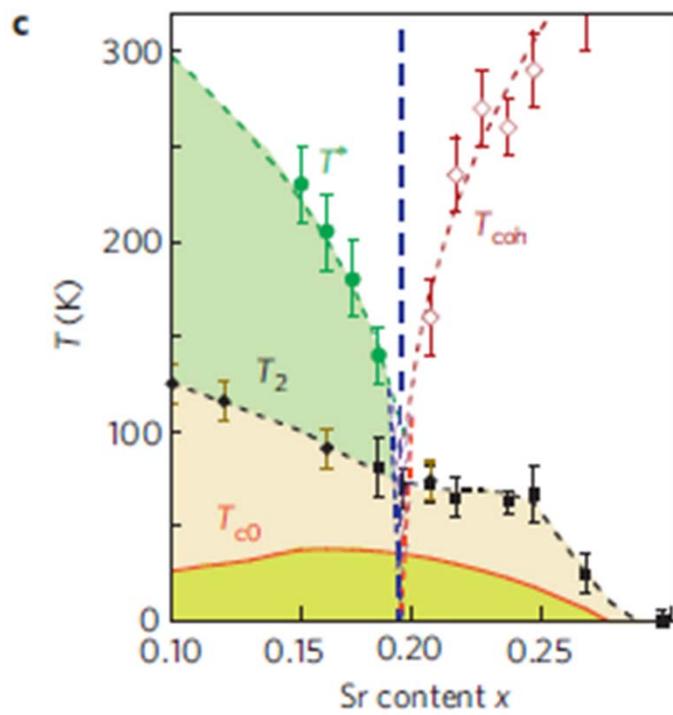
ARPES
Bi2212

Kondo, Takeshi, et al. Kaminski Nature
Physics **2011**, 7, 21-25



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T_2



Magnetoresistance, LSCO
Fluctuating vortices

Patrick M. Rourke, et al. Hussey Nature Physics 7, 455–458 (2011)



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Giant proximity effect

$$\begin{aligned}T_c &= 32 \text{ K} \\T_c &< 5 \text{ K}\end{aligned}$$

Morenzoni et al.,
Nature Comms. 2 (2011)

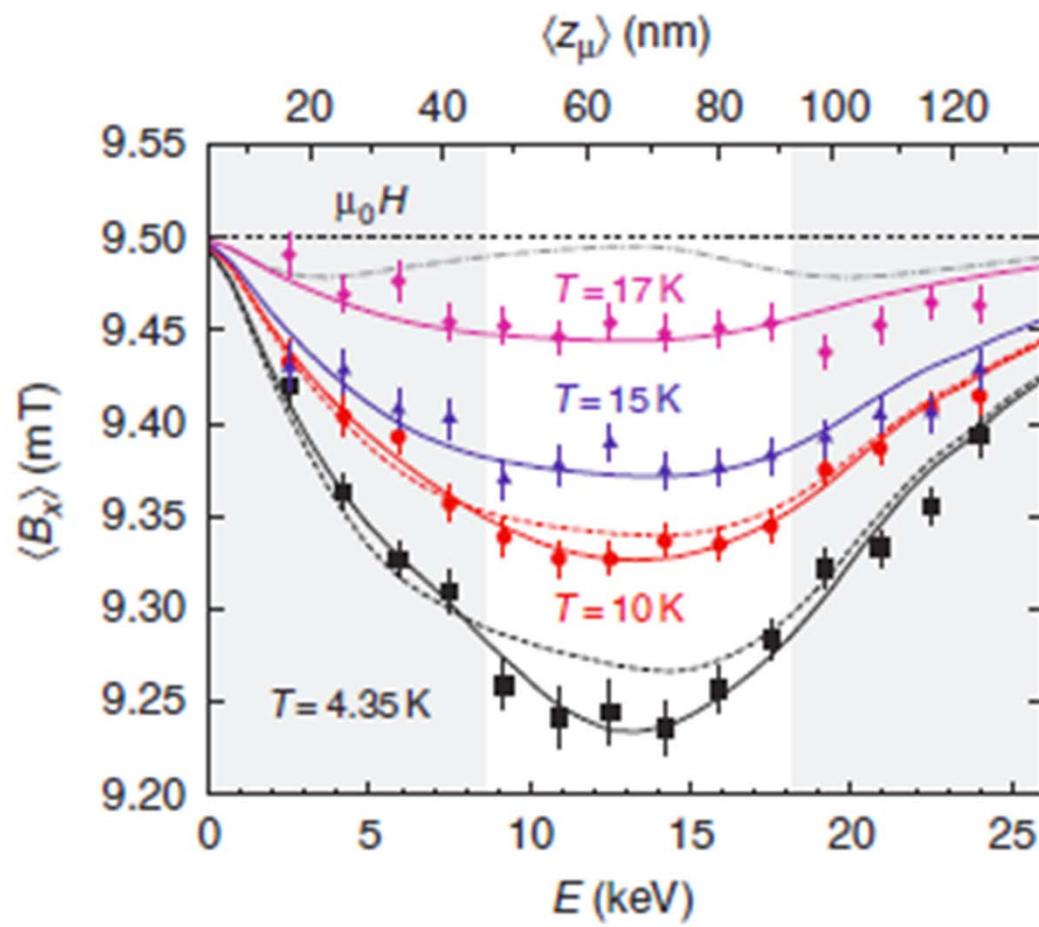
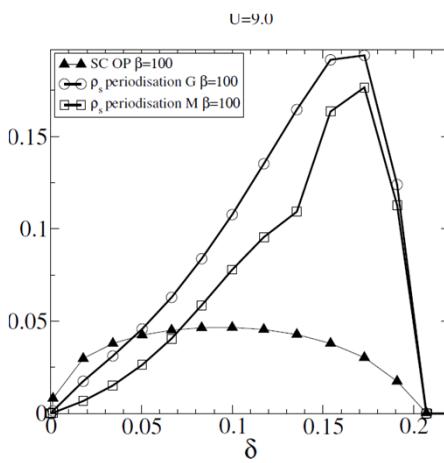
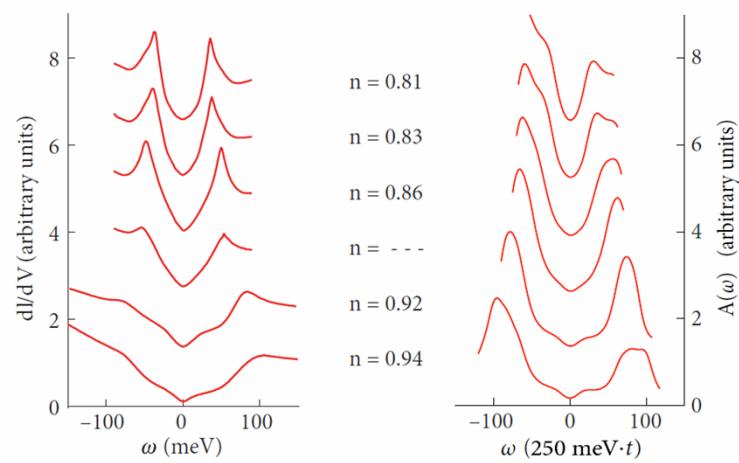
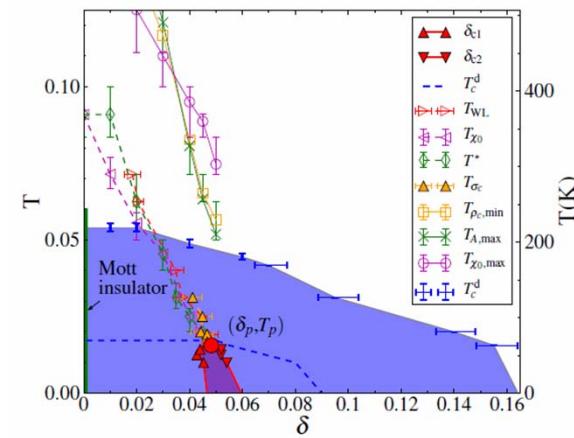
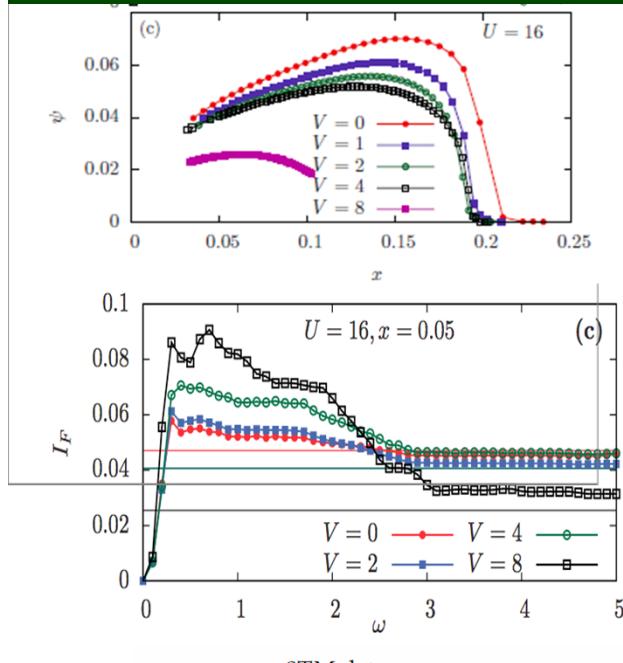


Figure 6 | Depth profile of the local field at different temperatures. The

Summary



Collaborators



Giovanni Sordi



David Sénéchal



Alexandre Day



Vincent Bouliane



Patrick Sémon



Kristjan Haule



Simon Verret

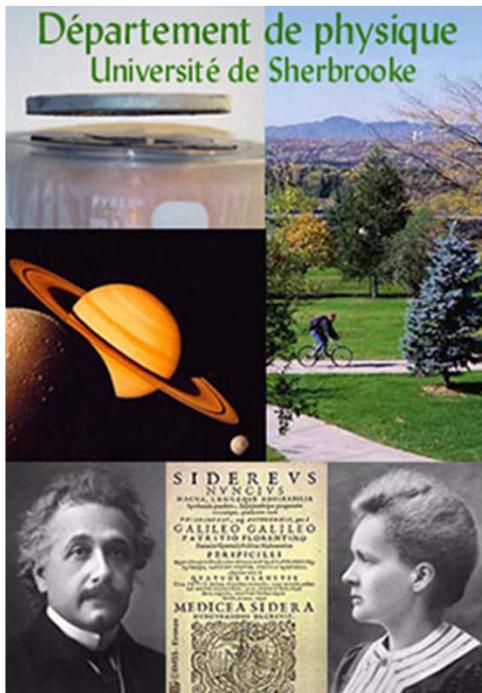


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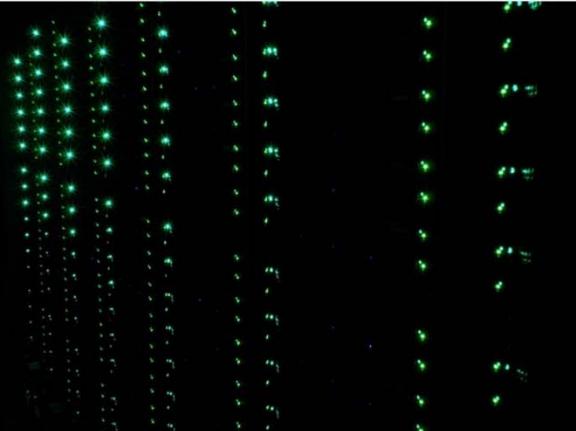
Le regroupement québécois sur les matériaux de pointe



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Review: A.-M.S.T. arXiv: 1310.1481



A.-M.S. Tremblay

“Strongly correlated superconductivity”

Chapt. 10 : *Emergent Phenomena in Correlated Matter Modeling and Simulation*, Vol. 3, E. Pavarini, E. Koch, and U. Schollwöck (eds.)

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