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Pseudogap and superconductivity in cuprates, a dynamical mean-field perspective

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Ringberg symposium on Unconventional Superconductivity and Spin Liquids, 14-18 Oct. 2019



USHERBROOKE.CA/IQ

Phase diagram YBa₂Cu₃O_{7-x}





Zhao et al. Nat. Phys. 13, 250 (2017).

A W. A Yan- Par

Knight shift (Spin susceptibility)



Nakano *et al.* Phys. Rev. B **49**, 16000 (1994) Alloul *et al* (1989)





Figure from: Marc-Henri Julien



Model









Hubbard model



Attn: Charge transfer insulator

Method

- **Dynamical Mean Field Theory** - clusters
- Concept: atomic-like localized correlations consistent with delocalized aspect

REVIEWS

Maier, Jarrell et al., RMP. (2005) Kotliar et al. RMP (2006) AMST et al. LTP (2006)











Dynamical Mean-Field Theory. The beginnings in d = infinity



W. Metzner and D. Vollhardt, PRL (1989)A. Georges and G. Kotliar, PRB (1992)M. Jarrell PRB (1992)

DMFT, (d = 3)

2d Hubbard: Quantum cluster method



+ and -

- Long range order:
 - No mean-field factorization on the cluster
 - Symmetry breaking allowed in the bath (mean-field)
- Included exactly:
 - Short-range dynamical and spatial correlations
- Missing:
 - Long wavelength p-h and p-p fluctuations
 - Hence good when the correponding correlation lengths are small

Some groups using these methods for cuprates

- Europe:
 - Georges, Parcollet, Ferrero, Civelli, Wu (Paris)
 - Lichtenstein, Potthoff, (Hamburg) Aichhorn (Graz),
 Liebsch (Jülich) de Medici (Grenoble) Capone (Italy)
- USA:
 - Gull (Michigan) Millis (Columbia)
 - Kotliar, Haule (Rutgers)
 - Jarrell (Louisiana)
 - Maier, Okamoto (Oakridge)
- Japan
 - Imada (Tokyo) Sakai, Tsunetsugu, Motome

Outline

- The model
- The method
- Part I: T = 0 phase diagram
- Part II: The pseudogap from Knight shift
- Part III: Specific heat in the strange metal
- Part IV: Strongly correlated superconductivity
- Part V: Perspective



Part I

T = 0 phase diagram







T = 0 phase diagram

$$U = 8t, t' = -0.3t, t'' = 0.2t$$



A. Foley *et al.* Phys. Rev. B **99**, 184510 (2019)
S. S. Kancharla, *et al.* Phys. Rev. B **77**, 184516 (2008)
D. Sénéchal, *et al.* Phys. Rev. Lett. **94**, (2005)
M. Jarrell *et al.* EPL **56** 563, (2001)

CDMFT 4 sites

Fall at half-filling without AFM

t' = 0 DCA, 8 site



Gull et al. Phys. Rev. Lett. 110, 216405 (2013)



Part II:

The pseudogap









Simon Bergeron



Maxime Charlebois

B

Patrick Sémon



Alexis Reymbaut

R. Garioud

The pseudogap from Knight shift

A. Reymbaut, et al. Phys. Rev. Research 1, 023015 (2019)



Marion Thénault





Thanks: Marc-Henri Julien

Knight shift (Q=0 spin susceptibility)



Fig. 3 Temperature and doping dependence of the q = 0 spin susceptibility. At the smaller dopings (larger filling $\langle n \rangle$), $\chi_s(T)$ exhibits a peak in the temperature dependence indicating the opening of a PG

DCA 12 sites, *t*'=0, *U* = 7

T.A. Maier, D.J. Scalapino, npj Quantum Materials (2019)

Comparison



Fig. 3 Temperature and doping dependence of the q = 0 spin susceptibility. At the smaller dopings (larger filling $\langle n \rangle$), $\chi_s(T)$ exhibits a peak in the temperature dependence indicating the opening of a PG

Knight shift



DCA 8 sites, U = 6, t' = -0.1t

Chen, LeBlanc, Gull, Nature Com. Apr. 2017

See also Jarrell et al. 2001, 2002

Spin susceptibility





G.Sordi et al. Phys. Rev. B 87, 041101(R) (2013)





W Wu, A Georges, M Ferrero Phys. Rev. X 8, 021048 (2018). Bragança, Sakai, Aguiar, Civelli, PRL **120**, 067002 (2018)

Results : effect of t' on T*



Doiron-Leyraud *et al.* Nature Comm. **8** 2044



 $p^* < p_{fs}$

Doiron-Leyraud *et al.* Nature Comm. **8** 2044



A. Reymbaut *et al.* Phys. Rev. Research **1**, 023015 (2019)



A. Reymbaut *et al.* Phys. Rev. Research **1**, 023015 (2019)



A. Reymbaut *et al.* Phys. Rev. Research **1**, 023015 (2019)



Phys. Rev. Research 1, 023015 (2019)

Results : effect of t' on T*



Doiron-Leyraud et al.A.Reymbaut, et al.Nature Comm.8 2044Phys. Rev. Research 1, 023015 (2019)

Results: van Hove singularity



Doiron-Leyraud *et al.* Nature Comm. **8** 2044

A.Reymbaut, *et al.* Phys. Rev. Research **1**, 023015 (2019) Sordi et al., Sci. Rep. 2 547 (2012);

Physics: Plaquette eigenstates



U = 6.2; t' = 0

Sordi et al., Sci. Rep. 2 547 (2012);

See also:

Michel Ferrero, P. S. Cornaglia, L. De Leo, O. Parcollet, G. Kotliar, A. Georges PRB **80**, 064501 (2009)

Part III:

Specific heat in the strange metal phase











Maxime Charlebois

Alexis Reymbaut

Specific heat in the strange metal phase

A. Reymbaut, et al. Phys. Rev. Research 1, 023015 (2019)



Marion Thénault

R. Garioud

Specific heat in the strange metal phase



B. Michon, C. Girod, Taillefer, Klein, Nature 567, 218 (2019)

Specific heat in the strange metal phase





Part IV:

Strongly correlated superconductivity









Giovanni Sordi

Patrick Sémon

Lorenzo Fratino

Superconductivity for large U

Sordi et al. PRL 108, 216401 (2012) Fratino et al. Sci. Rep. 6, 22715 (2016)

Superconducting transition temperature



T.A. Maier, D.J. Scalapino, npj Quantum Materials (2019)









Olivier Simard

Charles-David Hébert

Alexandre Foley

David Sénéchal

What causes T_c to drop near n = 1?

O. Simard, C.-D. Hébert, A. Foley, A.-M.S. Tremblay, D. Sénéchal, Phys. Rev. B **100**, 094506 (2019)

What causes T_c to drop?

Phase fluctuations? Emery Kivelson Nature 374 (1995)



Uemura, Y.J. *et al.*, PRL vol.62, (1989) Tallon *et al.*, PRB **68**, 180501(R) (2003)

Superfluid stiffness T=0

$$j = -\rho_s A$$
 $d = 2$

$$\frac{1}{\lambda^2} = \rho_s \mu_0$$

$$T_c^{KT} = \frac{\pi}{8e^2} \rho_s(T_c^{KT})$$

$$T_c^{KT} < \frac{\pi}{8e^2} \rho_s(T_c^{KT} = 0)$$

 $\hbar = 1; k_B = 1$

Emery Kivelson, Nature **374**, 434 (1995) Metzner, Yamase, Phys. Rev. B **100**, 014504 (2019) Hazra, Verma, and Randeria, Phys. Rev. X **9**, 031049 (2019)



O. Simard, *et al.* Phys. Rev. B **100**, 094506 (2019) See also E. Gull, A.J. Millis, Phys. Rev. B **88**, 075127 (2013)



What energy scale controls Tc ?

T_c controlled by J 0.08 (a) 0.06 (b)U=5.6t < U_{MIT} (d) (e) -(c) -(f) (g) $U=6.2t > U_{MIT}$ 300 $\delta = 0$ U=7.0t U=12.0t U=16.0t U=9.0t 0.05 T(K) 0.06 0.04 T/t 200 0.04 0.03

0 0.05 0.10 0.15 0 0.05 0.10 0.15

δ

Fratino et al. Sci. Rep. **6**, 22715

5.0

U/t

4.0

6.0 0 0.05 0.10 0.15

δ

0.02

Some experiments that suggest $T_c < T_{pair} < T^*$ T. Kondo *et al.* PRL 111 (2013) Kondo, Takeshi, et al. Kaminski Nature Physics 2011, 7, 21-25 A. Pushp, Parker, ... A. Yazdani, Science **364**, 1689 (2009) Lee ...Tajima (Osaka) https://arxiv.org/pdf/1612.08830 Patrick M. Rourke, et al. Hussey Nature Physics **7**, 455–458 (2011) Lee et al. J. Phys. Soc. Jpn. 86, 023701 (2017)

0 0.05 0.10 0.15

δ

0.02

0.01

0.00

100

0

O COA-

0 0.05 0.10 0.15

δ

0 0.05 0.10 0.15

Condensation energy



Fratino et al. Sci. Rep. **6**, 22715

Theory, see also Jarrel PRL (2004), Gull Millis PRB (2014) Experiments: Bontemps, Santander-Syro Van der Marel ...



Part V:

Perspective









Giovanni Sordi

Kristjan Haule

Influence of the Mott transition away from half-filling

Sordi et al., PRL 104, 226402 (2010) Sordi et al., PRB 84, 075161 (2011) Fratino et al., PRB 93, 245147 (2016) [Emery model] Sordi et al., Sci. Rep. 2 547 (2012); Sordi et al., PRB 87, 041101(R) (2013)

AFM phase diagram d=2, t'=0



L. Fratino, P. Sémon, M. Charlebois, G. Sordi, AMT Phys. Rev. B 95, 235109 (2017)

Change in potential energy due to large ξ



L. Fratino,¹ P. Sémon,² M. Charlebois,² G. Sordi,¹ and A.-M. S. Tremblay^{2, 3} arXiv:1702.01821

Influence of Mott transition away from half-filling

n = 1, d = 2 square lattice



Influence of Mott transition away from half-filling

n = 1, d = 2 square lattice







G. Sordi et al. Phys. Rev. Lett. 108, 216401/1-6 (2012)
G.Sordi et al. Phys. Rev. B 87, 041101(R)/1-5 (2013)
P. Sémon, G. Sordi, *et al.*, Phys. Rev. B 89, 165113/1-6 (2014)

c-axis resistivity





K. Takenaka, K. Mizuhashi, H. Takagi, and S. Uchida, Phys. Rev.B 50, 6534 (1994).

G.Sordi et al. Phys. Rev. B 87, 041101(R)/1-5 (2013)

Connecting the finite doping behavior to the Mott transition at half-filling



Conclusion

p* in Hubbard is the end of Mott physics

Mott transition and its finite doping extension is the organizing principle













Figure from: Marc-Henri Julien





Compute • calcul

High Performance Computing

CREATING KNOWLEDGE DRIVING INNOVATION BUILDING THE DIGITAL ECONOMY

Le calcul de haute performance

CRÉER LE SAVOIR ALIMENTER L'INNOVATION BÂTIR L'ÉCONOMIE NUMÉRIQUE Calcul Québec

Fondation canadienne pour l'innovation

Merci Thank you



USHERBROOKE.CA/IQ





D. Sénéchal

Bumsoo Kyung

The glue

Kyung, Sénéchal, Tremblay, Phys. Rev. B **80**, 205109 (2009) Sénéchal, Day, Bouliane, AMST, Phys. Rev. B **87**, 075123 (2013) A. Reymbaut *et al.* PRB **94** 155146 (2016)

Im Σ_{an} and electron-phonon in Pb

Maier, Poilblanc, Scalapino, PRL (2008)



The glue



The glue and neutrons



FIG. 3 (color online). **Q**-integrated dynamic structure factor $S(\omega)$ which is derived from the wide-*H* integrated profiles for LBCO 1/8 (squares), LSCO x = 0.25 (diamonds; filled for $E_i = 140 \text{ meV}$, open for $E_i = 80 \text{ meV}$), and x = 0.30 (filled circles) plotted over $S(\omega)$ for LBCO 1/8 (open circles) from [2]. The solid lines following data of LSCO x = 0.25 and 0.30 are guides to the eyes.

Wakimoto ... Birgeneau PRL (2007); PRL (2004)

